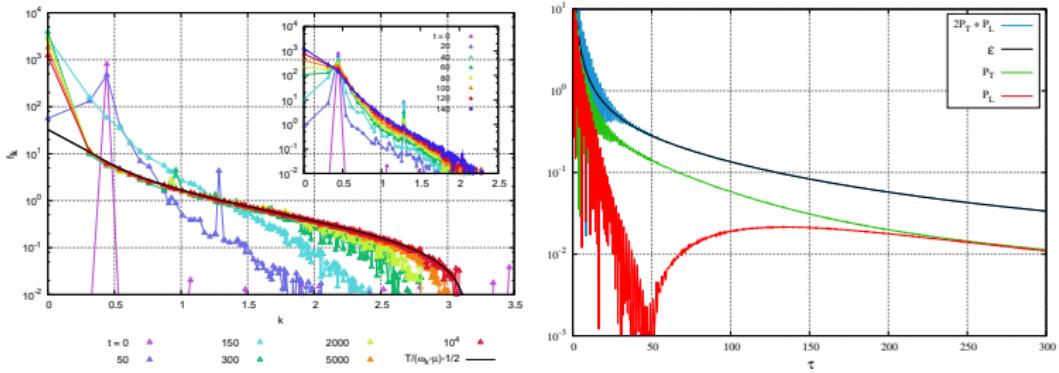


A First Step Towards Understanding Izotropization in HIC



Orsay, 26 October 2012

Thomas EPELBAUM
IPhT

① THEORETICAL FRAMEWORK

② FIXED VOLUME

③ EXPANDING VOLUME

④ CONCLUSION

① THEORETICAL FRAMEWORK

CGC and JIMWLK
Resummation Scheme

② FIXED VOLUME

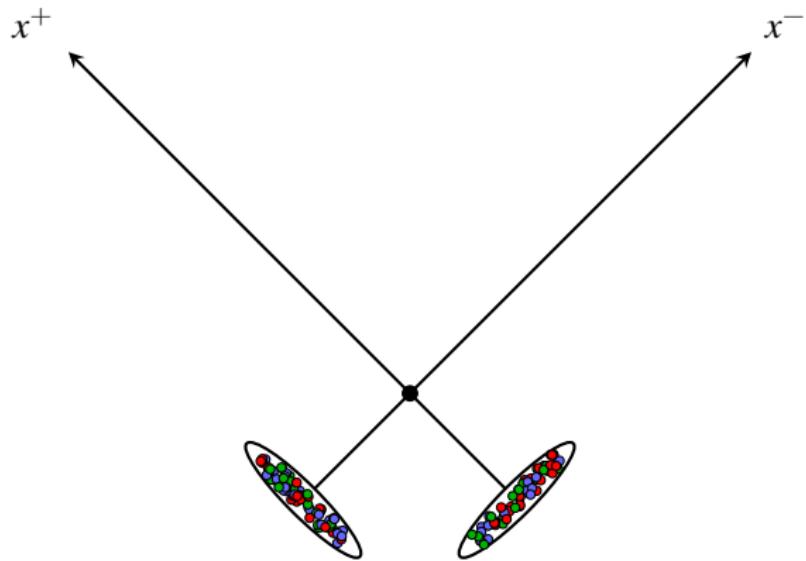
The Model
Energy-Momentum Tensor
Distribution Function
Bose-Einstein Condensation

③ EXPANDING VOLUME

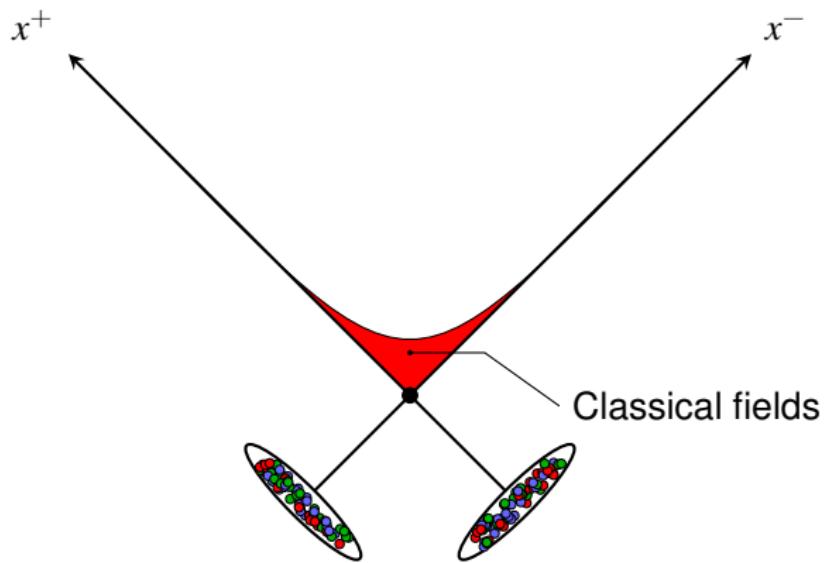
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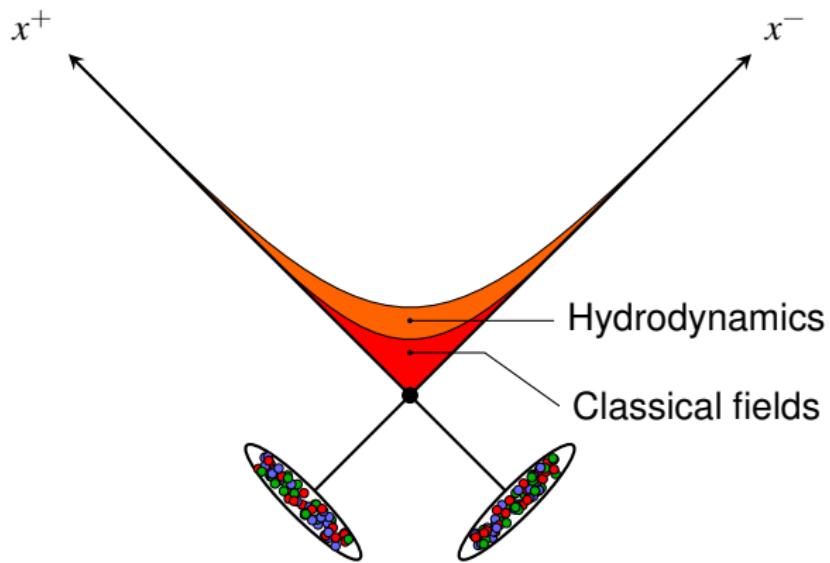
THE GENERAL PICTURE



