

Large anisotropies in the Little Bang

Heavy ion seminar
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Li Yan, JYO, arXiv:1312.6555, PRL 112 (2014) 082301
Li Yan, JYO, Art Poskanzer, in preparation

Anisotropic flow

- Particles are emitted with a *probability distribution* that is not isotropic in azimuthal angle

$$P(\phi) = 1 + 2 \sum_{n>0} v_n \cos(n(\phi - \Psi_n))$$

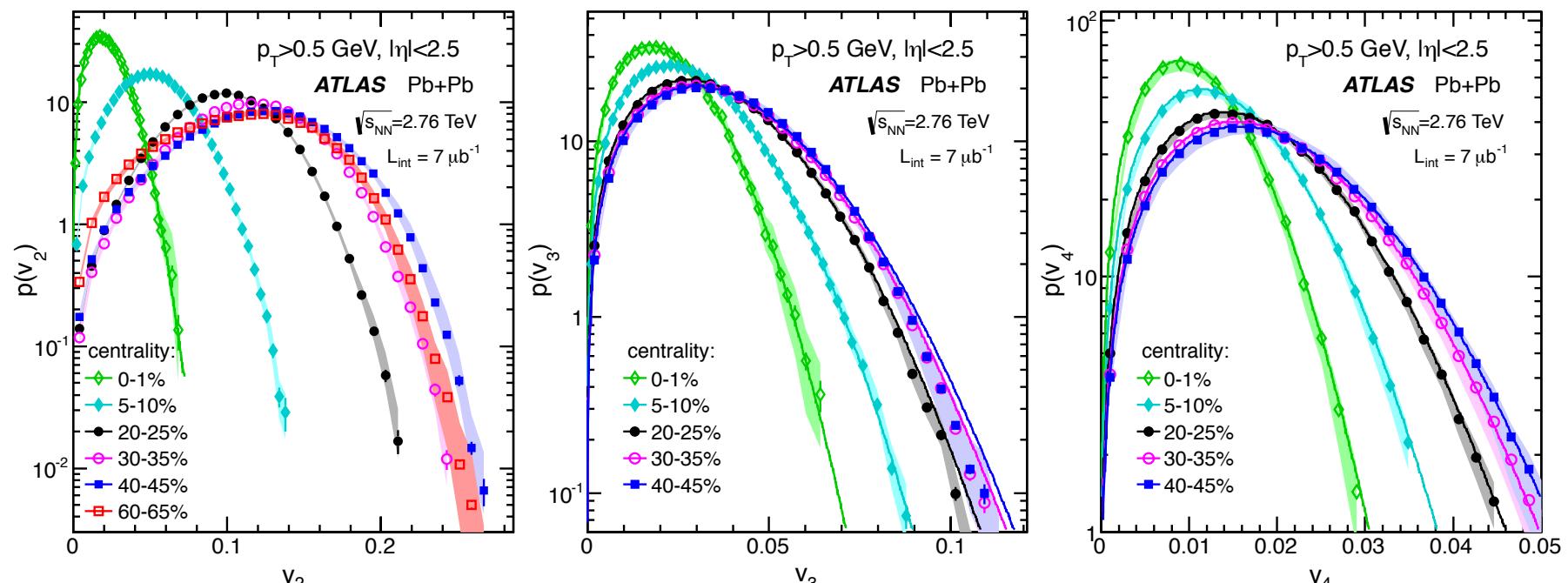
- v_n ≡ **anisotropic flow**
 v_2 ≡ **elliptic flow**
 v_3 ≡ **triangular flow...**
- Finite number of particles → trivial anisotropies from statistical fluctuations.
- v_n can be measured only after statistical fluctuations are subtracted (“unfolded”)

Flow fluctuations

- v_n fluctuates event to event
(PHOBOS, 2005)
- v_n itself has a *probability distribution* for a given system and centrality.

New data in Pb-Pb

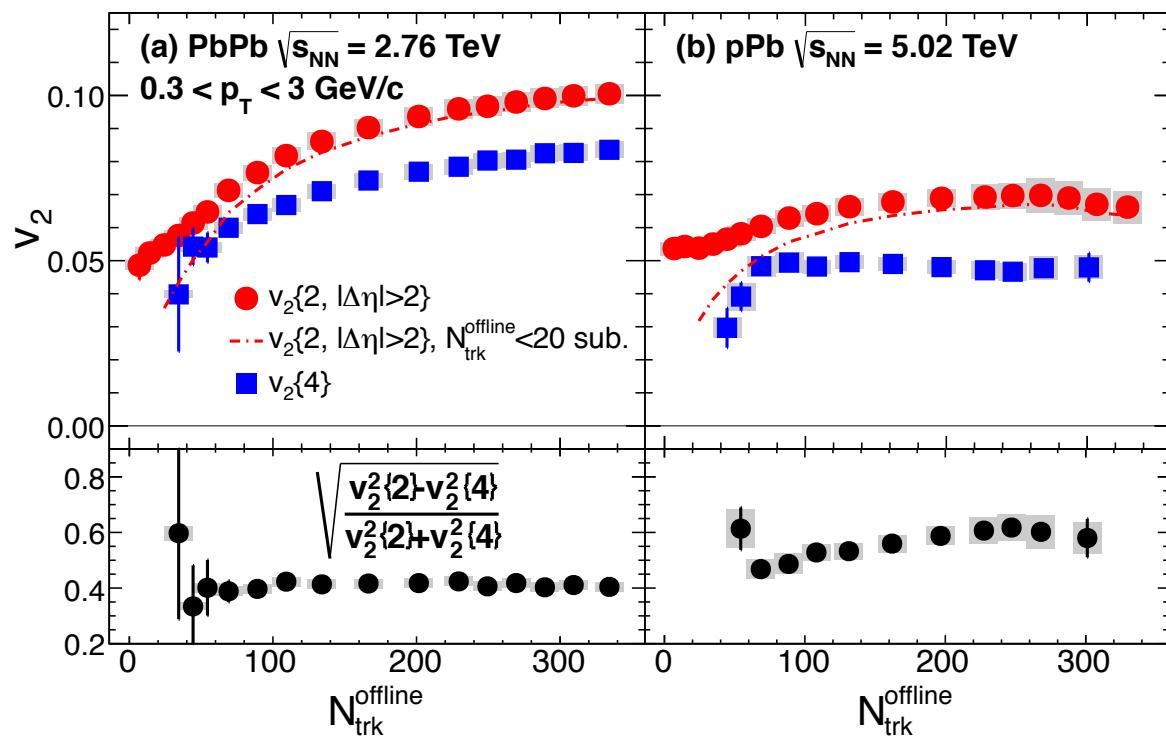
The probability distribution of v_2 , v_3 , v_4
for various centralities



ATLAS 1305.2942

New data in p-Pb

First 2 cumulants of the distribution of v_2
(less detailed than the full distribution)



$$v_2\{2\} \equiv (\langle v_2^2 \rangle)^{1/2}$$

$$v_2\{4\} \equiv (2\langle v_2^2 \rangle^2 - \langle v_2^4 \rangle)^{1/4}$$

If v_2 doesn't fluctuate,
 $v_2\{2\} = v_2\{4\} = v_2$

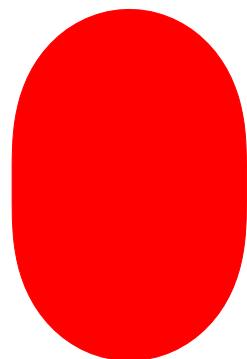
In general $v_2\{4\} < v_2\{2\}$

CMS 1305.0609

- Do we understand these new data?
- What can we learn from them?

The origin of anisotropic flow

*Initial transverse
density profile*

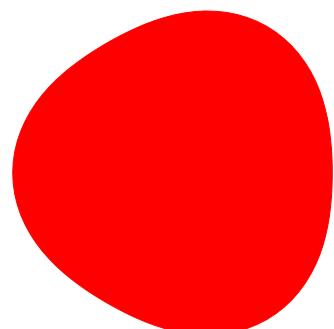


Expansion



Final distribution

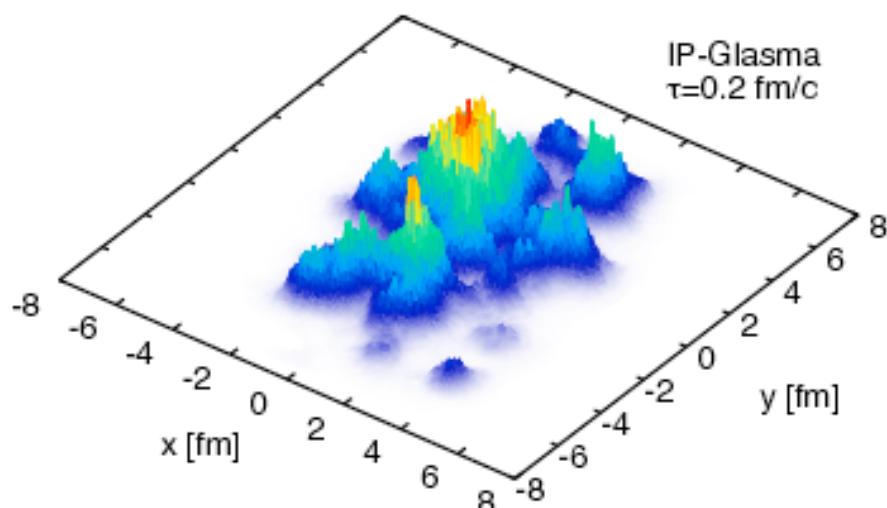
Elliptic flow v_2



Triangular flow v_3

Initial anisotropies

= Fourier decomposition of the initial density profile $\rho(x,y)$



$$\varepsilon_n \equiv \frac{\left| \int r^n e^{in\phi} \rho(r, \phi) r dr d\phi \right|}{\int r^n \rho(r, \phi) r dr d\phi}$$

$\varepsilon_2 \equiv$ initial eccentricity

$\varepsilon_3 \equiv$ initial triangularity

Gale Jeon Schenke 1301.5893

$|\varepsilon_n| < 1$ by definition

Anisotropic flow \approx initial anisotropy

$$v_2 \approx K_2 \epsilon_2$$

$$v_3 \approx K_3 \epsilon_3$$

response coefficients

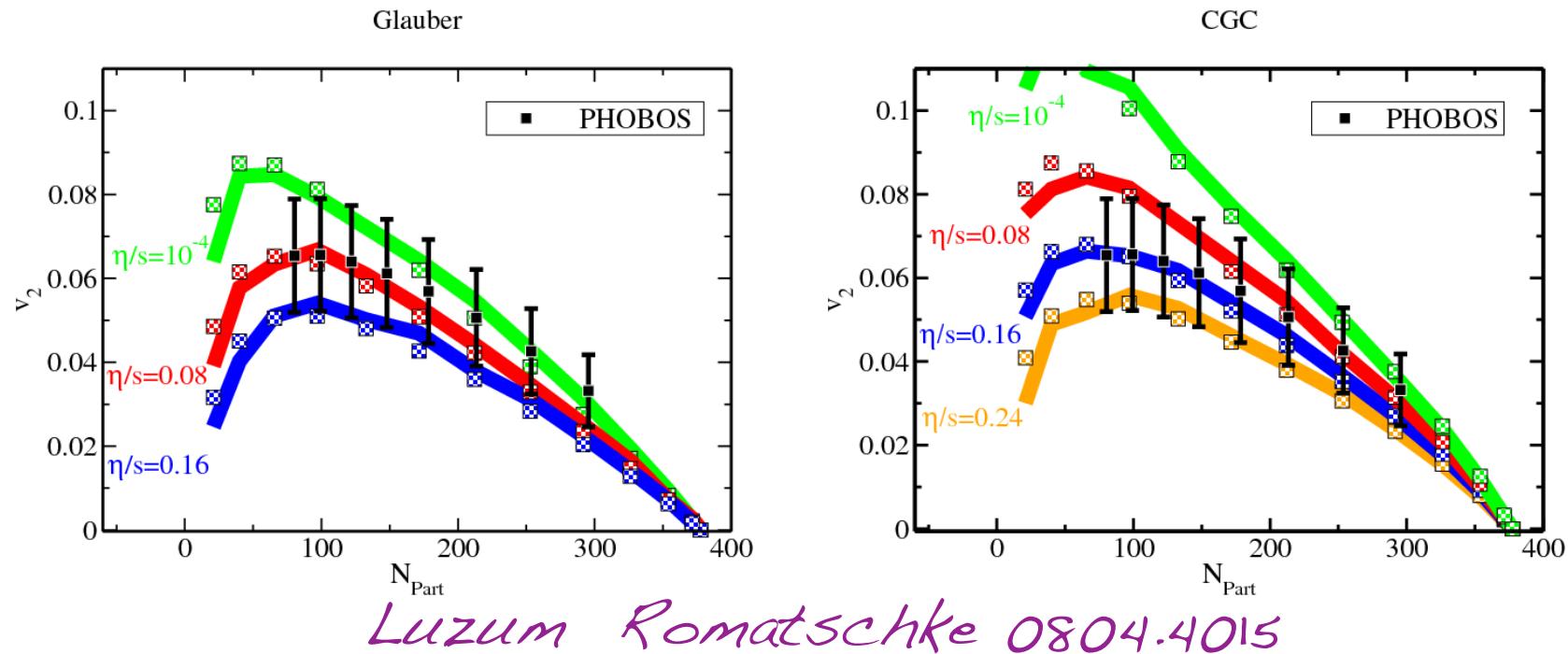
fluctuate event to event.

depend on system and centrality

in hydro, depend on viscosity

v_n fluctuations are due to ϵ_n fluctuations

Problem: can we disentangle the initial anisotropy from the response?



