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Searches for double Higgs production or decay using the CMS detector

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- 1. Introduction
- 2. Double Higgs searches in CMS:
	- A. bbyy
	- B. bbbb
	- C. bbWW
	- D. bb $\tau\tau$
- 3. Results and Conclusions

Motivations: Resonant searches ons: Resonant searc

MSSM/2HDM: Additional Higgs doublet→CP-even scalar H.

was an probe are fort fighters and region where bri_ller in • We can probe the low mA/low tan β region where BR(H \rightarrow h(125)h(125)) is sizeable.

Singlet Model. Additional Higgs singlet with an extra scalar H. Singlet model: Additional Higgs singlet with an extra scalar H.

□ sizable BR beyond 2xmtop, non negligible width at high mH • Sizeable BR beyond 2xmtop, non negligible width at high m_H.

Warped Extra Dimensions: spin-2 (Resonances of the 2 (Resonances of the 2 (Resonances of the 2 (Resonances of the 2 (Resonance of the 2 (Resonance of the 2 (Resonance of the 2 (Resonances of the 2 (Resonances of the 2 (Res

 $\sum_{i=1}^n a_i$ and spin-variation (bulk resonances $\sum_{i=1}^n a_i$ spin-2 (KK-graviton) and spin-0 (radion) resonances.

• Different phenomenology if SM particles are allowed (bulk RS) or not (RSI model) in the extra dimensional bulk

Motivations: Non-resonant searches

 $(1, 1, 1, 1)$ σ^{SM} _{hh}(13TeV) = 33.45fb^{+4.3%}-6.0%(scale unc.) ±3.1%(PDF+ α s unc)^[1]

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coupling (λ_{hhh}). Even if in Run2 we not have full sensitivity to "measure" SM λ_{hhh} $\prod_{i=1}^{n}$ rodud The non-resonant double Higgs production allows to directly probe the Higgs trilinear

 \rightarrow The BSM physics can be modelled in EFT adding dim-6 operators^[2] to the SM Lagrangian, and the physics can be described with 5 parameters: λ _{hhh}, y_t, c_{2,} c_{2g}, c_g the
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 \cdot New diagrams and couplings in the gam an sen
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- Non SM top Yukawa and λ_{hhh} couplings
- New diagrams and couplings in the game

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interactions

CMS searches

- 4 different searches presented today:
	- bbbb, bbWW, bb $\tau\tau$, bbyy
- At least one h→bb to have large enough BR $\frac{33}{17}$
- Rare processes, low σ , complex environment
- •Resonant and non-resonant searches performed in Run1 and Run2 • Run1:
	- bbbb Resonant: PLB 749 (2015) 560, arXiv:1602:08762
	- $$
	- bbyy Resonant and Non-resonant: arxiv:1603.06896
- Run2:
	- bbbb Resonant: PAS-HIG-16-002, PAS-B2G-16-008
	- **PEW bb T Resonant: PAS-HIG-16-029, AREW Non-resonant PAS-HIG-16-028 NEW Non-resonant PAS-HIG-16-028**
		- bbWW Resonant PAS-HIG-16-011, **NEW** Non-resonant: PAS-HIG-16-024

I will focus on the results at $\sqrt{s} = 13 \text{TeV}$

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Searches: how and where

- 3 Datasets used for this presentation:
- Run I, $\sqrt{s} = 8$ TeV, $\mathcal{L} = 17.9 19.7$ fb⁻¹
- Run2, 2015, \sqrt{s} =13 TeV, \mathscr{L} =2.3-2.7 fb⁻¹
- Run2, 2016, \sqrt{s} =13 TeV, \mathscr{L} =12.9 fb⁻¹
- B-tagging algorithm to identify b-jets from jet constituents
- CSVv2: Based on displaced tracks+secondary vertexes MVA[1]
- At high mH→boosted regime→merged jets
- **COV20 TIGGS**
• Reconstruction using substructure information for jets, b-tag
- \bullet bbbb, bb $\tau\tau$ channels
- Trade-off between BR and contamination, complementarity among channels
- \cdot bbbb: highest BR, high QCD/tt contamination
- \cdot bbWW: high BR, large irreducible $t\bar{t}$ background
- \cdot bb τ : relatively low background and BR
- bbyy: high purity, very low BR

[1]JINST 8(2013) P04013

hh→bb: run1 results

hh→bbbb

b-tagging at trigger level, \geq 4 b-jets offline

Low Mass Region (m_H <400) and High Mass Region $(400< m_H < 1200)$ studied separately

Background shape estimation from data in LMR, HMR ·

2 analysis strategies:

• *double b-tagger*: BDT from jet properties + background estimation from multiple sidebands •*subjet b-tag*: background fit + 3 categories based on number of b-tagged sub-jets

hh→bbbb: results

No evidence for the presence of new resonances so far over large mass range Sensitive to Radion (below 2TeV) and Graviton production (below 800GeV) Boosted analysis:

- double b-tagger: at low/high mass
- •sub-jets b-tagging: for 1200<mH<2000 GeV

hh→bbWW

- Search for $hh \rightarrow bbWW \rightarrow bb2l2\nu$.
- 2 isolated OS leptons + 2 b-jets in the final state
- 2015 dataset at \sqrt{s} =13 TeV
- Final BR for $bb212\nu$ final state: 1.22%
- Main backgrounds: $t\overline{t}$, DY, single top
- 2 BDT discriminants (h masses, angles, transverse mass) to separate signal from background at low $(m_H$ <450) and high mass $(m_H$ >450). Optimised for m_H =400 and m_H =650. I single BDT trained for non-resonant searches.
- Resonant: cut&count experiment in 4 categories: (mbb-peak,mbb-sidebands) x (low BDT, high BDT)
- Non resonant: 2D fit in [m(bb), BDT score] to extract the limits

hh→bbWW: results

Spin-2 RS1 KK-graviton excluded below 600 GeV Non-resonant analysis sensitive to O(400xSM)

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 m_{x}^{spin-0} (GeV)

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Observed 95% upper limit ected 95% upper limit Expected limit ± 1 std. deviation $\text{eted limit} \pm 2 \text{ std. deviation}$

 $radion$ (Λ _p=1.0 TeV, kL=35)

Type I Type II

 m_{12} =0 GeV, m =125 GeV, m =m_H=m 2HDM, tanβ=0.10, cos(β-α)=0.10

$hh \rightarrow bbb\tau\tau$

- Intermediate BR, fully reconstructed final state
- $|\tau_H+1|$ isolated leptons (e, μ , τ_H)+2 b-jets in the final state
- 3 final states: $e\tau_H$, $\mu\tau_H$, $\tau_H\tau_H$
- Main bkgs: $t\bar{t}$ (from MC), QCD multijet (from data in control regions)

Resonant search:

Limit extraction on kinematic fit of the 4-body invariant mass; 3 categories: 1bjet, 2bjet, boosted b-jets category

Non-resonant analysis:

- kinematic BDT discriminant to reduce $t\overline{t}$, only angular information
- visible mass as final variable

Only results on 2016 data shown. Results with 2015 data: CMS-PAS-HIG-012 CMS-PAS-HIG-013

hh→bb: results

Non resonant limits starts to make dents in part of the 5D EFT model phase space No significant excess observed in the resonant analysis

Summary of Run1 results

- Several analysis performed at CMS
- Coverage ranges from 2xmh to few TeVs
- hMSSM: Effective MSSM model with $m_h = m_{H0}$ [1]
- $H \rightarrow hh$ searches are providing an important coverage of the low mA/low tanβ region

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Summary of Run2 results

hh \rightarrow bbWW \rightarrow bb2l2 ν : 2015 data

(2.3/fb) at \sqrt{s} =13 TeV

• 2 mass regions, optimised for mH=400 and mH=650

hh \rightarrow bbbb: 2015 data (2.3/fb) at \sqrt{s} =13 TeV

•Boosted regime not shown hh \rightarrow bb $\tau\tau$: 2016 data (12.9/fb) at \sqrt{s} =13 TeV

Non-resonant production exclusion

Several competing analyses in different final states under study in CMS, providing excellent coverage in different decay modes.

- Non resonant double Higgs production is the main way to measure Higgs self-coupling.
	- At the moment, we can probe O(10-100xSM). Much larger luminosity is needed to reach SM sensitivity, but we are starting to probe BSM and to constraint exotic BSM

Resonant searches can already provide important constrain on BSM physics (MSSM, WED, heavy scalars).

• KK-graviton excluded below 800 GeV, ΛR=1TeV Radion below 2 TeV

Further improvement awaited with end-of-the-year luminosity and the combination of the results among all channels

Exciting prospects for double Higgs searches

BACKUP

gg→hh parametrization * *g* * *g* $M²$ and → hh narametrization factor into the suppression factor into the suppression factor in the coefficients. The coefficients of $\frac{1}{2}$ We thus obtain the following interactions in terms of the Higgs boson scalar *h*, relevant to α ² and → **hh** narametrization factor into the suppression factor into the suppression of α We thus obtain the following interactions in terms of the Higgs boson scalar *h*, relevant to $H = H + H + H + H + H$ **bigaria metrization** <u>L</u>
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The relevant lagrangian terms of gg→HH production in D=6 EFT **b** relevant I *h* **Letter Concerns** to the end of $\overline{\mathbf{r}}$ **h**
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\mathcal{L}_{hh} = -\frac{m_h^2}{2v} \left(1 - \frac{3}{2} c_H + c_6 \right) h^3 + \frac{\alpha_s c_g}{4\pi} \left(\frac{h}{v} + \frac{h^2}{2v^2} \right) G^a_{\mu\nu} G^{\mu\nu}_a \n- \left[\frac{m_t}{v} \left(1 - \frac{c_H}{2} + c_t \right) \bar{t}_L t_R h + \text{h.c.} \right] - \left[\frac{m_t}{v^2} \left(\frac{3c_t}{2} - \frac{c_H}{2} \right) \bar{t}_L t_R h^2 + \text{h.c.} \right]_{\text{arXiv:1410.3471}}
$$

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Higgs and the gluons denoted with a cross.