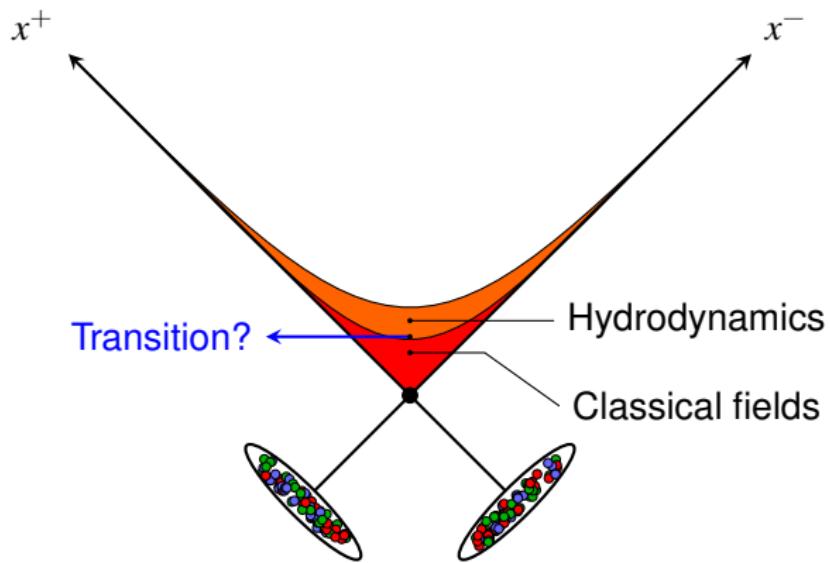
THE GENERAL PICTURE



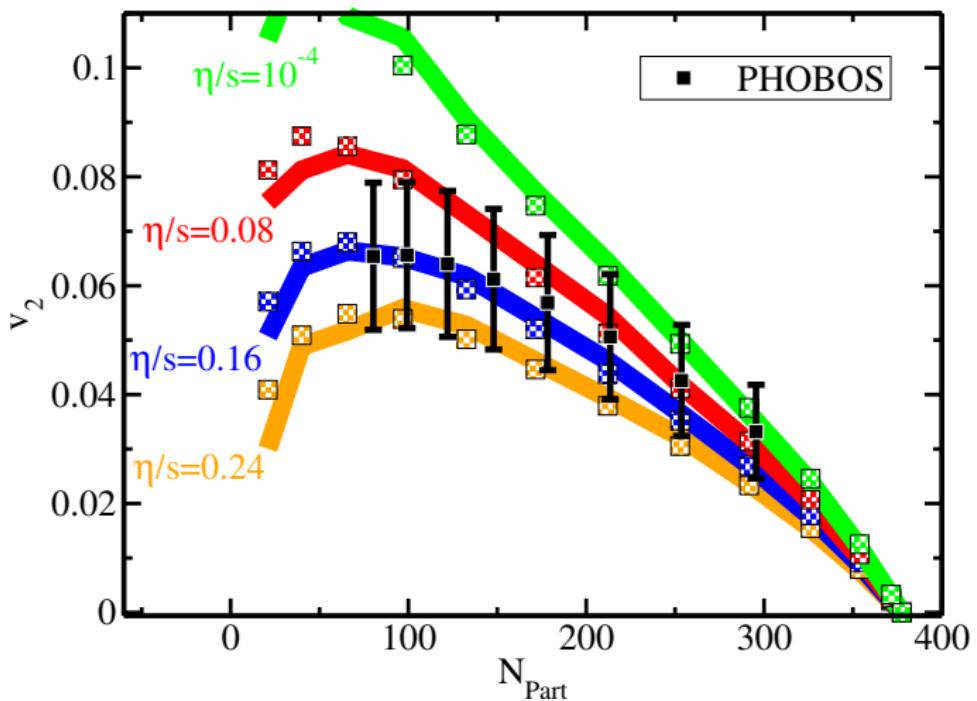
THE GENERAL PICTURE



THE GENERAL PICTURE

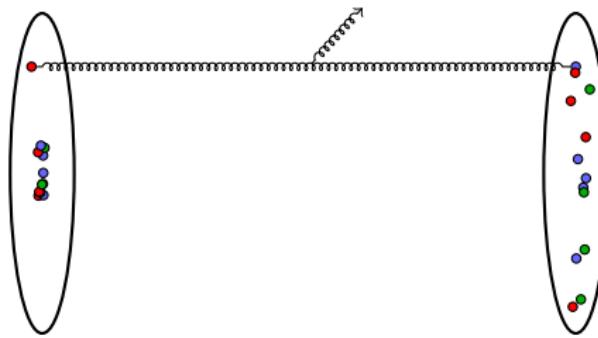


CGC

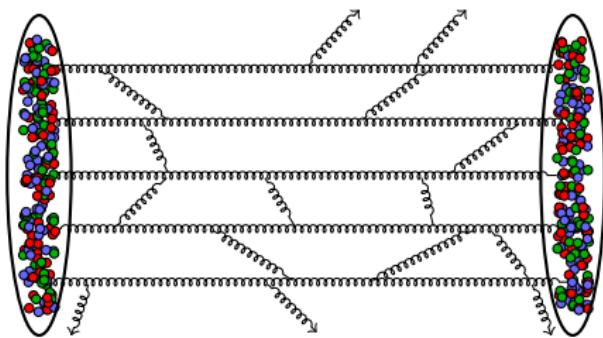


M. Luzum, P. Romatschke, Phys. Rev. C 78, (2008).

Dilute Regime



Dense regime



HYDRO PREREQUISITES

- Equation of state (EOS), small anisotropy
- Initial conditions: energy density, pressure
- Transport coefficients: viscosity...

THEORETICAL FRAMEWORK

- Color Glass Condensate (CGC)
- JIMWLK Equation
- Resummation Scheme

① THEORETICAL FRAMEWORK

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[MCLERRAN, VENUGOPALAN]

[JALILIAN-MARIAN, IANCU, MCLERRAN, WEIGERT, LEONIDOV, KOVNER]

Schematical vision of the degree of freedom

CUTOFF DEPENDANCETheory Λ dependant $\rightarrow \log(\Lambda)$ appears at NLO

[MCLERRAN, VENUGOPALAN]

[JALILIAN-MARIAN, IANCU, MCLERRAN, WEIGERT, LEONIDOV, KOVNER]

Schematical vision of the degree of freedom

CUTOFF DEPENDANCE



Theory Λ dependant $\rightarrow \log(\Lambda)$ appears at NLO

RENORMALIZATION GROUP EQUATION

Absorb the Λ dependance in the source

[MCLERRAN, VENUGOPALAN]

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Schematical vision of the degree of freedom

CUTOFF DEPENDANCE



Theory Λ dependant $\rightarrow \log(\Lambda)$ appears at NLO

RENORMALIZATION GROUP EQUATION

Absorb the Λ dependance in the source

Issues:

- Very Anisotropic system at $\tau = 0^+$
- Secular divergences.