A long-standing problem in heavy-ion physics:
for any model of initial conditions (Glauber and CGC), i.e.,
for any ϵ_n , one can tune the viscosity — the response K_n —
to match the observed v_n

Is there a general law that describes anisotropy fluctuations?

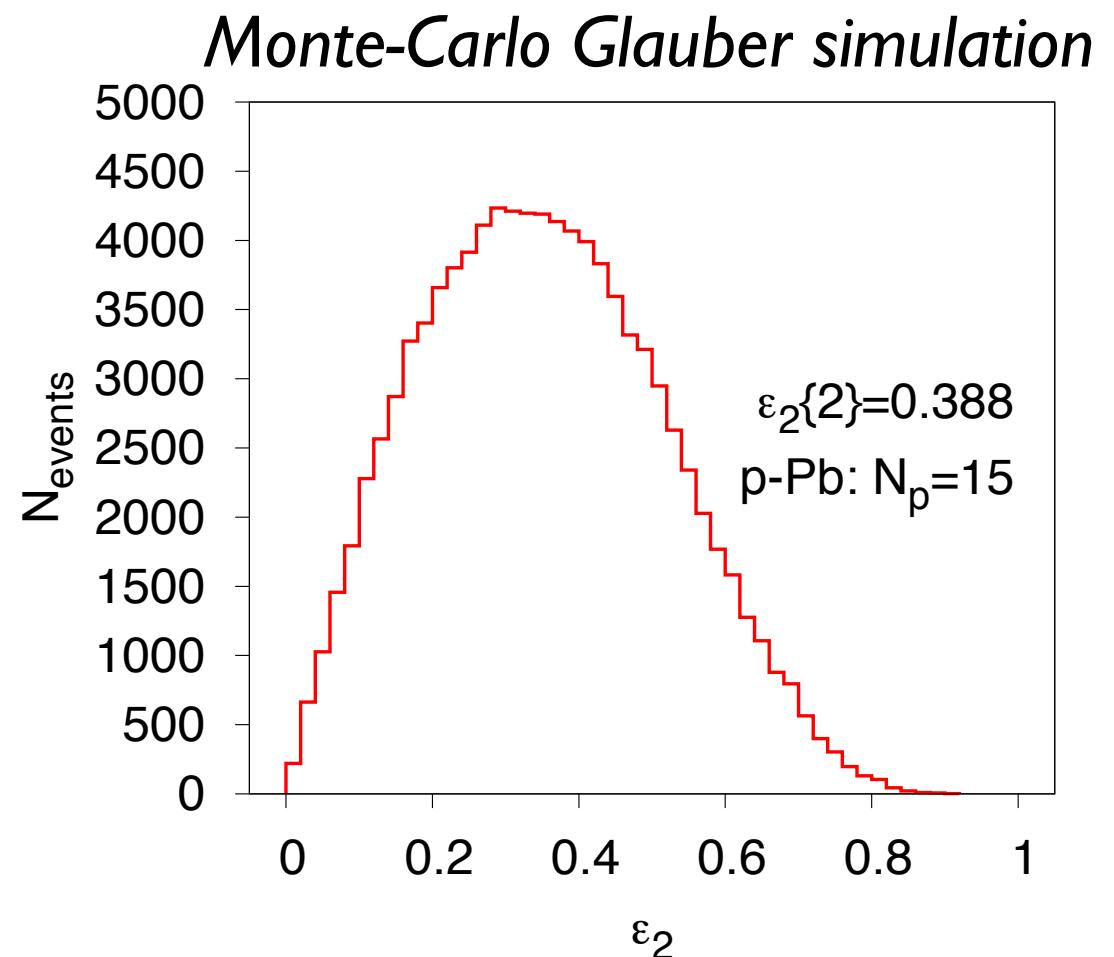
- If we know the statistics of the initial ε_n , then the distribution of observed v_n is the distribution of ε_n , rescaled by the response K_n
- State of the art (as of 2013): Gaussian fluctuations
 $P(\varepsilon_n) \propto \varepsilon_n \exp(-\varepsilon_n^2/\sigma^2)$ *Voloshin et al 0708.0800*
- Then the distribution of v_n is also a Gaussian, of width $K_n \times \sigma$: we are still unable to disentangle the initial state from the response.

The statistics of initial fluctuations

$$\varepsilon_2 = \frac{|\int r^2 e^{2i\phi} \rho(r, \phi) r dr d\phi|}{\int r^2 \rho(r, \phi) r dr d\phi}$$

central p+Pb collision:
initial density $\rho(r, \phi)$ =
independent of ϕ up to
fluctuations

small system: large
fluctuations & *anisotropies*



Is there a simple law that describes this distribution?

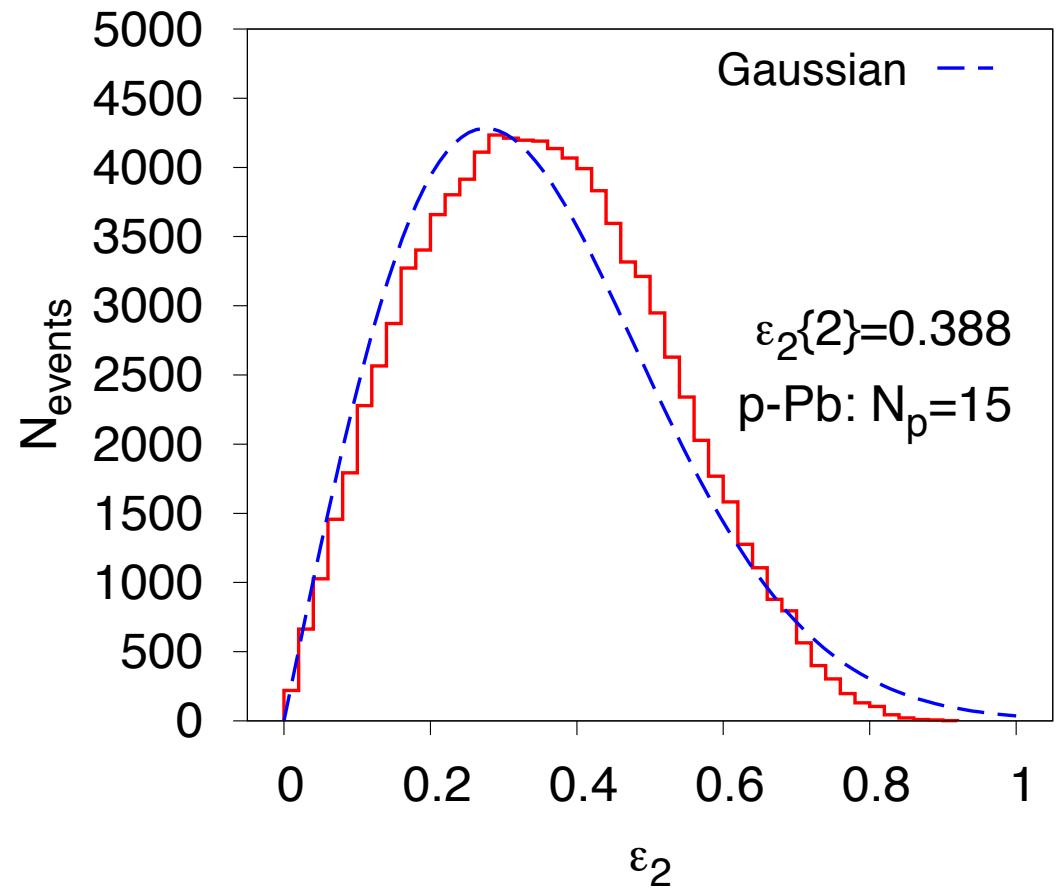
Gaussian?

Central limit theorem

$$P(\varepsilon_2) = 2(\varepsilon_2/\sigma^2) \exp(-\varepsilon_2^2/\sigma^2)$$

Not a good fit.

Does not implement
the condition $\varepsilon_2 < 1$



New “Power” distribution

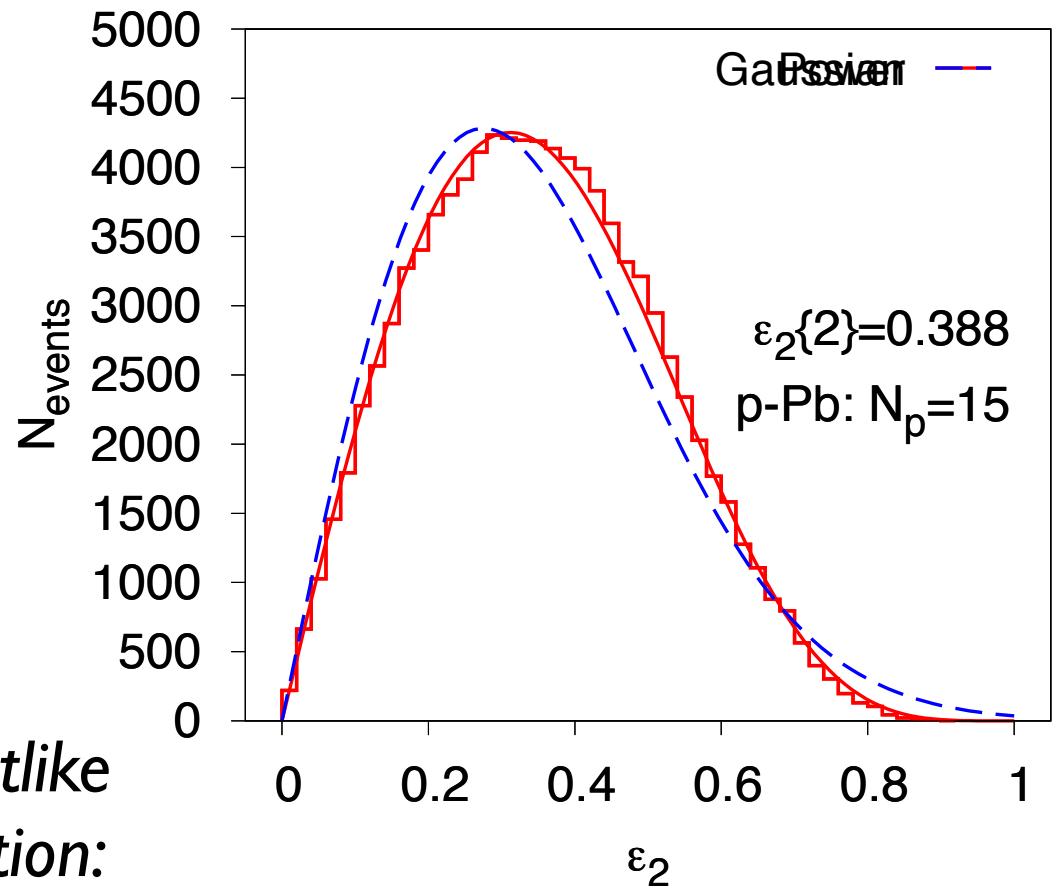
$$P(\varepsilon_2) = 2\alpha\varepsilon_2(1-\varepsilon_2^2)^{\alpha-1}$$

Equivalent to Gaussian for
 $\alpha \gg 1$

Naturally implements the
condition $\varepsilon_2 < 1$.