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FIG. 1: Cartoon of the region in the plane (*g*⇤*, /g*⇤), defined by Eqs. (13),(14), that can be probed

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Naively all the Wilson coefficients in Eq. (3.1) should be bounded from perturbativity ar-

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c $\left\{\begin{matrix} 1 \\ 0 \end{matrix}\right\}$

Implementation for hh 2924 we can take advantage of this property of the ratio between the ratio between the cross, and 200 2911 different EFT parameters and the SM cross section with the SM cross section with the \mathbb{R} **An EFT implementation for hh**

 H iggs production cross $\left\| R_{hh} \right\|$ section can be written as a function of the 5 EFT parameters: λ_{hhh}, y_t, e₂, c_{2g}, c_g *hh* σ_{hh}^{SM} The double Higgs production cross $\frac{R_{hh} \equiv \frac{\sigma_{hh}^{S}}{\sigma_{hh}^{S}}}{\sigma_{hh}^{S}}$

LO $\stackrel{LO}{=} A_1 \kappa_t^4 + A_2 c_2^2 + (A_3 \kappa_t^2 + A_4 c_g^2) \kappa_\lambda^2 + A_5 c_{2g}^2 + (A_6 c_2 + A_7 \kappa_t \kappa_\lambda) \kappa_t^2$ *t*

 $+(A_8\kappa_t\kappa_\lambda+A_9c_g\kappa_\lambda)c_2+A_{10}c_2c_{2g}+(A_{11}c_g\kappa_\lambda+A_{12}c_{2g})\kappa_t^2$ *t* $+(A_{13}\kappa_{\lambda}c_q + A_{14}c_{2q})\kappa_t\kappa_{\lambda} + A_{15}c_qc_{2q}\kappa_{\lambda}.$

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 \mathbf{S} ^r) signal snapes from different points in the 5D EFT phase space are clustered together.

12 clusters are identified according to represented by a unique "shape benchmark" benchmark there kinematical properties

 $\frac{1}{2}$ cluster, a representative shape is identified, as the one with the minimum distance (in the test statistics). from all other shapes in the cluster Inside each cluster, a representative

 Luca Cadamuro (LLR) CMS central approval 27/07/2016 8 Each point of the phase space can be mapped by means of its cross-section and representative shape

hh→bb

- Lowest BR of all channels considered, but excellent resolution on myy
- Selection: p_T^{γ} ¹/m_{$\gamma\gamma$} < 1/3, p_T^{γ} ²/m $\gamma\gamma$ < 1/4 + mass cuts
- Two categories: $1b$ -jet (low purity category), ≥ 2 btagged jets (high purity)
- Different signal regions at low and high mass
- 2 \cdot m_H<400 GeV: b-jet regression $+$ 2D signal fit $\pmb{\mathsf{Q}}$ extraction based on $(m_{jj}, m_{\gamma\gamma})$ 19.7 fb⁻¹ (8 TeV Events / 20 GeV Ge **CMS** $pp \rightarrow X \rightarrow HH \rightarrow \gamma \gamma bb$ 12 \cdot 400 < m_H < 1100 GeV: **GeV** Events / 10 GeV Medium-purity 35 Events / 20 \rightarrow Data \rightarrow Background model kinematic fit of the 10 68% CL 95% CL \overline{C} 30 4-body invariant mass 8 Events 25 6 20 Non-resonant analysis: 15 4 10 2 2 b-tag cat. X 2 cat. myy 5 • $m_{\gamma\gamma}$ <350 && $|cos\theta^{CS}|$ < 0.65 0 400 500 600 700 800 900 1000 1100 1200 $m_{\gamma \gamma \text{ii}}^{\text{kin}}$ (GeV) $^{0-}_{60}$ • $m_{\gamma\gamma}$ > 350 && $|cos\theta^{CS}|$ <0.9

hh→bb

- Lowest BR of all channels considered, but excellent resolution on myy
- Selection: p_T^{γ} ¹/m_{$\gamma\gamma$} < 1/3, p_T^{γ} ²/m $\gamma\gamma$ < 1/4 + mass cuts
- Two categories: 1b-jet (low purity category), ≥ 2 btagged jets (high purity)
- Different signal regions at low and high mass
- m_H <400 GeV: b-jet regression $+$ 2D signal fit extraction based on $(m_{ii}, m_{\gamma\gamma})$
- \cdot 400 <m H < 1100 GeV: kinematic fit of the 4-body invariant mass

Non-resonant analysis:

2 b-tag cat. X 2 cat. m $\gamma\gamma$ • $m_{\gamma\gamma}$ <350 && $|cos\theta^{CS}|$ < 0.65 • $m_{\gamma\gamma}$ > 350 && $|cos\theta^{CS}|$ <0.9