ANISOTROPY

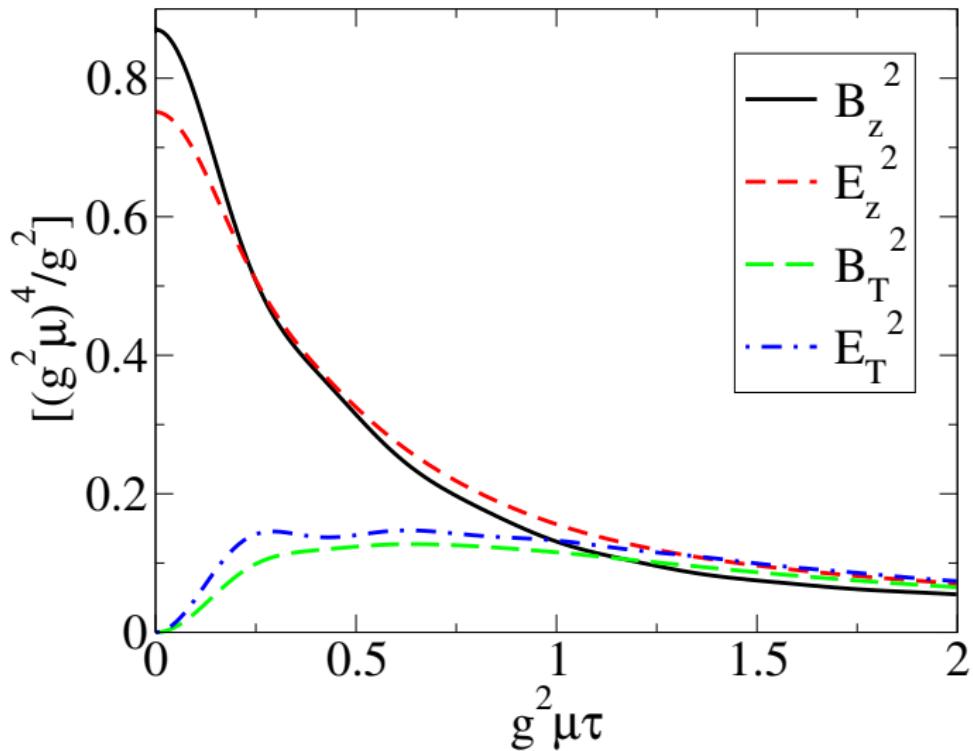


$$\epsilon = E_T^2 + B_T^2 + E_L^2 + B_L^2$$

$$P_T = E_L^2 + B_L^2$$

$$P_L = E_T^2 + B_T^2 - E_L^2 - B_L^2$$

ANISOTROPY



T. Lappi, L. McLerran Nucl. Phys. A772, (2006).

ANISOTROPY



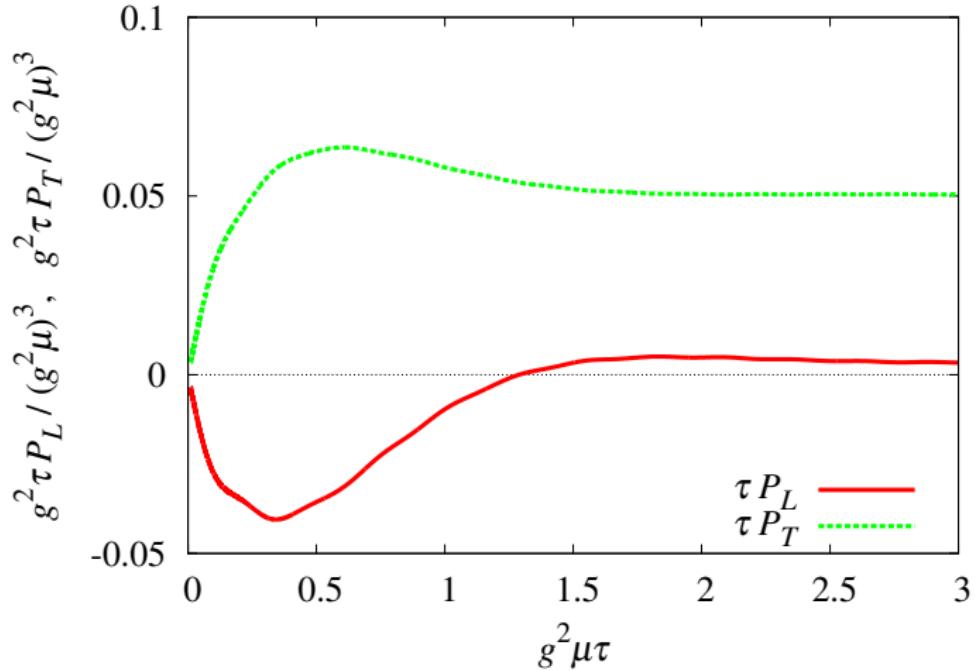
$$\epsilon = \underbrace{E_T^2}_0 + \underbrace{B_T^2}_0 + E_L^2 + B_L^2$$

$$P_T = E_L^2 + B_L^2$$

$$P_L = \underbrace{E_T^2}_0 + \underbrace{B_T^2}_0 - E_L^2 - B_L^2$$

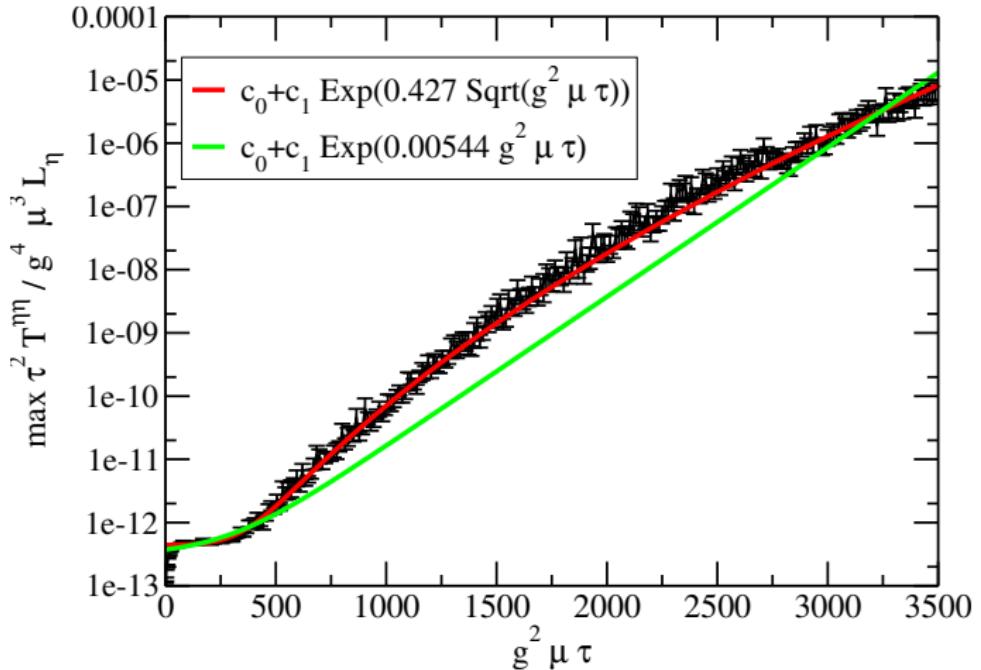
Initial $T^{\mu\nu}$ is $(\epsilon, \epsilon, \epsilon, -\epsilon)$!

ANISOTROPY



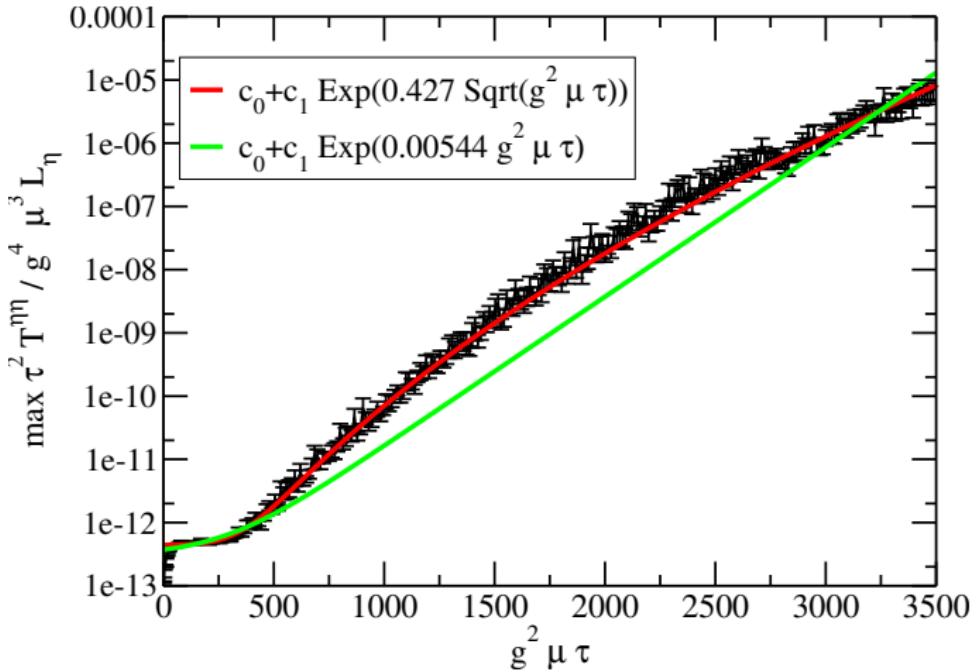
K. Fukushima, F. Gelis Nucl. Phys. **A874**, (2012).

[MROWCZYNSKI]



P. Romatschke, R. Venugopalan Phys. Rev. Lett. **96**, (2006).
 Fluctuations grows like $e^{\sqrt{\mu\tau}}$!

[MROWCZYNSKI]



P. Romatschke, R. Venugopalan Phys. Rev. Lett. **96**, (2006).
 when $g e^{\sqrt{\mu\tau}} \approx 1$, perturbative expansion breaks down!