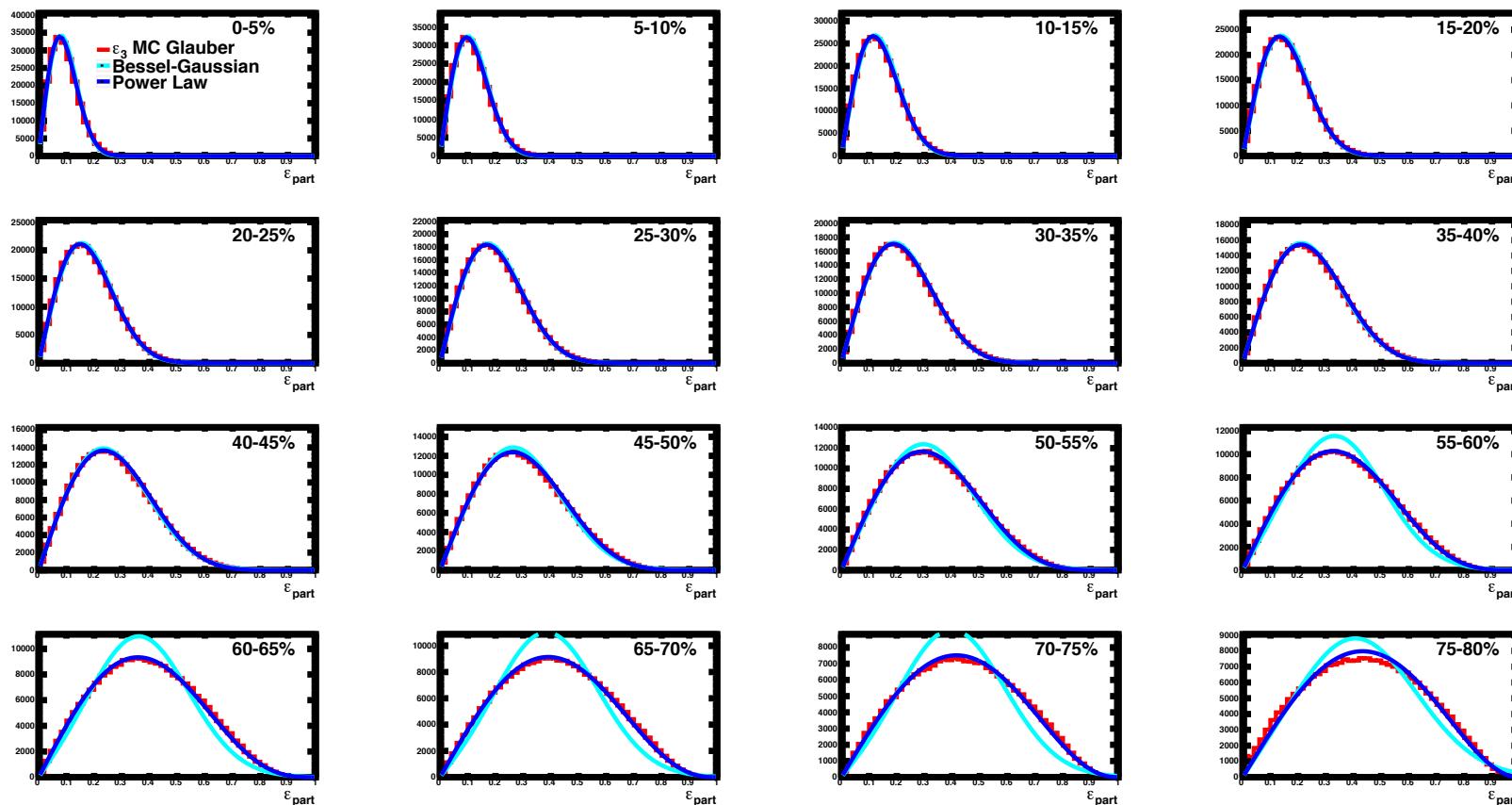
Exact result for $N=2\alpha+1$ pointlike
sources with Gaussian distribution:

JYO, PRD 46 (1992) 229



Much better fit to Monte-Carlo results!

Testing the Power distribution for ε_3 in Au-Au collisions



fits to Monte-Carlo Glauber by Art Poskanzer

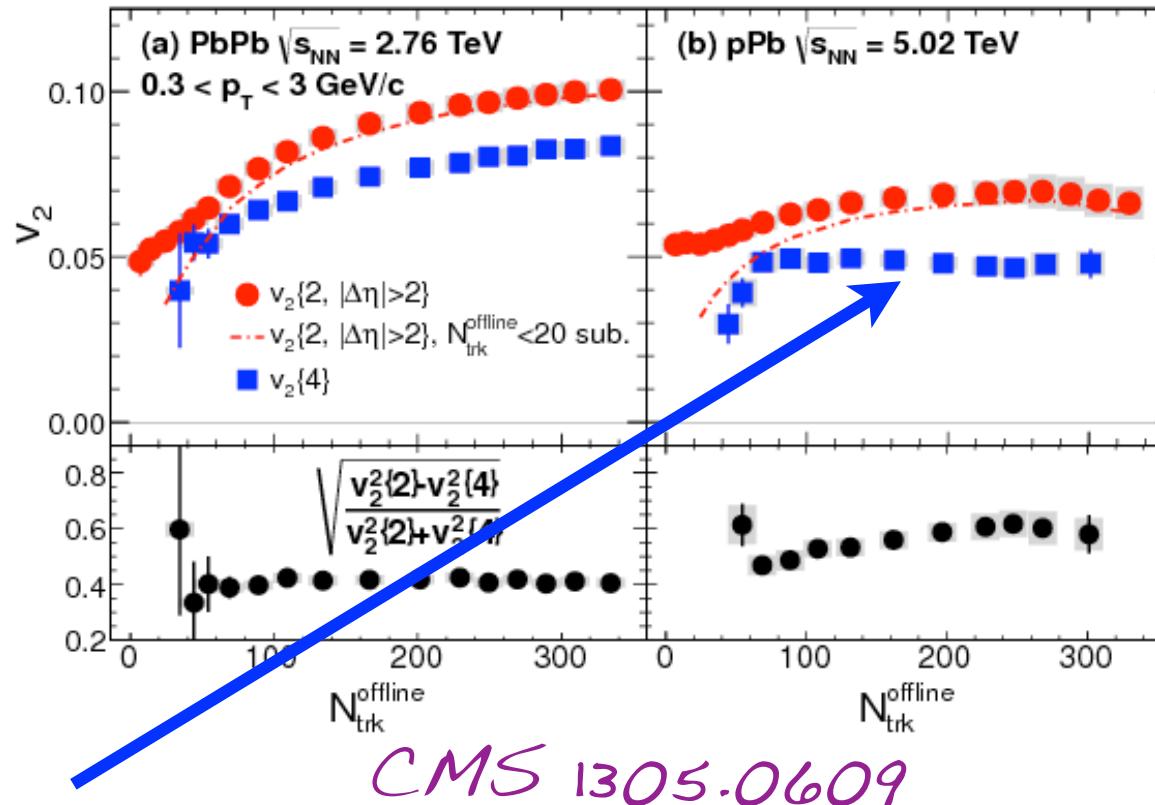
Universality of initial anisotropy fluctuations

- The *Power* distribution fits several models of initial conditions (MC Glauber, MC KLN, IP-Glasma, DIPSY) when the anisotropy is solely created by fluctuations: ε_2 in p-p collisions, ε_2 and ε_3 in p-Pb collisions, ε_3 in Pb-Pb or Au-Au collisions.

Li Yan, JYI, PRL 112 (2014) 082301

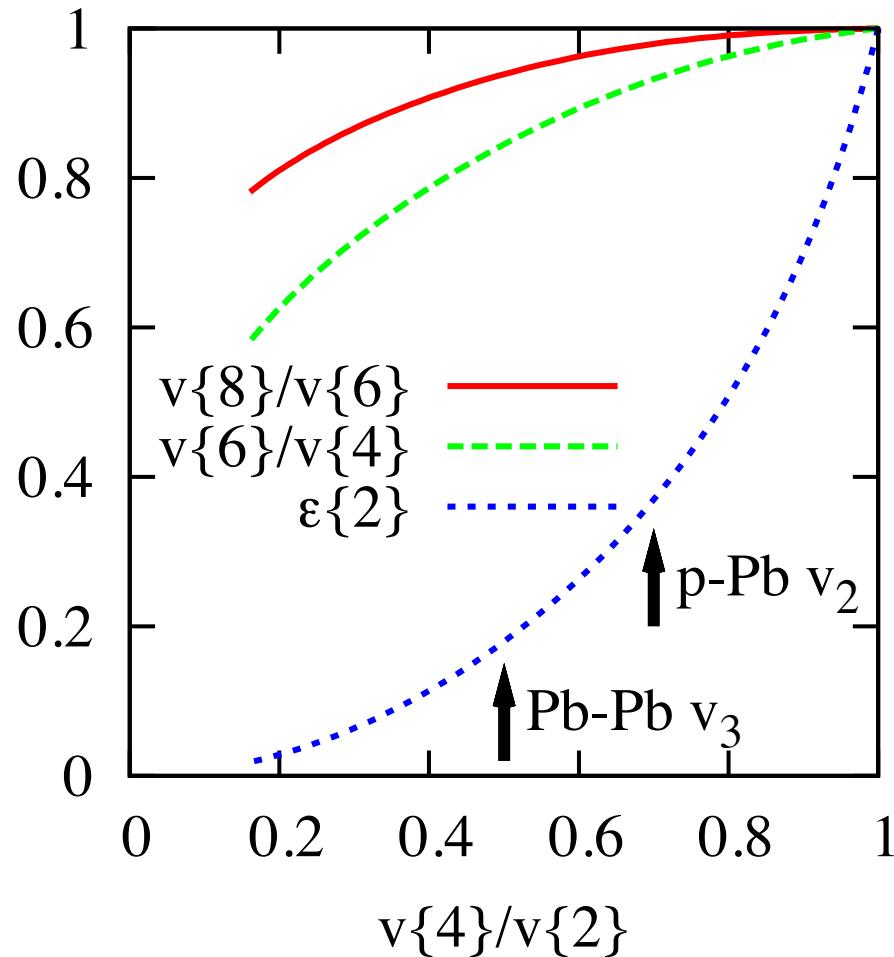
- We postulate that it is universal, to a good approximation.

Natural explanation for $v_2\{4\}$ in pPb



Gaussian fluctuations give $v_2\{4\}=0$.
Our new Power distribution naturally predicts a large $v_2\{4\}$ in p-Pb.

