

# Directed flow at midrapidity in $\sqrt{s}=2.76\text{ TeV}$ PbPb collisions\*



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\* [ARXIV:1203.0931v1](https://arxiv.org/abs/1203.0931v1), TO APPEAR IN PRL SOON!

# Outline:

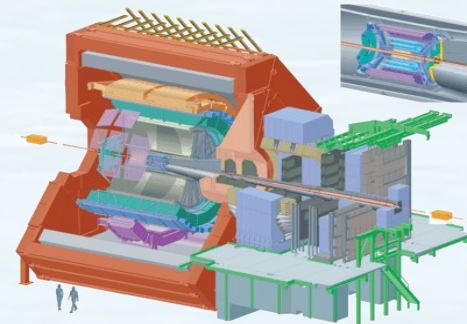
- ✓ Dihadron correlations and factorisation
- ✓ Momentum conservation coefficient: fit parameter
- ✓ Momentum conservation coefficient: estimation
- ✓ Viscous hydro calculations: LHC and RHIC calculations
- ✓ Conclusions



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# Measurements at ALICE



DISTRIBUTIONS OF ANGLES  $\Delta\phi$  AND/OR  $\Delta\eta$  BETWEEN:

- A “TRIGGER” PARTICLE AT TRANSVERSE MOMENTUM  $P_T^T$
- AN “ASSOCIATED” PARTNER AT  $P_T^A$

TWO-PARTICLE CORRELATIONS:

$$V_{n\Delta} = \langle \cos n(\Delta\phi) \rangle$$



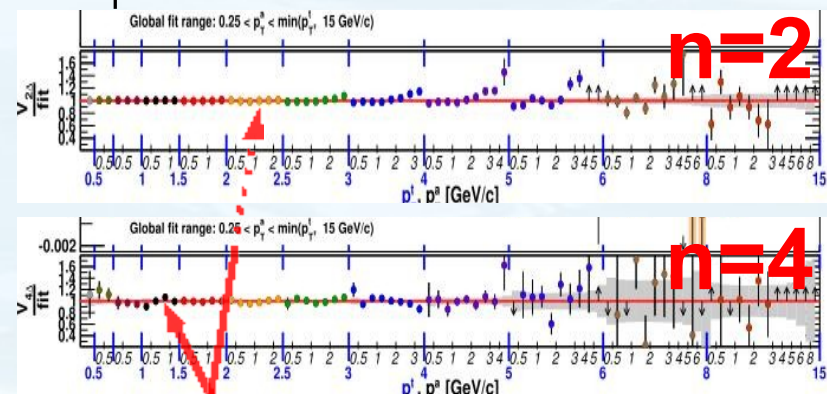
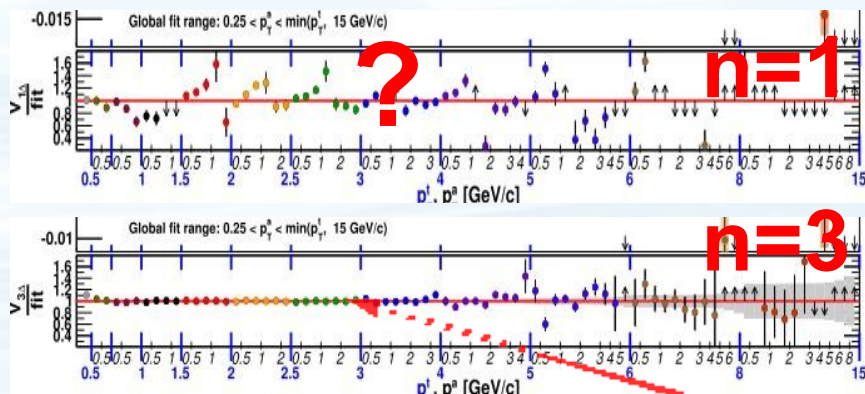
# ALICE analysis