① THEORETICAL FRAMEWORK

CGC and JIMWLK

Resummation Scheme

RESUMMATION

$$T_{\text{resum}}^{\mu\nu} = e^{\hat{O}} T_{\text{LO}}^{\mu\nu}[\varphi_0] = T_{\text{LO}}^{\mu\nu}[\varphi_0] + T_{\text{NLO}}^{\mu\nu}[\varphi_0] + \dots$$

New perturbative development parameter $g e^{\sqrt{\mu\tau}}$

INITIAL ϕ

$$\phi = \varphi_0 + \int_k c_k a_k + c_k^* a_{-k}$$

with c_k random gaussian numbers.

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LAGRANGIAN OF THE THEORY

$$\mathcal{L}(\phi) = \frac{1}{2}(\partial_\mu \phi)(\partial^\mu \phi) - \underbrace{\frac{g^2}{4!}\phi^4}_{V(\phi)}$$

WHY DO WE USE THIS MODEL?

- Scale invariance in $3 + 1$ dimensions
- Parametric resonance
- A lot simpler!

② FIXED VOLUME

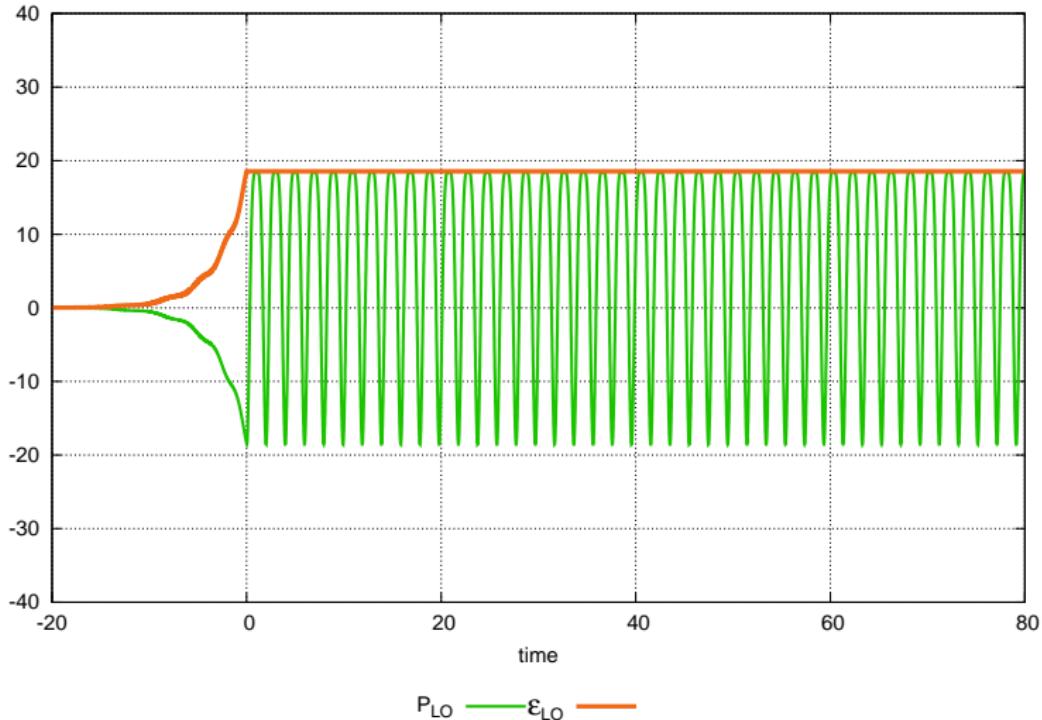
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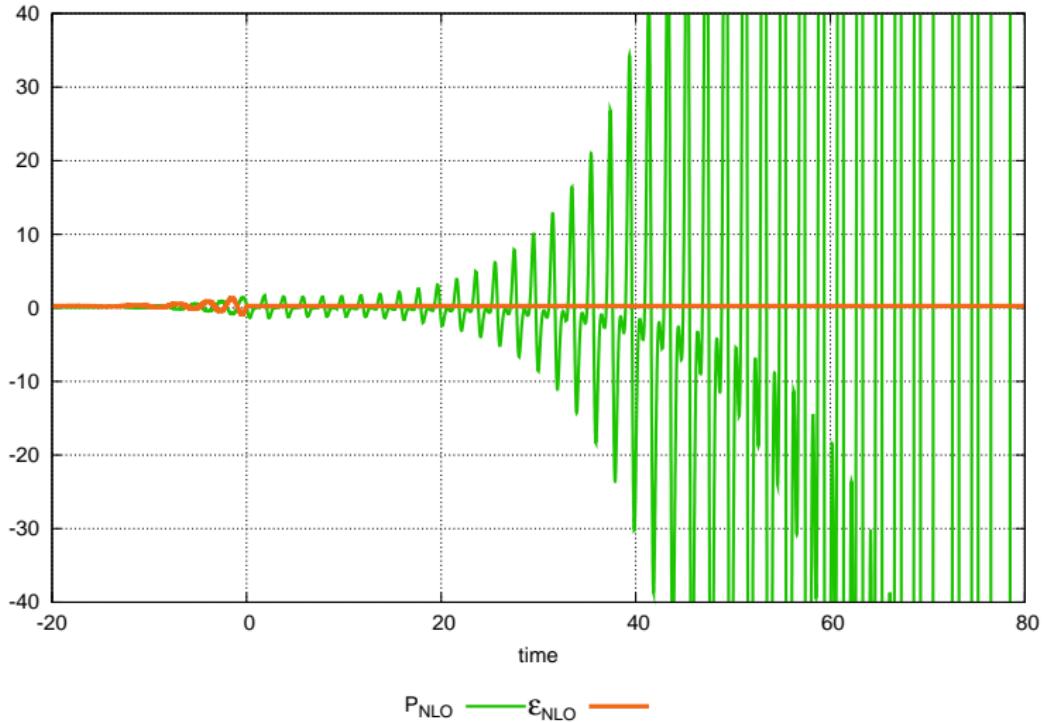
$$T_{\text{LO}}^{\mu\nu}$$