Consequences & predictions



- Using as input the experimentally measured ratio $v_n\{4\}/v_n\{2\}$
- Quantitative prediction for higher-order cumulants $v_n\{6\}$ and $v_n\{8\}$
- We can read off the rms anisotropy $\epsilon_n\{2\}$, a property of the initial state, directly from experimental data

Generalization to ε_2 in Pb-Pb

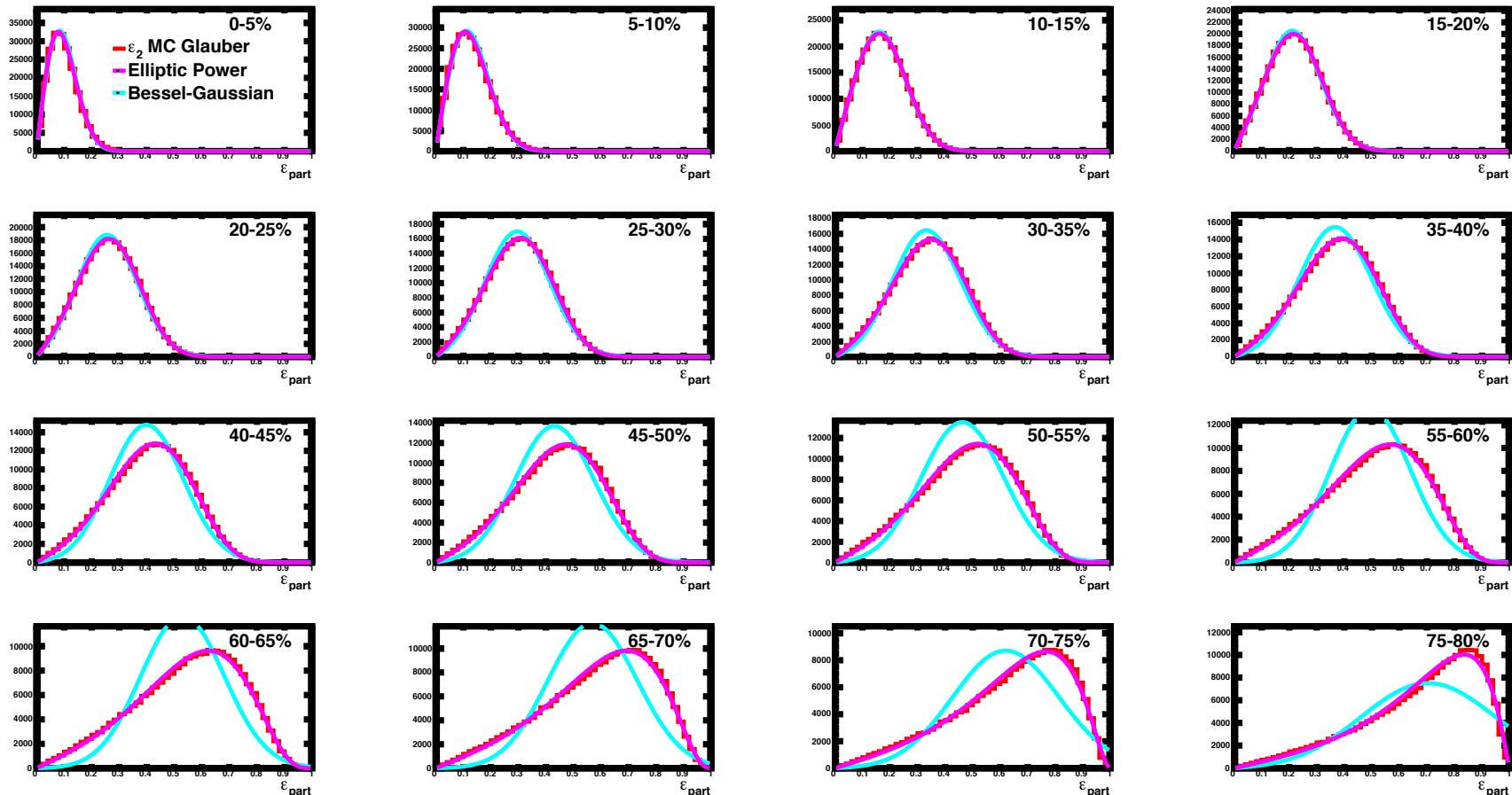
- For ε_2 in non-central Pb-Pb or Au-Au collisions, there is a **mean anisotropy in the reaction plane** in addition to **fluctuations**: requires a generalized distribution with 1 extra parameter: the *Elliptic Power* distribution

$$\frac{dn}{d\varepsilon} = \frac{2}{\pi} \varepsilon^\alpha (1 - \varepsilon^2)^{(\alpha-1)} (1 - \varepsilon_0^2)^{(\alpha+1/2)} \int_0^\pi (1 - \varepsilon_0 \varepsilon \cos \phi)^{-(1+2\alpha)} d\phi$$

Reduces to the Power distribution for $\varepsilon_0 = 0$

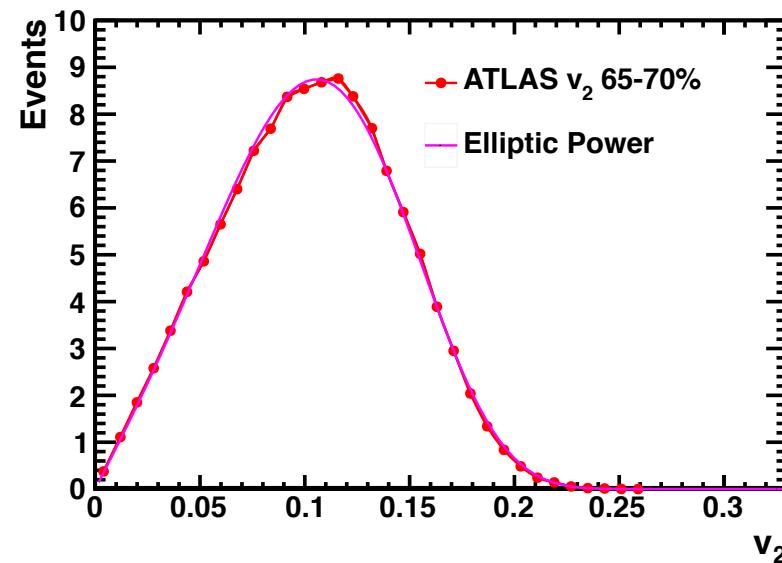
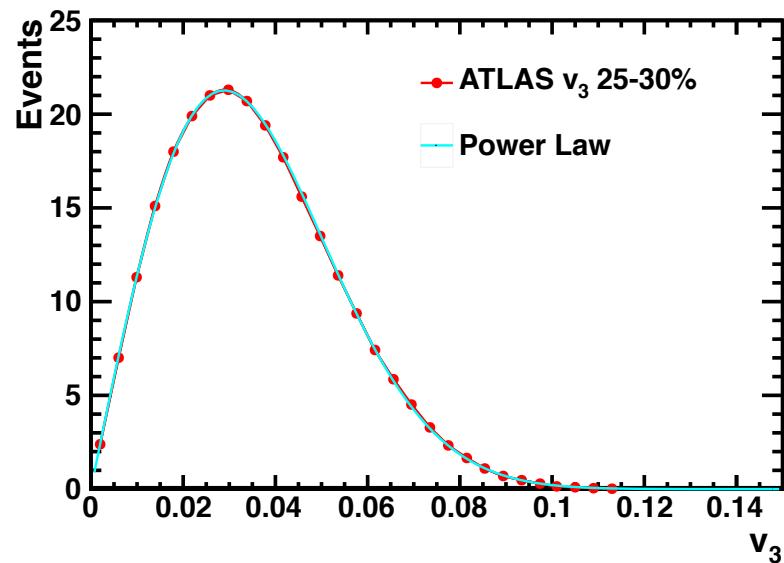
Li Yan, JYO, Art Poskanzer, in preparation

Testing the *Elliptic Power* distribution for ϵ_2



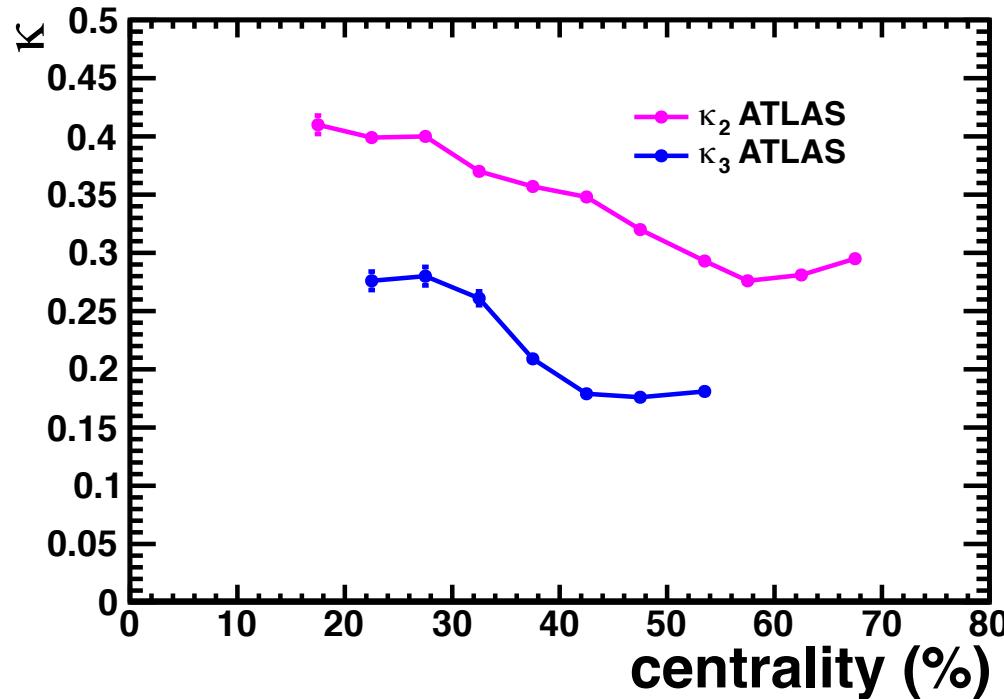
fits to Monte-Carlo Glauber by Art Poskanzer

Fitting ATLAS v_3 and v_2 distributions with rescaled *Power* and *Elliptic-Power*



We obtain good fits to ATLAS data for
 v_2 and v_3 for all centralities

Extracting the hydro response from ATLAS data

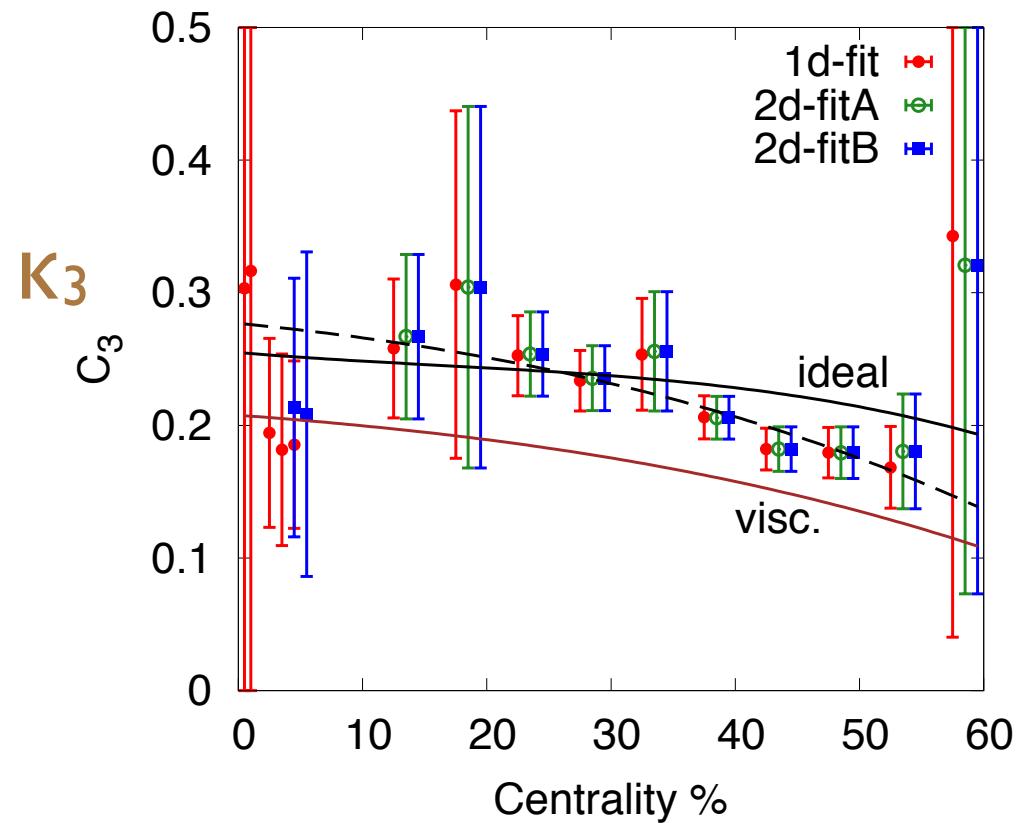


Art Poskanzer,
preliminary
(stat errors only)

As expected from viscous effects, the response decreases for more peripheral collisions.