IT WAS FOUND RECENTLY, THAT TWO-PARTICLE CORRELATION  
FACTORIZES IN LONG-RANGE CORRELATIONS WITH  $|\Delta\eta| > 0.8$ :

$$V_{n\Delta} = V_n(p_T^+) * V_n(p_T^a)$$

ALICE: FIT  $N \times N$  MATRIX WITH **N** PARAMETERS OF  $V_n$ :

N-NUMBER OF  $p_T$  BINS



ARXIV:1109.2501v2

FIT IS GOOD FOR  $n > 1$

# Factorization

$$V_{n\Delta} = v_n(p_T^t) v_n(p_T^a)$$

How do we understand this?

PARTICLES ARE EMITTED INDEPENDENTLY WITH DISTRIBUTION:

$$\frac{dN}{d\phi} = \frac{N}{2\pi} (1 + 2 \sum v_n \cos(n\phi - n\psi_n))$$

WHERE  $\langle e^{in\phi} \rangle = v_n e^{in\psi_n}$

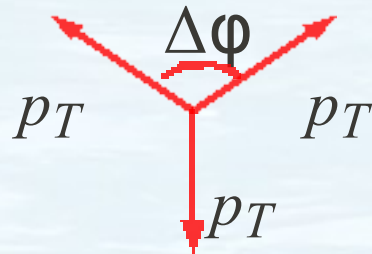
$$V_{n\Delta} = \langle \cos n\Delta\phi \rangle = \langle e^{in(\phi_t - \phi_a)} \rangle = \langle e^{in\phi_t} \rangle \langle e^{-in\phi_a} \rangle = v_n^a v_n^t$$



# Momentum conservation

Factorization doesn't work for  $n=1$ :

We add one nonflow term due to global momentum conservation



$$\langle \cos(\Delta\phi) \rangle_{mom.cons.} = -k p_T^t p_T^a < 0$$

$$\sum \vec{p}_T = 0$$

momentum  
conservation

$$V_{1\Delta} = v_1(p_T^t) v_1(p_T^a) - \boxed{k p_T^t p_T^a}$$

Two possibilities to find  $k$ :

→ As fit parameter

→ Calculate it as  $1/\langle \sum p_T^2 \rangle$

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ARXIV:NUCL-TH/0004026v2

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# Comparison of two fit functions:

OUR FIT WITH **N+1**

PARAMETERS ( **$v_1 + \kappa$** ):

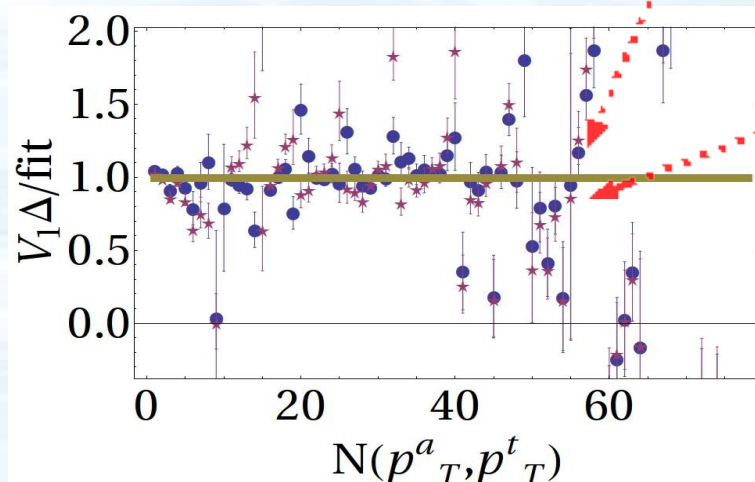
$$V_{1\Delta} = v_1(p_T^t) v_1(p_T^a) - \kappa p_T^t p_T^a$$

ALICE FIT WITH **N**

PARAMETERS ( **$v_1$** ):