SECULAR DIVERGENCES

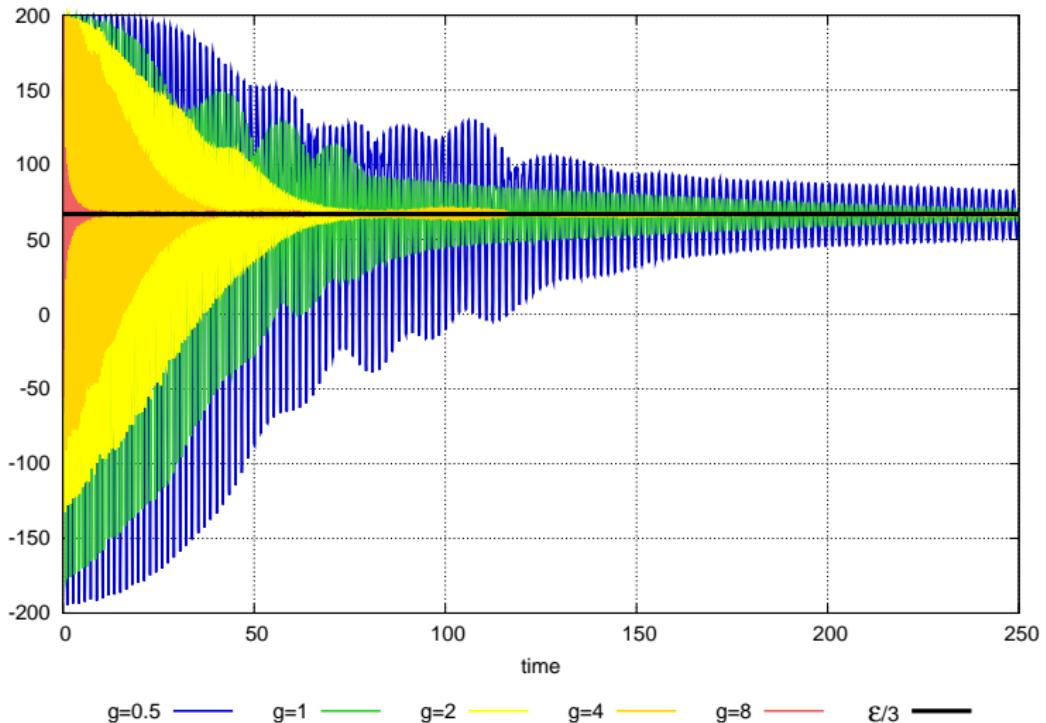


$T_{\text{NLO}}^{\mu\nu}$



PRESSURE EQUILIBRATION

$$T_{\text{resum}}^{\mu\nu}$$



g=0.5 ————— g=1 ————— g=2 ————— g=4 ————— g=8 ————— $\varepsilon/3$ —————

② FIXED VOLUME

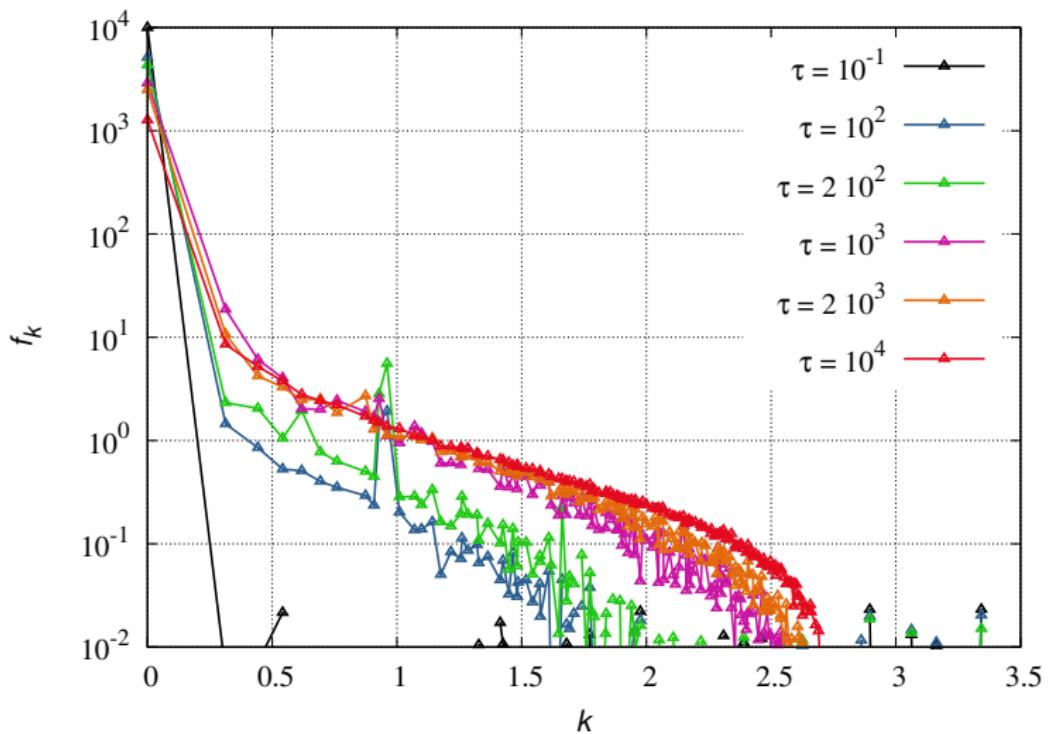
The Model

Energy-Momentum Tensor

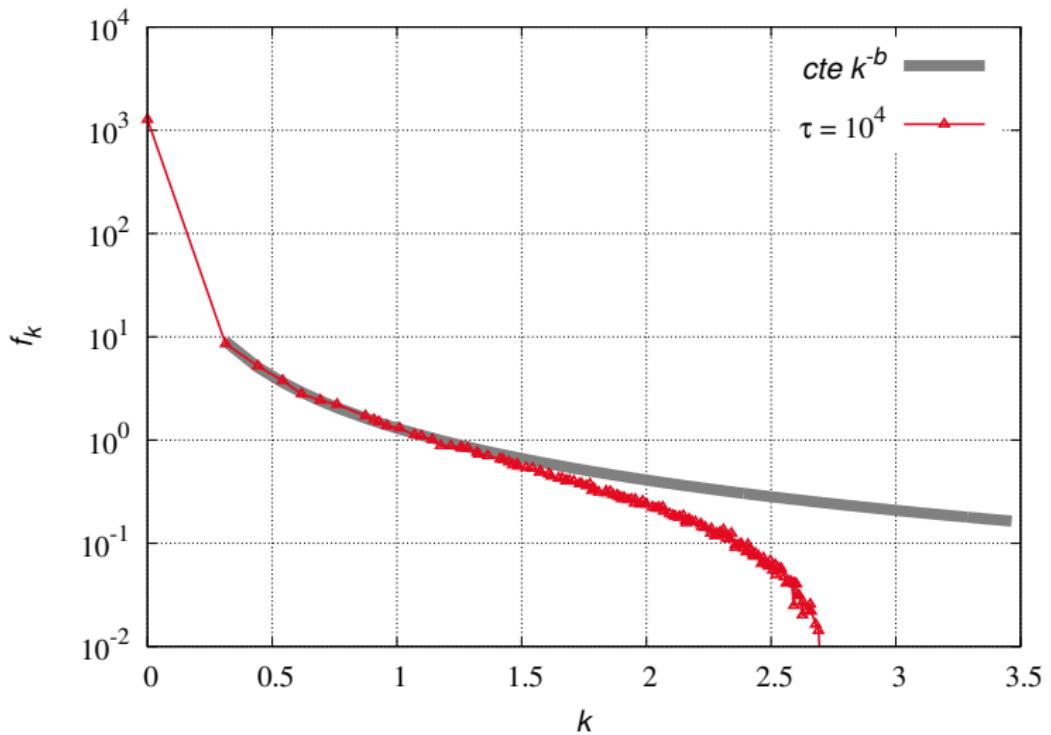
Distribution Function

Bose-Einstein Condensation

TIME EVOLUTION OF THE OCCUPATION NUMBER [TE, GELIS (2011)]