As expected, viscous decrease is faster for v_3 than for v_2

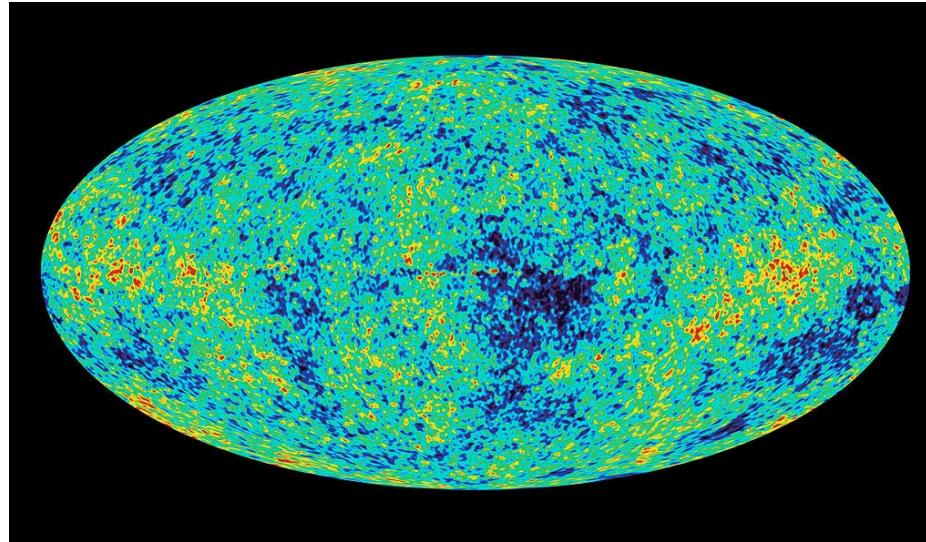
Extracting the QGP viscosity from ATLAS data



Li Yan (preliminary)
with syst errors

Viscous hydro fits return a value of η/s close to 0.1.

Big Bang versus Little Bang



WMAP

Small anisotropies observed in the cosmic microwave background are thought to originate from quantum fluctuations in the early Universe.

Anisotropic flow at RHIC and LHC is a similar phenomenon, occurring within a tiny system with large fluctuations.

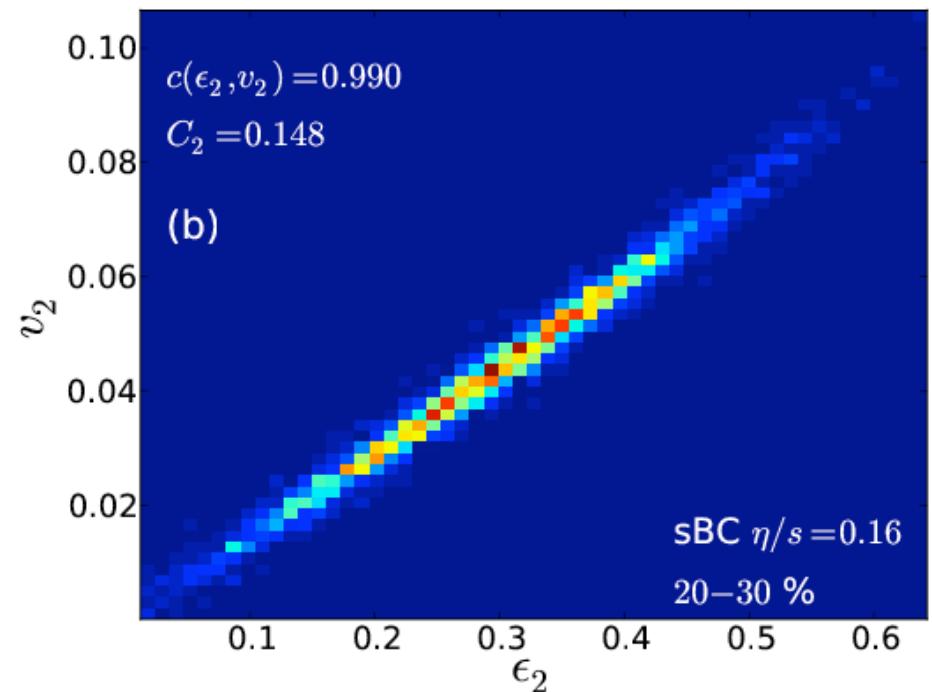
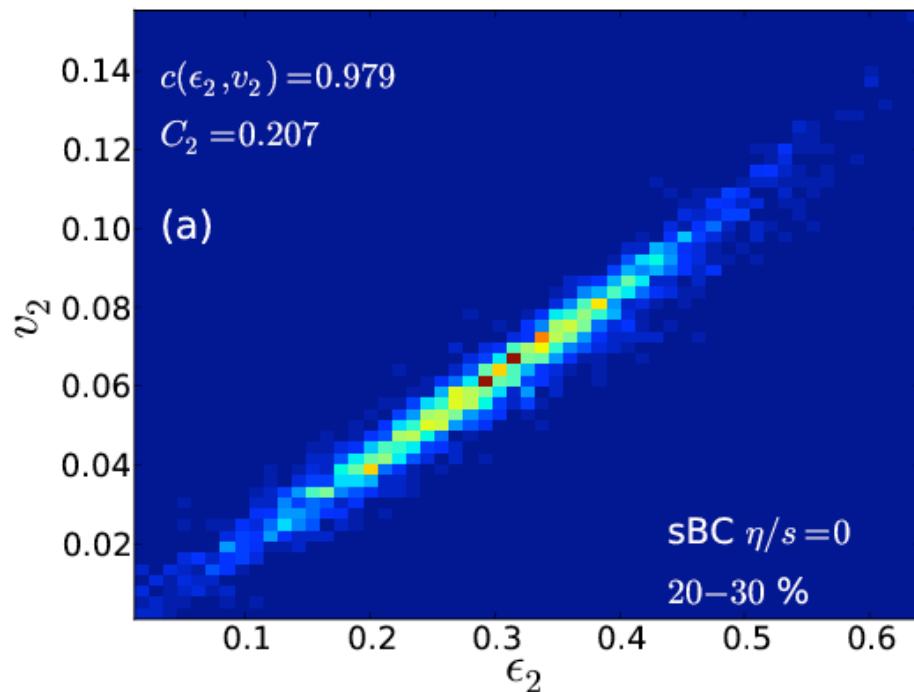
The non-Gaussianity of these fluctuations, and the fact that they are universal, allows us to disentangle initial fluctuations from the response.

Conclusions, perspectives

- Direct evidence from experimental data that anisotropic flow in p-Pb and Pb-Pb collisions is driven by **large anisotropies** in the initial state: the statistics of ε_n hits the boundary $\varepsilon_n < 1$
- The statistics of large fluctuations is not described by the central limit theorem but nevertheless **universal** to a good approximation
- We can extract both the initial anisotropy and the “hydrodynamic” response K_n from experimental data without any prior assumption about the initial state. Toward the first robust measurement of the viscosity of the QGP (work in progress).

Backup

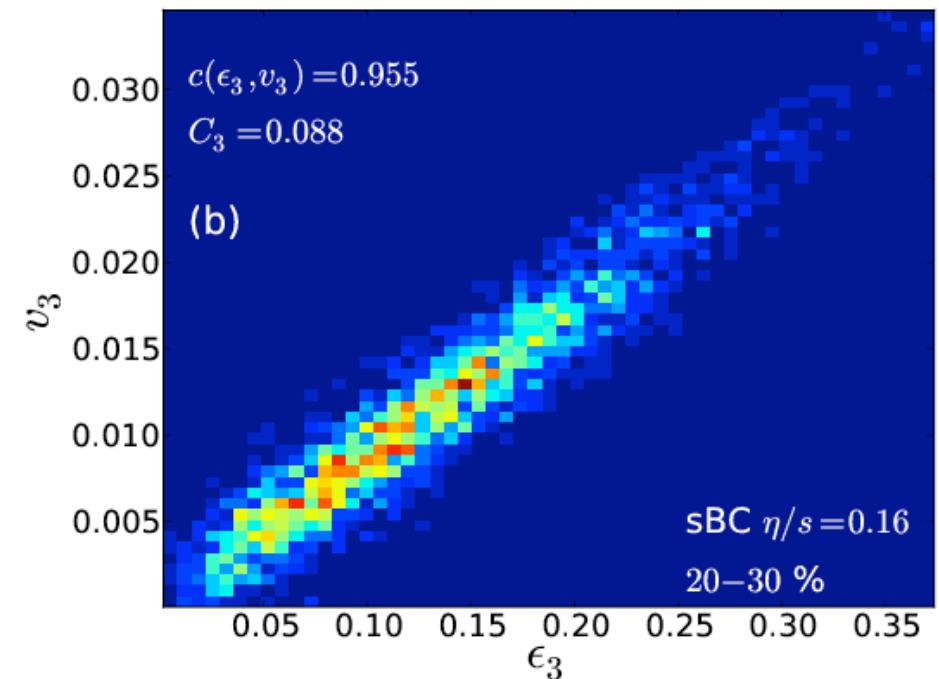
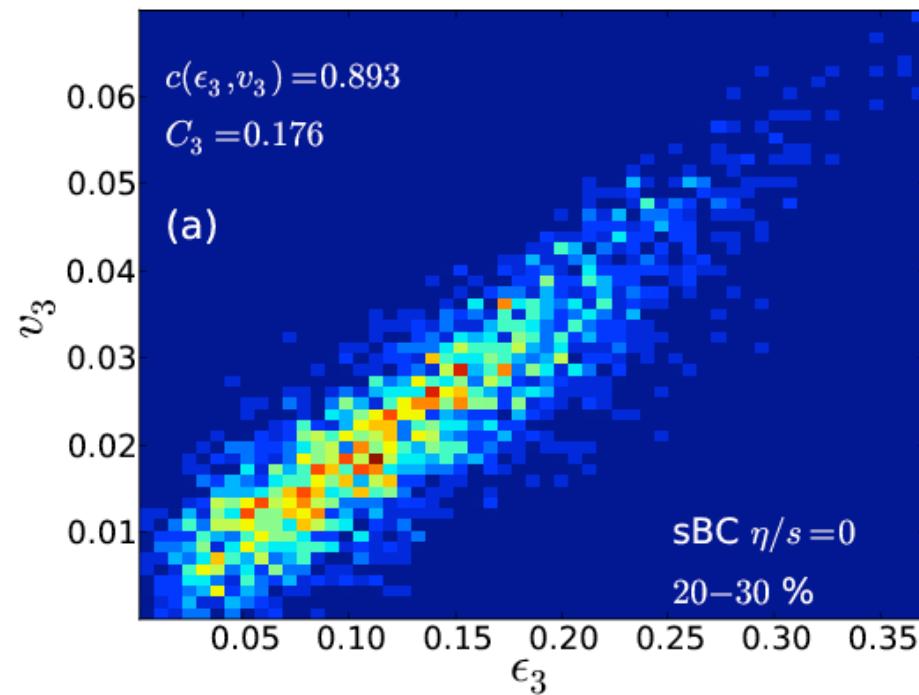
Elliptic flow v_2 versus initial eccentricity ϵ_2



Niemi Denicol Holopainen Huovinen 1212.1008

Each point=different initial density profile.
 v_2 is almost perfectly linear in ϵ_2