VS

$$V_{1\Delta} = v_1(p_T^t) v_1(p_T^a)$$

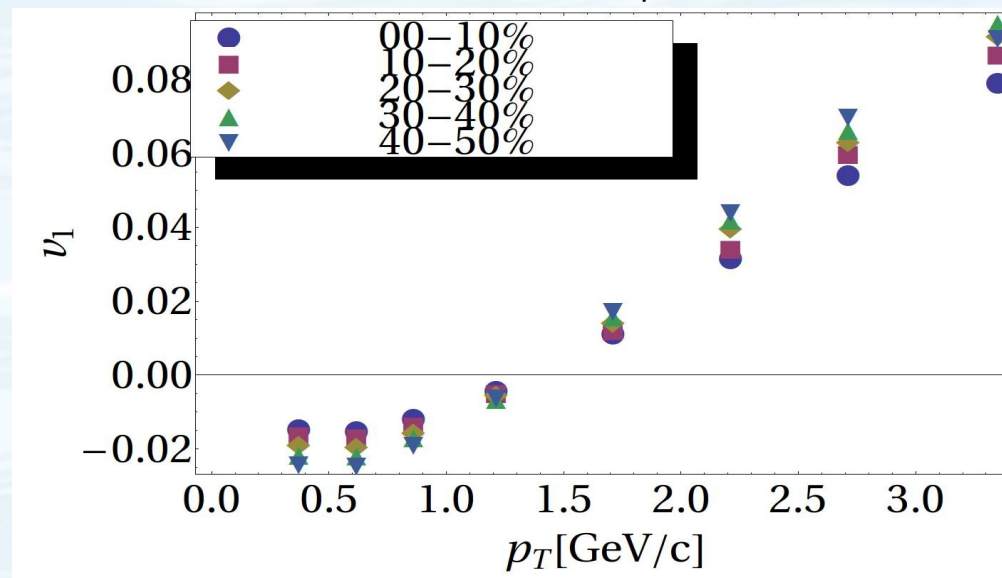


centrality	$\chi^2 / \text{d.o.f.}, \kappa=0$	$\chi^2 / \text{d.o.f.}$
0-10%	6	2
10-20%	17	1.7
20-30%	45	2.2
30-40%	75	2.3
40-50%	126	2.5

The quality of the fit is much better!

# First measurement of $v_1$ at the LHC

EXTRACTED VALUES FROM THE FIT GIVE US THE FIRST  
MEASUREMENT OF  $v_1$  AT THE LHC!



NO NET TRANSVERSE MOMENTUM  $\rightarrow$  LOW  $p_T$  PARTICLES FLOW IN THE  
OPPOSITE DIRECTION TO HIGH  $p_T$  !



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# How to estimate

$$k = 1 / \langle \Sigma p_T^2 \rangle ?$$

## What we know:

- $p_T$  SPECTRA OF  $\pi, K, p$  AT MIDRAPIDITY IN A LIMITED  $p_T$  RANGE
- TOTAL CHARGED MULTIPLICITY  $N_{CH}$ : EXTRAPOLATION MADE BY ALICE



SUM RUNS OVER  
**ALL** THE  
PARTICLES!

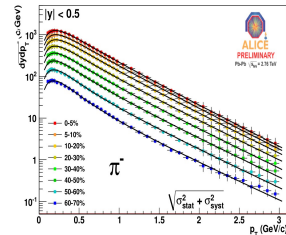
## What we don't know:

- ALICE DOESN'T MEASURE NUMBER OF NEUTRAL PARTICLES
- $p_T$  SPECTRA OUTSIDE MIDRAPIDITY



# Calculating k

arXiv:1109.2501v2



- FIT  $p_T$  SPECTRA BY LEVY FUNCTION TO EXTRAPOLATE FROM 0 TO  $\infty$

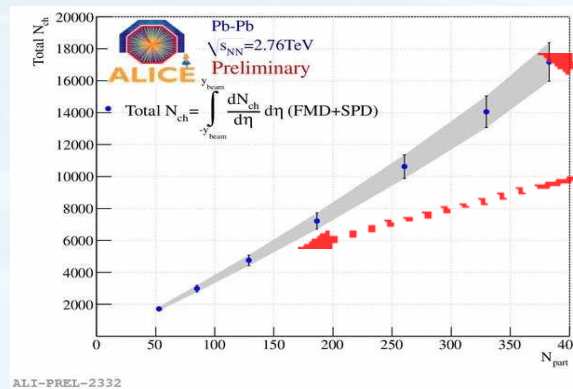
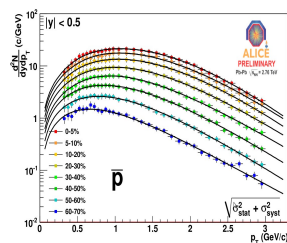
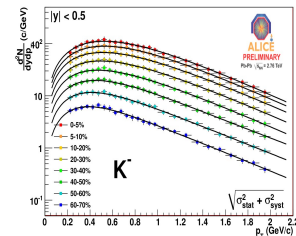
$$\frac{dN}{dy dp_T} = \frac{dN}{dy} \frac{p_T \cdot (n-1) \cdot (n-2)}{(n \cdot C \cdot (n \cdot C + m \cdot (n-2)))} \cdot \left(1 + \frac{(\sqrt{p_T^2 + m^2} - m)}{(n \cdot C)}\right)^{-n}$$

- INTEGRATE FUNCTION TO GET  $dN/dy$ ,  $\langle p_T^2 \rangle$

$$C, \frac{dN}{dy} \text{ } n\text{-Levy parameters}$$

$$\langle p_T^2 \rangle = \frac{\int p_T^2 \frac{dN}{dp_T dy} dp_T}{\int \frac{dN}{dp_T dy} dp_T}$$

- NEUTRAL PARTICLES ARE TAKEN INTO ACCOUNT ASSUMING TO ISOSPIN SYMMETRY



total  $N_{ch}$

arXiv:1107.1973v1

# Comparison of $k$ -coefficients

centrality	$k \text{ fit}, \times 10^{-5}$	$k \text{ est}, \times 10^{-5}$
0-10%	2.5	6.1
10-20%	4.7	8.8
20-30%	10.3	13
30-40%	21	21
40-50%	42	35



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# Types of $v_1$

Corresponding to:

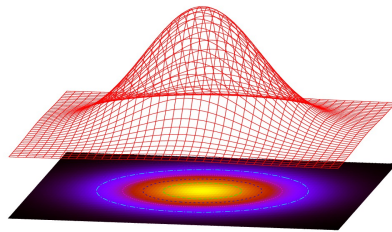
- event-plane method, **odd** in rapidity - **already studied**
- fluctuations in energy-density profile, **even** in rapidity - **our study**



# Next step: viscous hydro

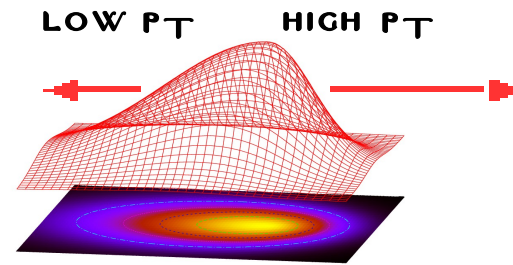
WE USE A SMOOTH, SYMMETRIC DENSITY PROFILE WHICH WE DEFORM TO INTRODUCE A DIPOLE ASYMMETRY  $\varepsilon_1$  OF THE DESIRED SIZE AND ORIENTATION.

$$\varepsilon_1 = \frac{\left| \left\{ r^3 e^{i\varphi} \right\} \right|}{\left\{ r^3 \right\}}$$