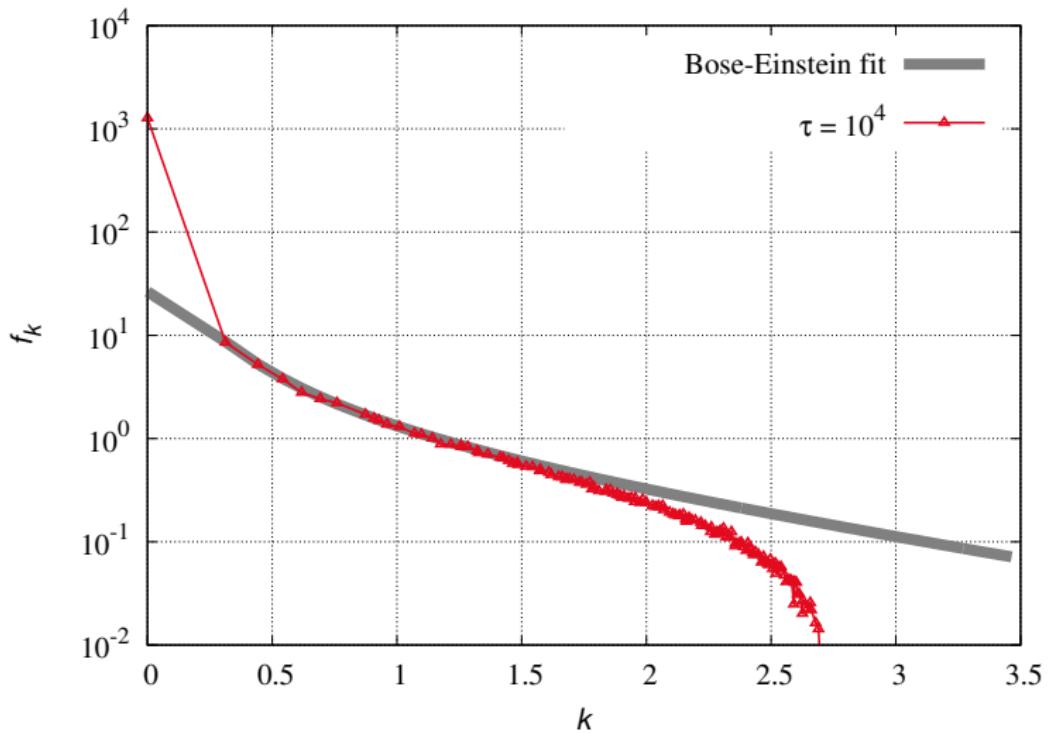
KOLMOGOROV SCALING AT LATE TIMES?



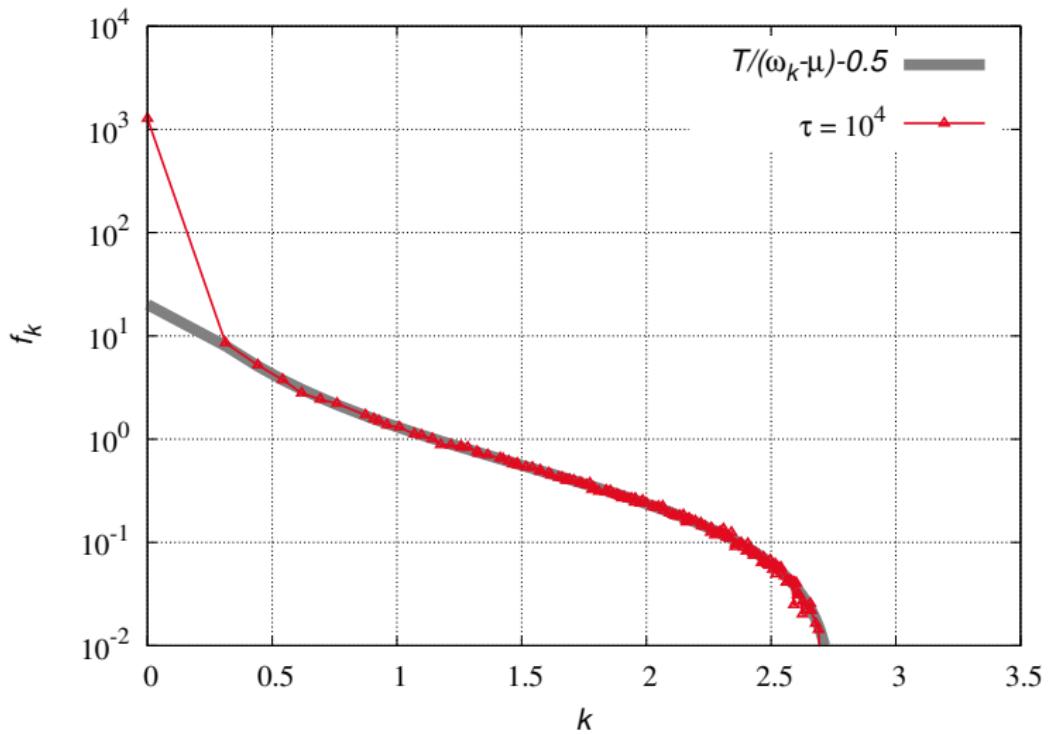
See also:

- R. Micha, I. Tkachev, Phys. Rev. **D 70**, (2004).
- J. Berges, D. Sexty, Phys. Rev. Lett. **108**, (2012).

BOSE-EINSTEIN EQUILIBRIUM DISTRIBUTION?



"CLASSICAL" EQUILIBRIUM DISTRIBUTION



② FIXED VOLUME

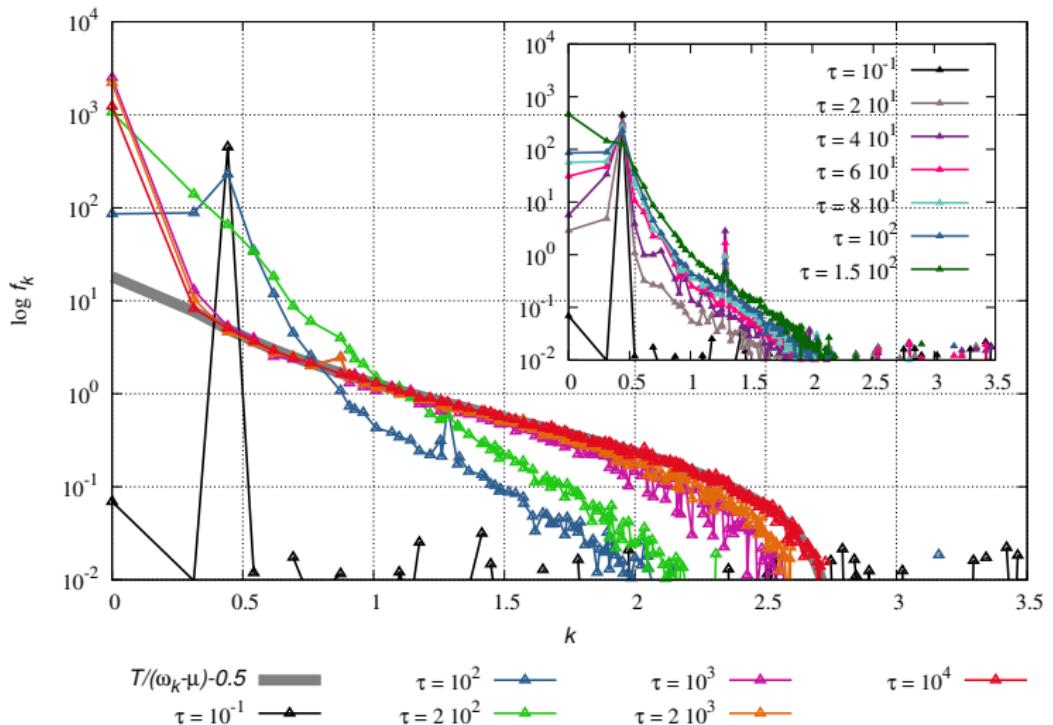
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NON-ZERO INITIAL MODE: $\varphi_0 \sim \cos k.x$