Triangular flow v_3 versus initial triangularity ϵ_3



Niemi Denicol Holopainen Huovinen 1212.1008

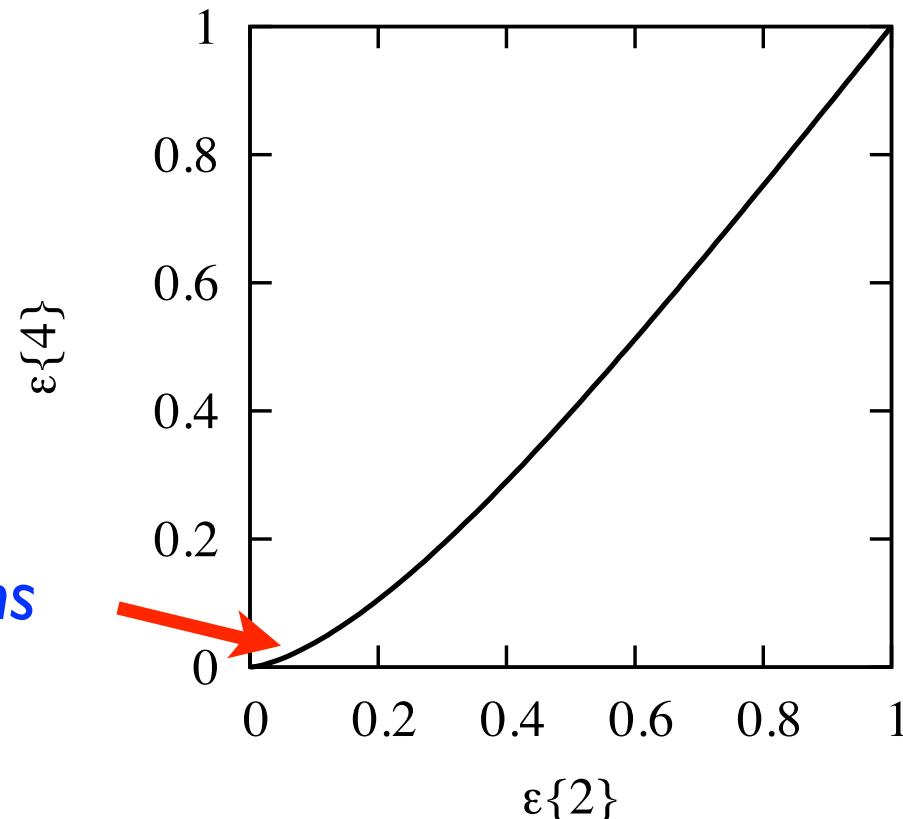
v_3 is also strongly correlated with ϵ_3

Cumulants

- 2-dimensional Gaussian: Wick's theorem
 $\langle \varepsilon^4 \rangle = 2 \langle \varepsilon^2 \rangle^2$ where $\langle \dots \rangle \equiv$ average over events
- Define $\varepsilon\{2\} \equiv \langle \varepsilon^2 \rangle^{1/2}$ (rms anisotropy)
$$\varepsilon\{4\} = (2\langle \varepsilon^2 \rangle^2 - \langle \varepsilon^4 \rangle)^{1/4}$$
- $\varepsilon\{4\} = 0$ for Gaussian.
- The power distribution predicts a universal, relation between $\varepsilon\{4\}$ and $\varepsilon\{2\}$

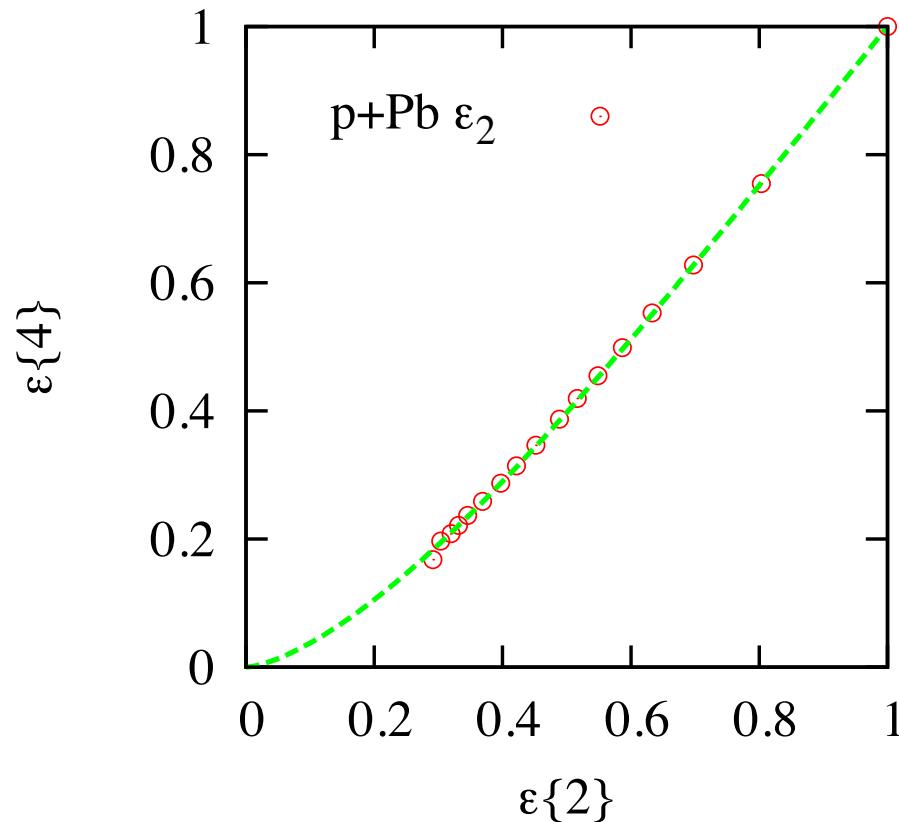
Testing universality with cumulants

Central limit:
large system,
small fluctuations
 $\varepsilon\{2\} \ll 1$ and
 $\varepsilon\{4\} \ll \varepsilon\{2\}$



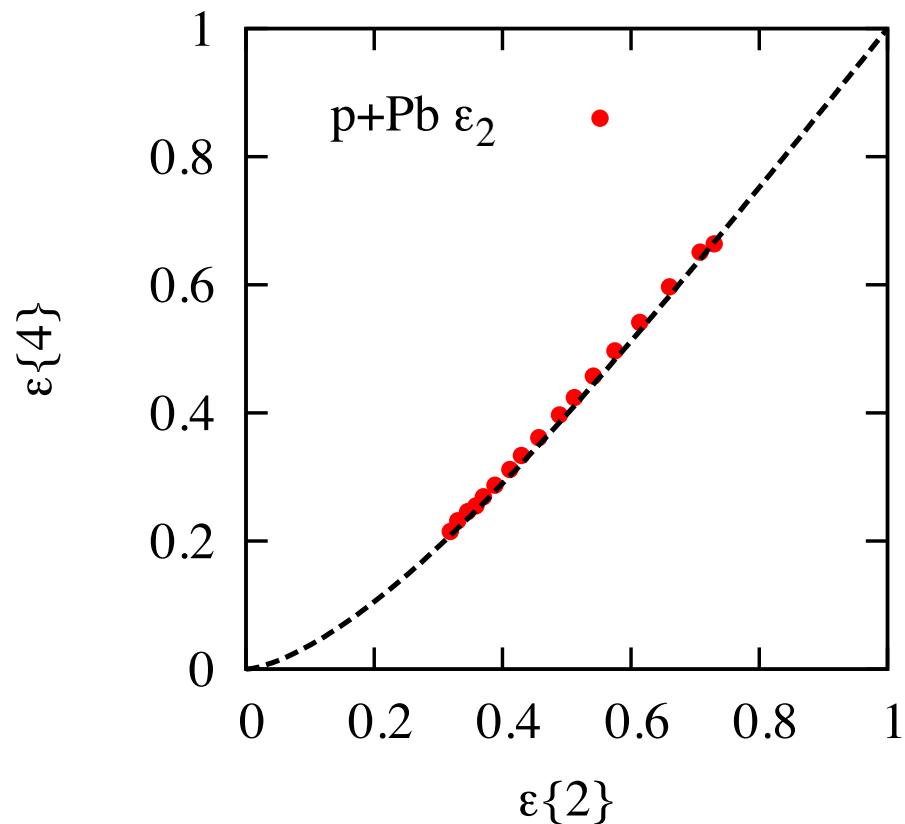
Prediction of the power distribution

Testing universality with cumulants



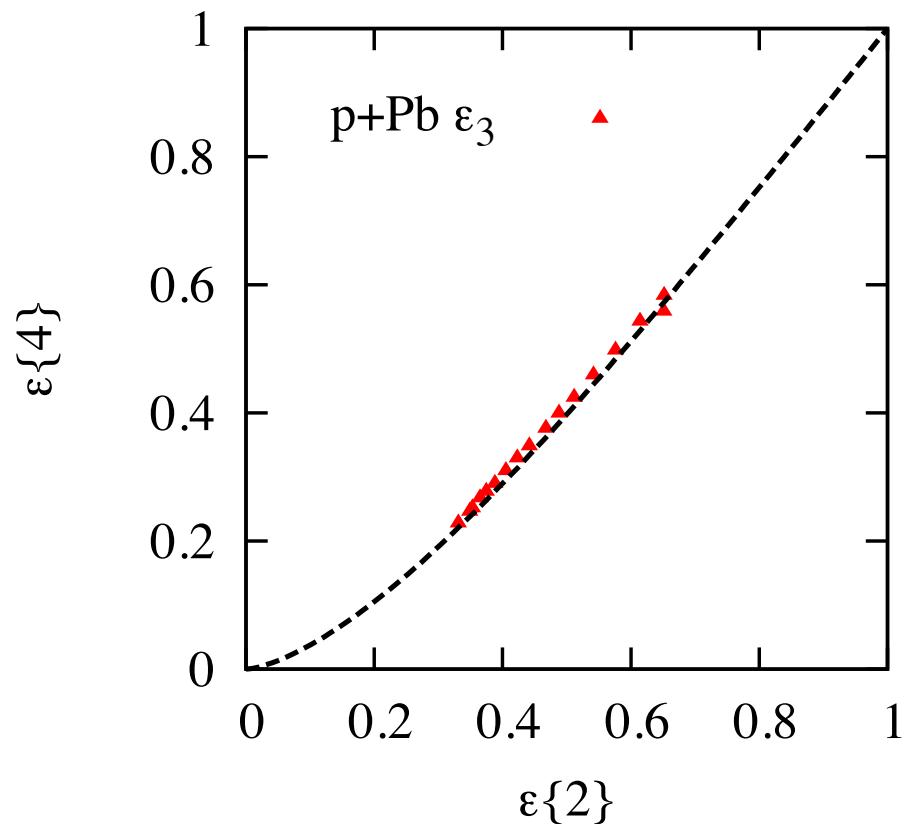
Pointlike sources with Gaussian distribution:
power distribution=exact=test of Monte-Carlo

Testing universality with cumulants



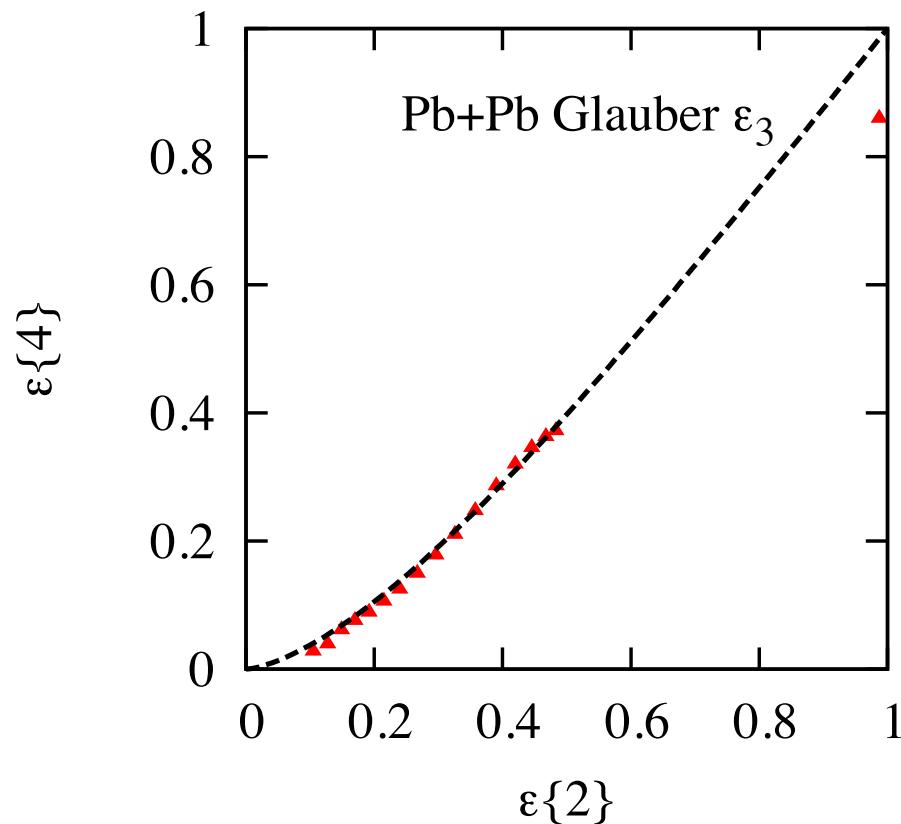
Each point: different number of hit nucleons in target