**ENERGY DENSITY PROFILE  
WITHOUT FLUCTUATIONS**

FLUCTUATIONS



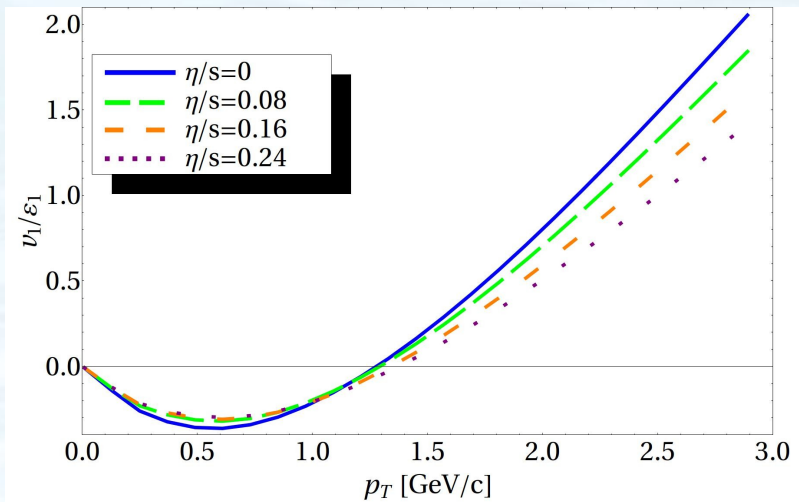
**ENERGY DENSITY PROFILE  
WITH FLUCTUATIONS**

OUR CALCULATION IS A 2+1D VISCOUS HYDRODYNAMIC USES AS INITIAL CONDITION THE TRANSVERSE ENERGY DENSITY ( $\epsilon(r, \phi)$ ) PROFILE FROM AN OPTICAL GLAUBER MODEL:

$$\boxed{\epsilon(r, \phi) \rightarrow \epsilon(r \sqrt{1 + \delta \cos(\phi - \Psi_1)}, \phi)}$$

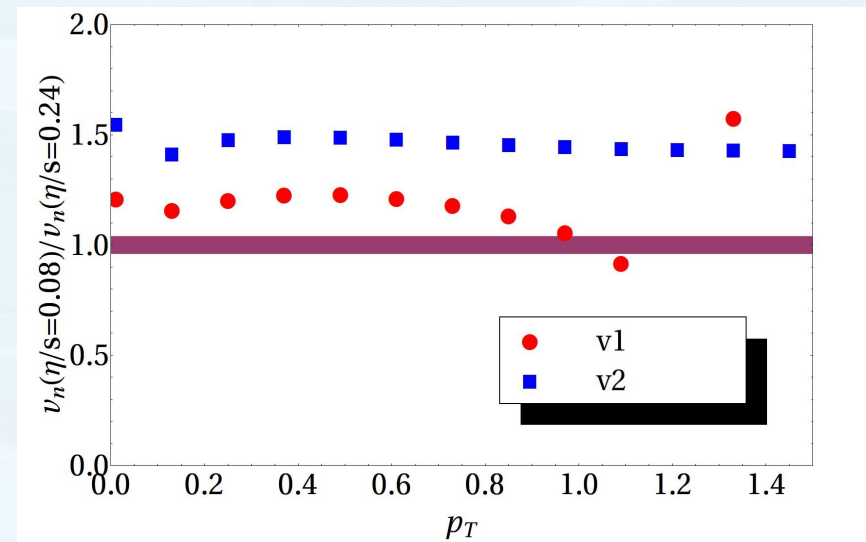
$$\delta \ll 1 \quad V_1 \sim \varepsilon_1 \sim \delta$$