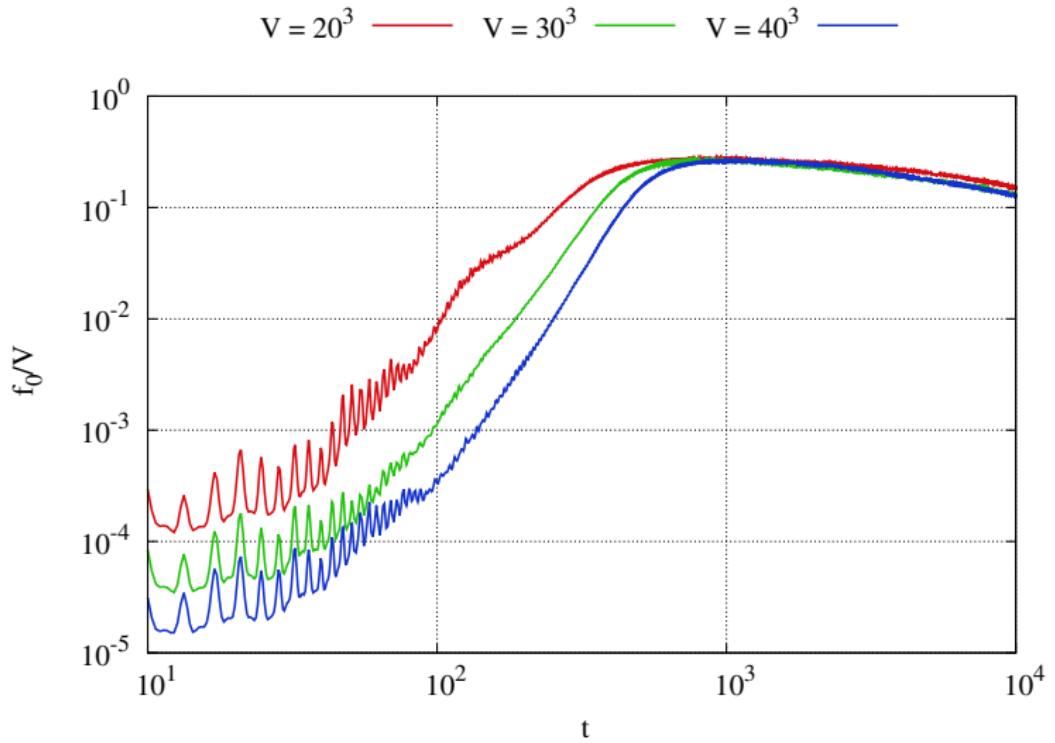


$$f_{\mathbf{k}} = \frac{T}{\omega_{\mathbf{k}} - \mu} - \frac{1}{2} + n_0 \delta(\mathbf{k})$$

implies

$$\frac{f_0}{V} = \text{cte}$$

EVOLUTION OF THE CONDENSATE



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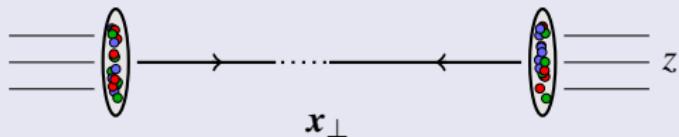
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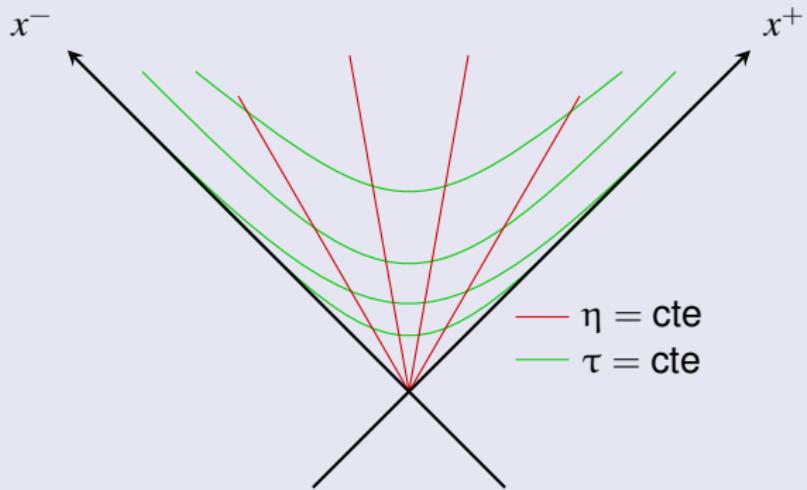
Comparison with Hydro

ADAPTED COORDINATE SYSTEM TO DESCRIBE A HEAVY ION COLLISION?

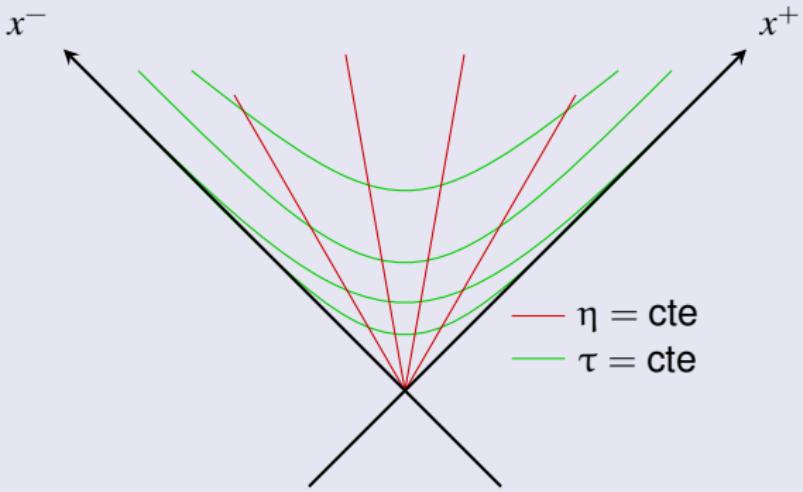


System boost invariant in z direction

PROPER TIME/RAPIDITY COORDINATE SYSTEM



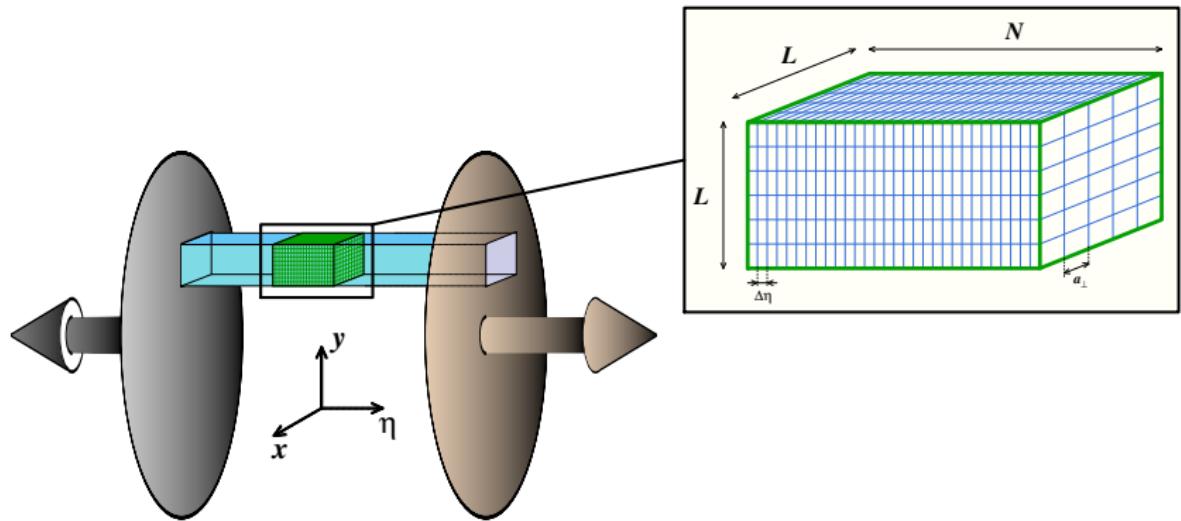
PROPER TIME/RAPIDITY COORDINATE SYSTEM



EOM FOR A BOOST-INVARIANT FIELD

$$\left[\frac{\partial^2}{\partial \tau^2} + \frac{1}{\tau} \frac{\partial}{\partial \tau} - \nabla_{\perp}^2 \right] \varphi_0(\tau, \mathbf{x}_{\perp}) + V'(\varphi_0(\tau, \mathbf{x}_{\perp})) = 0$$