Testing universality with cumulants



Each point: different number of hit nucleons in target

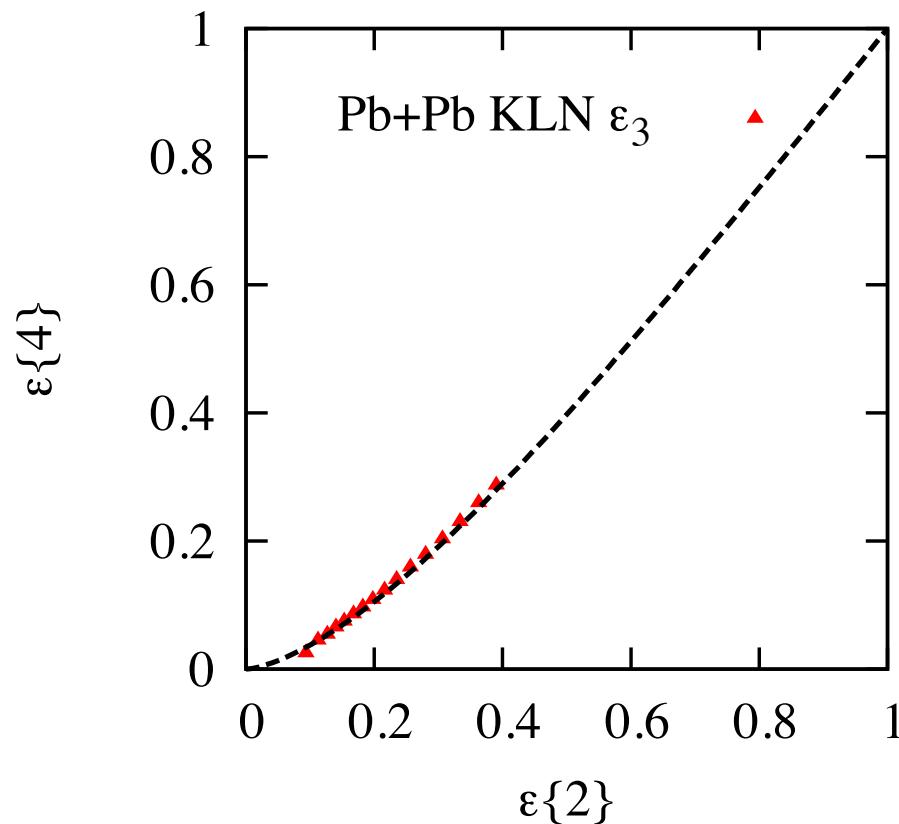
Testing universality with cumulants



Each point: different centrality

Pb-Pb: Larger system: smaller anisotropies

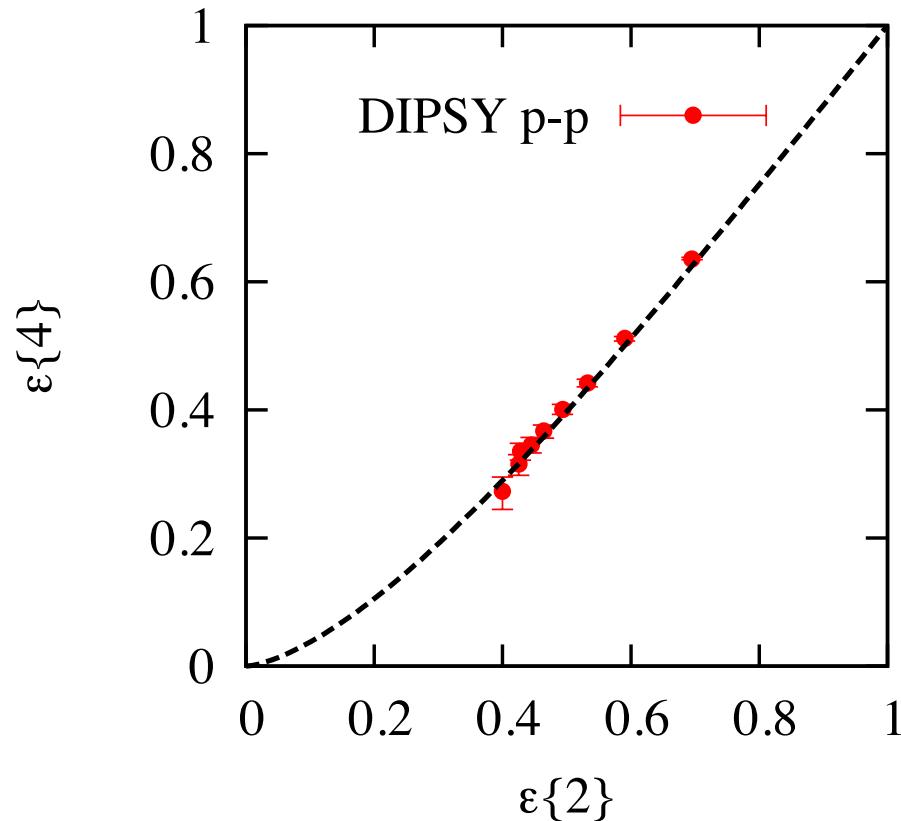
Testing universality with cumulants



Each point: different centrality

Pb-Pb: Larger system: smaller anisotropies

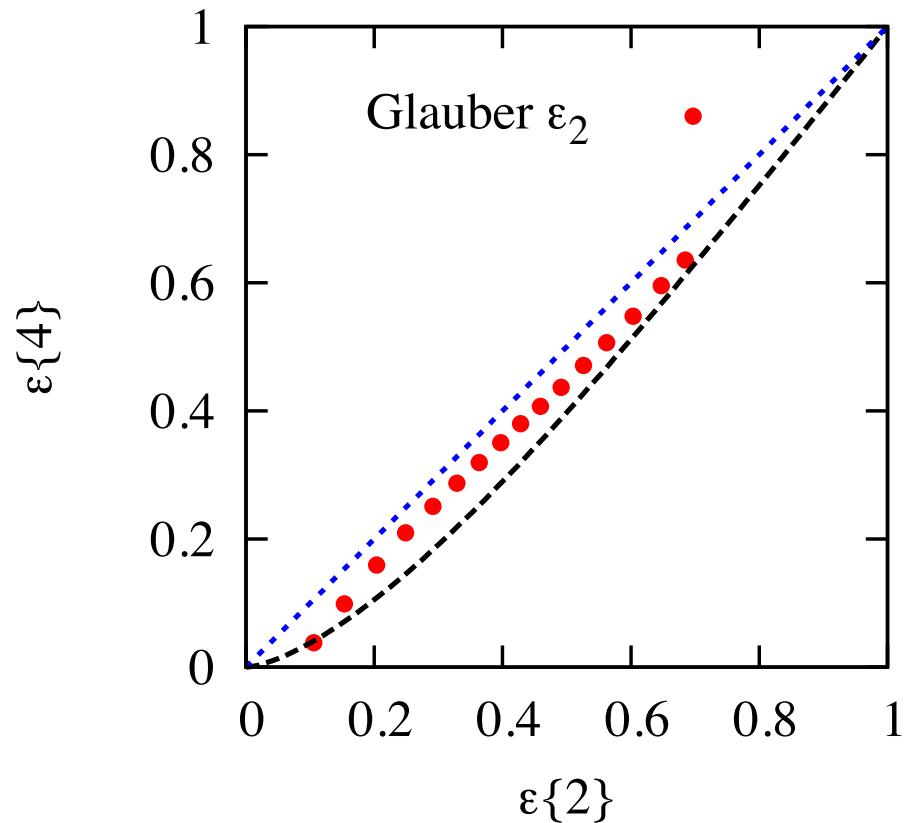
Testing universality with cumulants



data from Avsar Flensburg Hatta JYO Ueda 1009.5643

Each point: different parton multiplicity

Elliptic anisotropy in Pb-Pb



Driven by almond shape of overlap area, not fluctuations:
Deviates from the power distribution

Applying the power distribution to experimental data

If $v_n = K_n \varepsilon_n$, with constant K_n

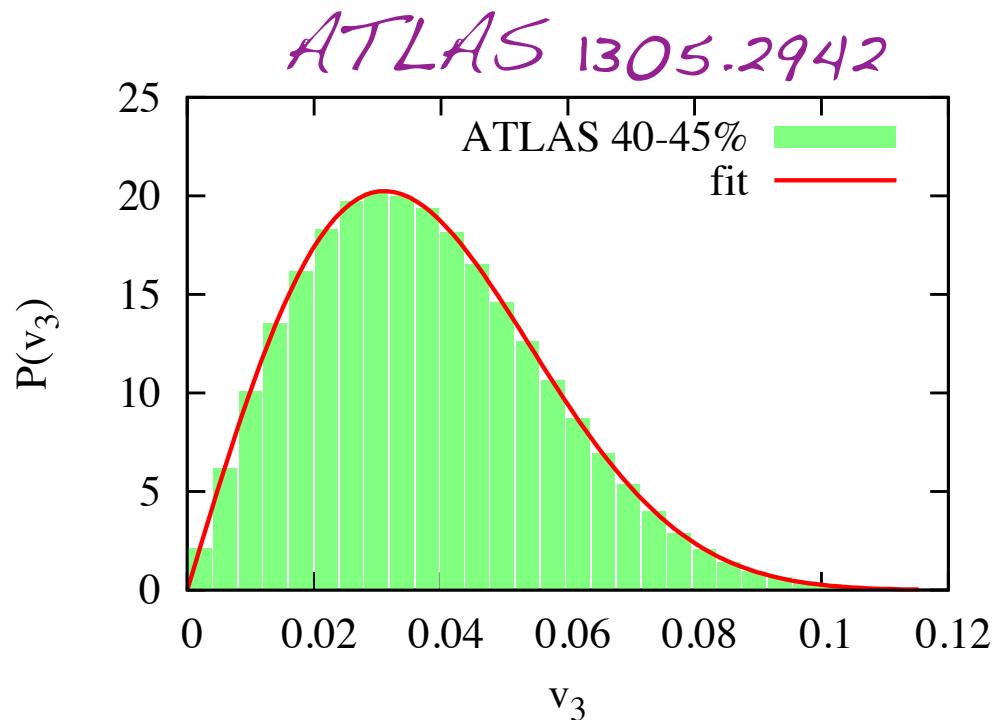
then $v_n\{4\}/v_n\{2\} = \varepsilon_n\{4\}/\varepsilon_n\{2\}$

we can read off the parameter α from the experimentally-measured ratio $v_n\{4\}/v_n\{2\}$

Fitting the distribution of v_n

The ATLAS distribution has published the distribution of v_n with $n=2,3,4$ in Pb-Pb collisions. We can fit these data assuming $v_3 \approx K_3 \varepsilon_3$ and a power distribution for ε_3

*The fit returns
 $K_3 = 0.18 \pm 0.02$,
in agreement with
viscous hydrodynamics*