# Dependence on viscosity of $v_n/\epsilon_n$



$v_1$  HAS A WEAKER DEPENDENCE ON  
VISCOSITY THAN  $v_2$

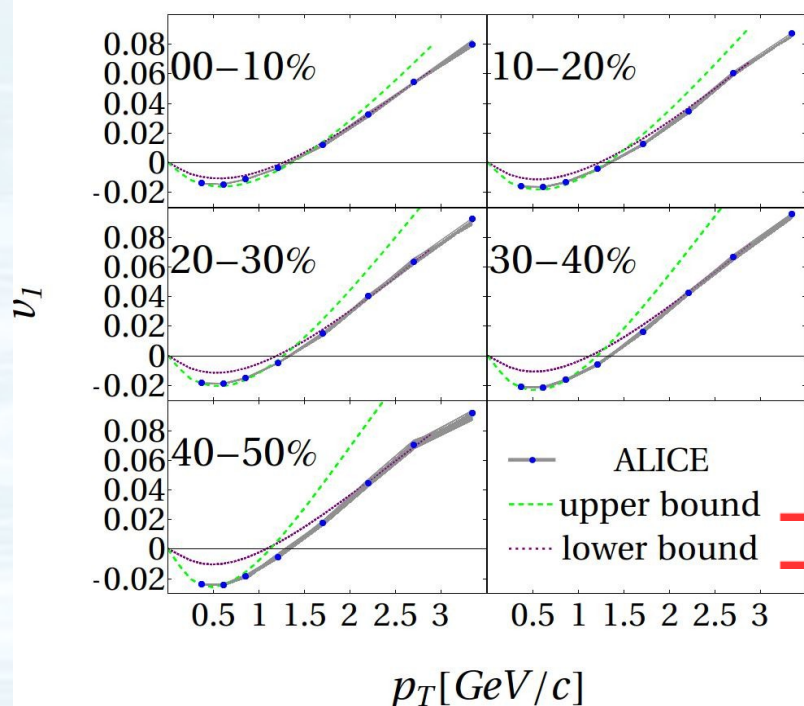
$v_1/\epsilon_1 (p_T)$





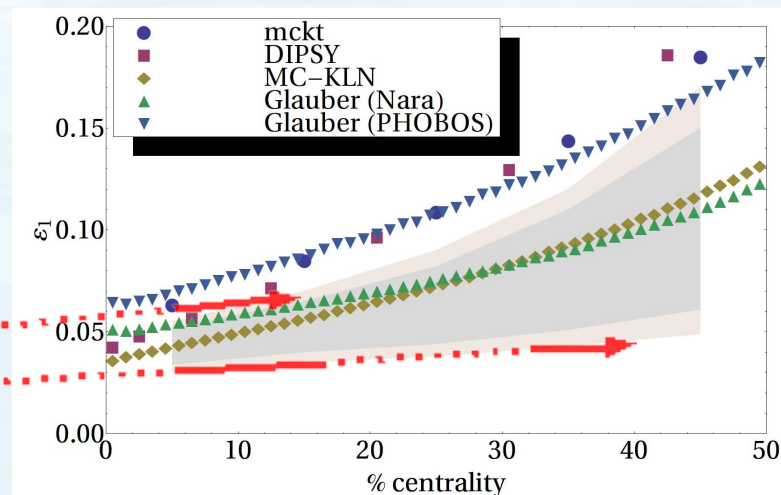
# Hydro+experimental data

→ CHOOSE  $\epsilon_1$  TO MATCH THE DATA FROM BELOW OR ABOVE



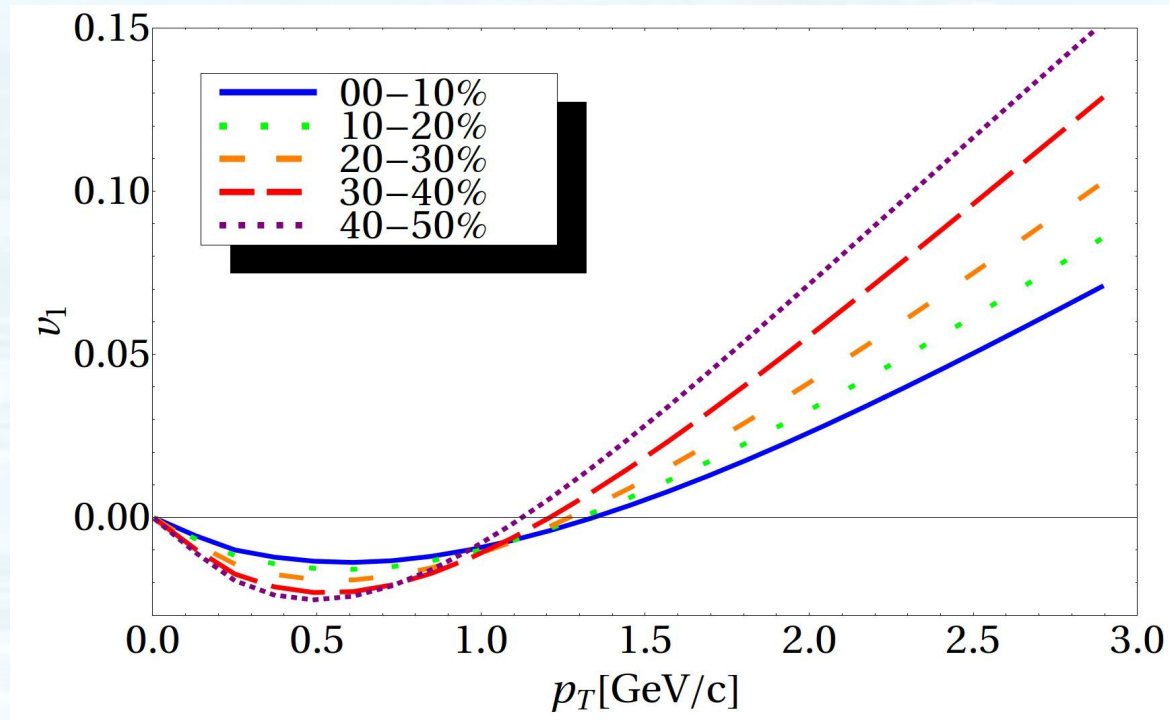
WITH HYDRO+EXPERIMENTAL DATA, WE CAN CONSTRAIN  $\epsilon_1$

→ COMPARISON OF MONTE-CARLO MODELS



THROUGH  $\epsilon_1$ , WE CONSTRAIN MODELS OF INITIAL STATE FLUCTUATIONS

# RHIC prediction for $v_1$



$$\epsilon_{1\text{LHC}} * \left( \frac{v_1}{\epsilon_{1\text{RHIC}}} \right)$$



# Conclusions:

- ▶ first measurement of directed flow,  $v_1$ , at midrapidity at the LHC,

similar analysis later by ATLAS: [arXiv:1203.3087v2](#)

- ▶ first viscous hydrodynamic calculation of directed flow

- ▶  $v_1$  depends less on viscosity than  $v_2$  and  $v_3$

- ▶ data on  $v_1$  constrain the fluctuations of the early-time system  $\rightarrow$  rule out certain current theoretical models

- ▶ first prediction made for directed flow at midrapidity in lower-energy collisions at RHIC

*Backup slides*



# Estimated value of $k$

estimated value  $k$ :  $k = \frac{1}{\langle \Sigma p_T^2 \rangle}$

$$\langle \Sigma p_t^2 \rangle = N_{ch} \cdot \left( \frac{3 \cdot [\langle p_t^2 \rangle dN/dy]_{pion} + 4 \cdot [\langle p_t^2 \rangle dN/dy]_{kaon} + 4 \cdot [\langle p_t^2 \rangle dN/dy]_{proton}}{2 \cdot [dN/dy]_{pion} + 2 \cdot [dN/dy]_{kaon} + 2 \cdot [dN/dy]_{proton}} \right)$$

Total  $N_{ch}$