Yan, Poskanzer, JYO, in preparation

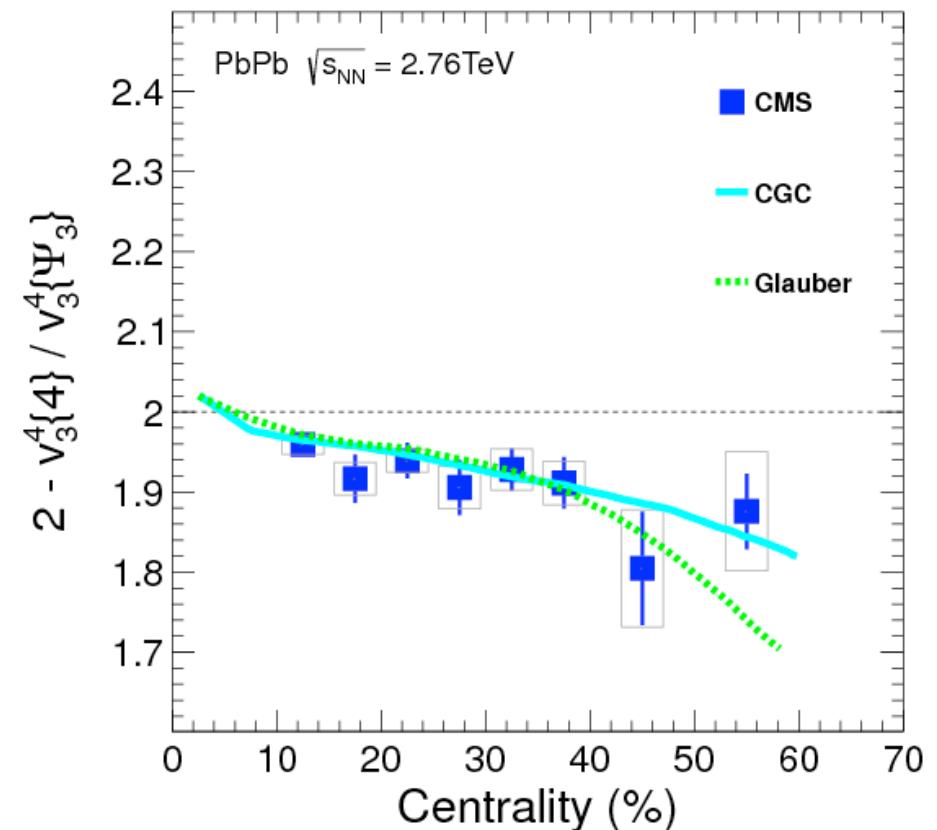
Simple predictions from eccentricity scaling

- Experimentally, one can measure moments (or cumulants) of the distribution of v_n .
- Eccentricity scaling implies that, e.g.
$$\langle v_n^4 \rangle / \langle v_n^2 \rangle^2 = \langle \varepsilon_n^4 \rangle / \langle \varepsilon_n^2 \rangle^2$$
- Thus one can check if a particular model of the initial state is compatible with data.

Eccentricity scaling versus data

data: $\langle v_3^4 \rangle / \langle v_3^2 \rangle^2$

models: $\langle \varepsilon_3^4 \rangle / \langle \varepsilon_3^2 \rangle^2$



CMS 1310.8651

Higher-order cumulants

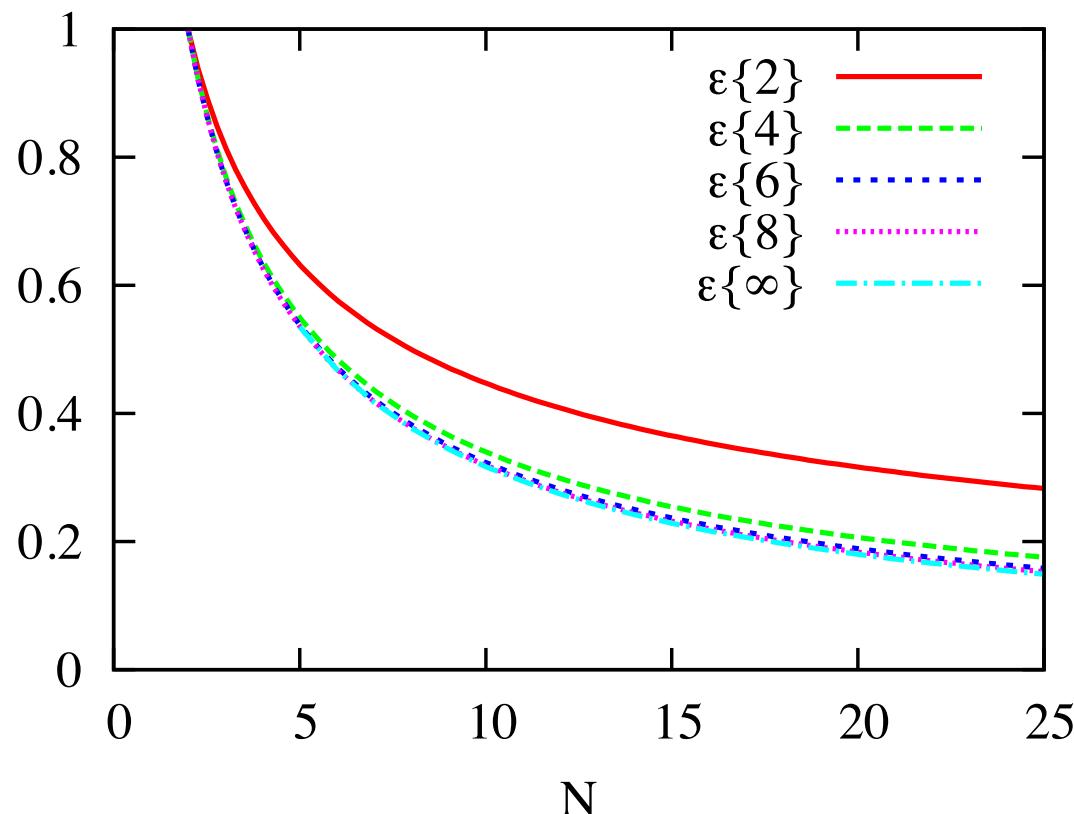
Expand the generating function

$$G(k) = \ln \langle \exp(ik \cdot \varepsilon) \rangle$$

where k and ε are 2-d vectors in the transverse plane, to order k^{2n} .

Asymptotic behavior for large n = singularity of $G(k)$ = zero of the Fourier transform of the distribution of ε .

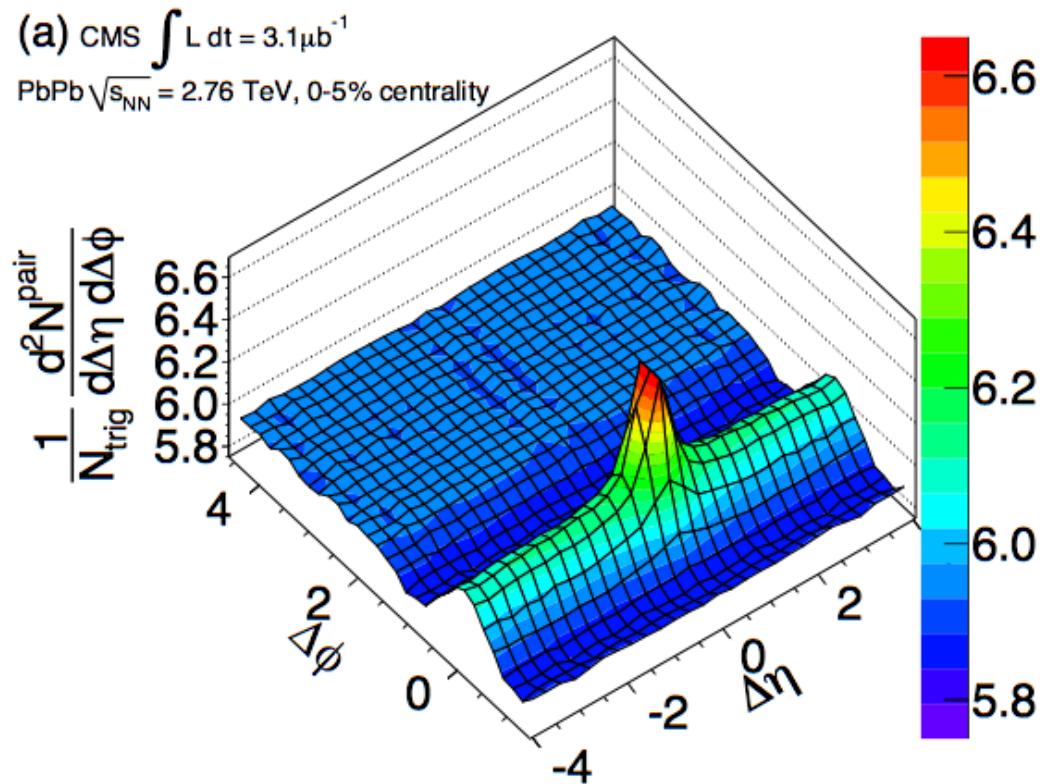
Higher-order cumulants (predicted by the power distribution)



$\varepsilon\{n\}$ quickly converges as order n increases

Anisotropic flow

⇒ correlations at large $\Delta\eta$

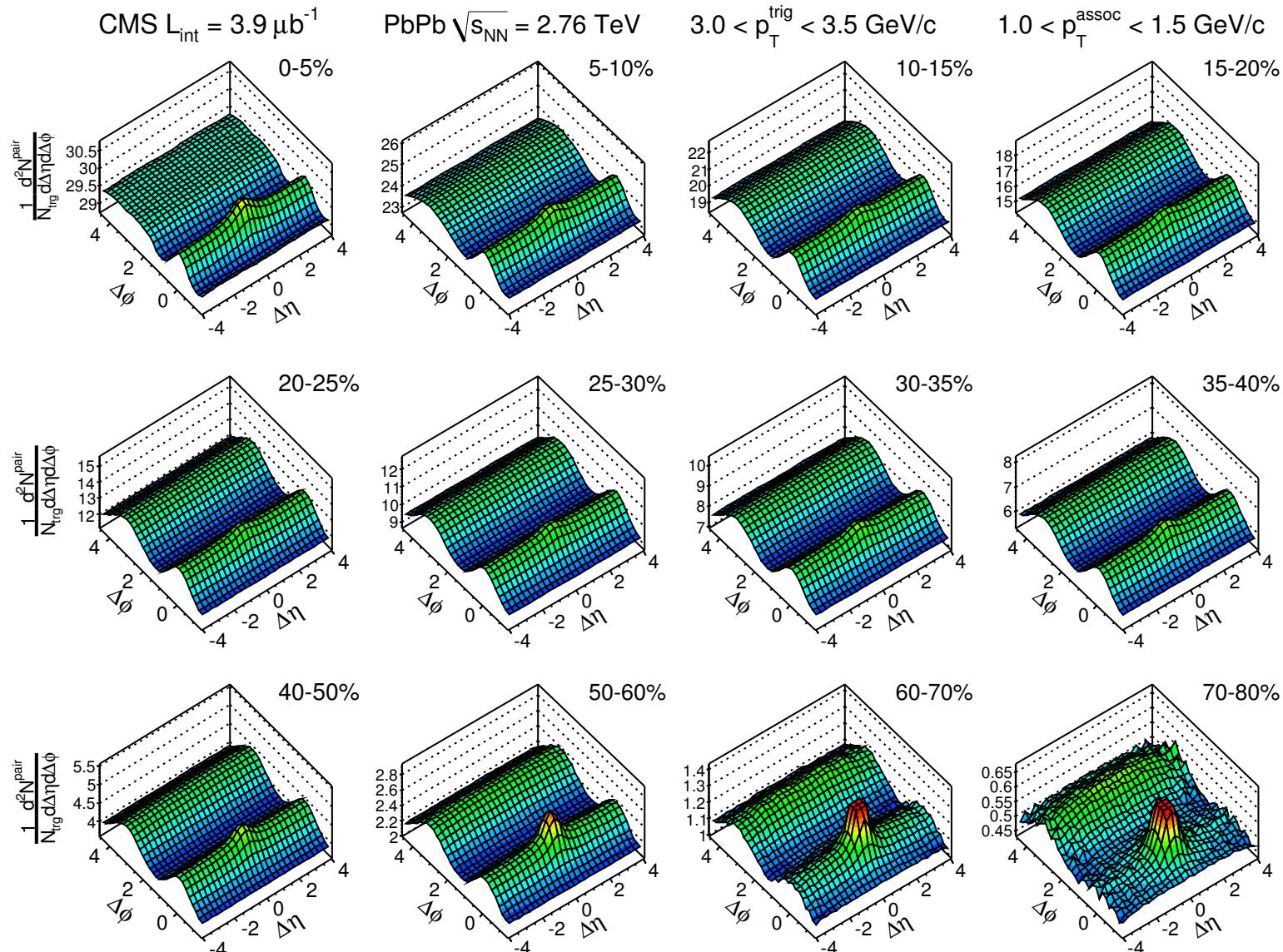


Number of pairs of particles versus relative azimuthal angle and pseudorapidity (~polar angle) in **central Pb-Pb collisions**

CMS arXiv:1105.2438

Anisotropic flow

 central



 peripheral

CMS 1201.3158