SCALAR FIELD THEORY

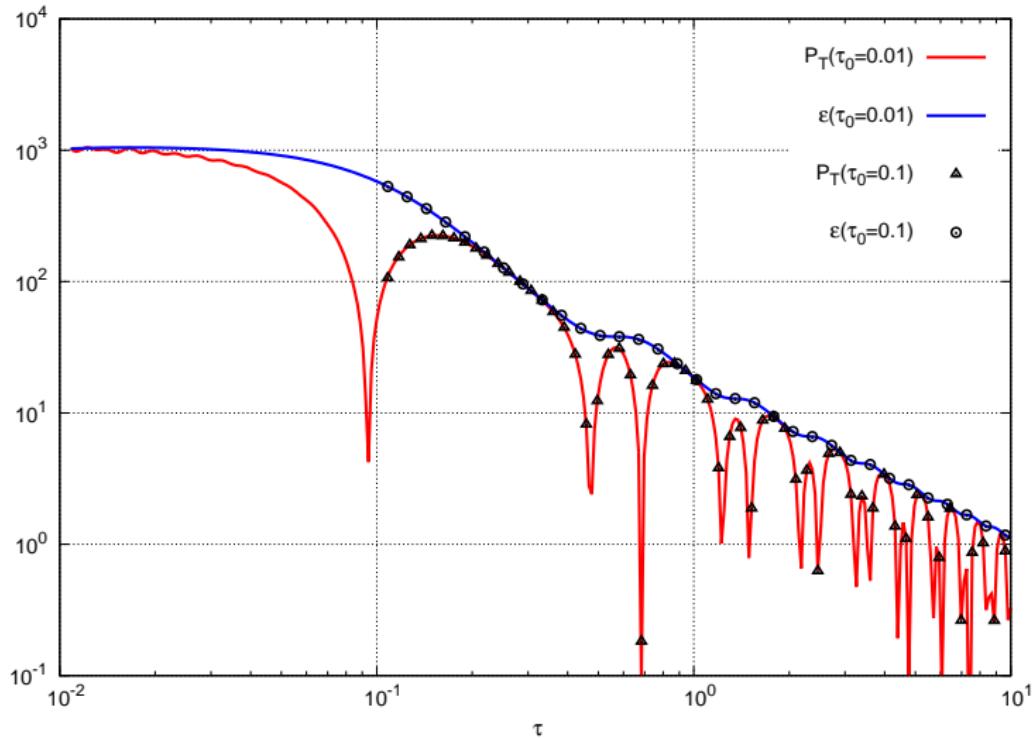


③ EXPANDING VOLUME

The model

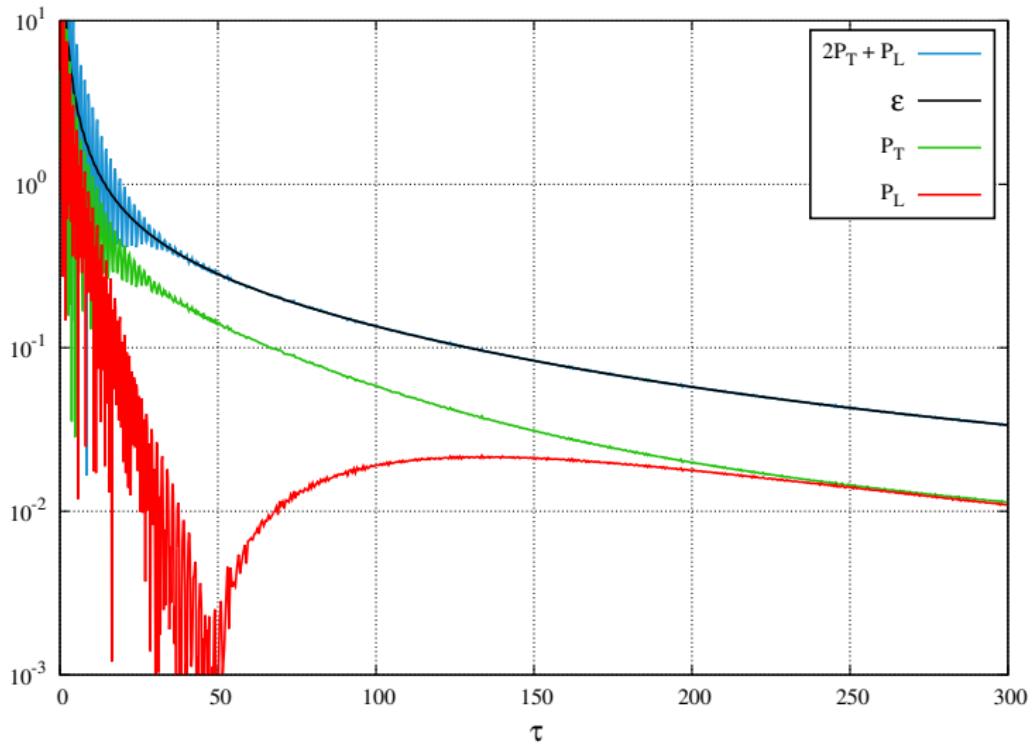
Izotropization

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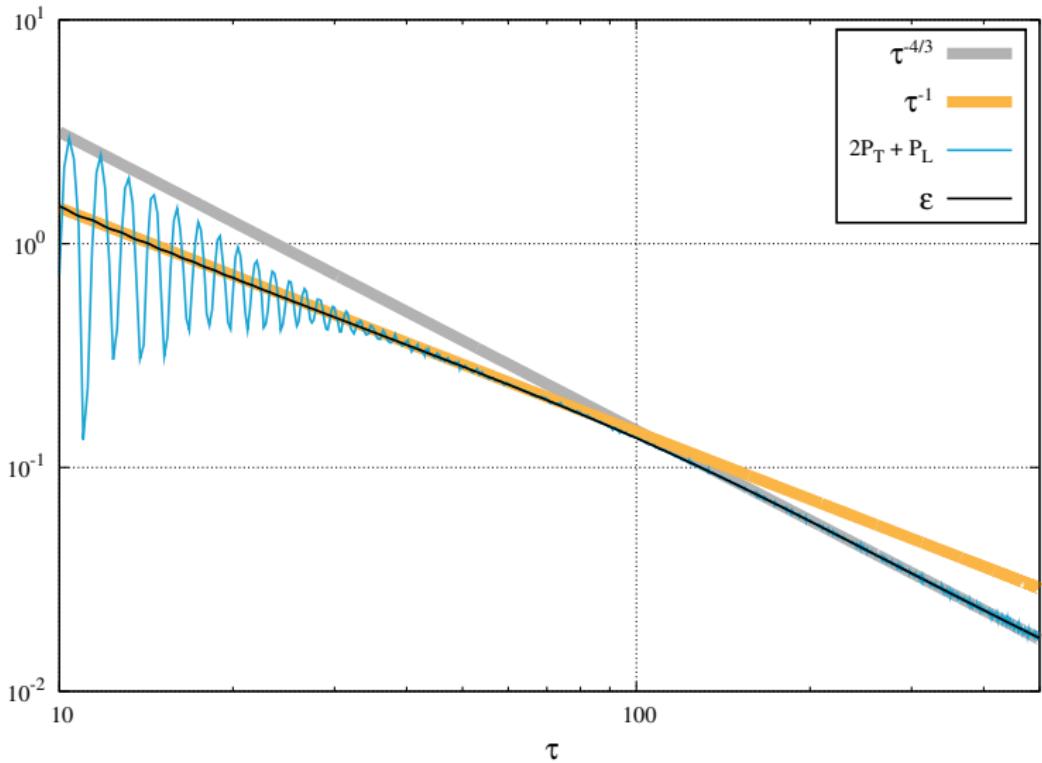


$$T_{\text{RESUM}}^{\mu\nu}$$

[DUSLING, TE, GELIS, VENUGOPALAN (2012)]



ϵ BEHAVIOUR

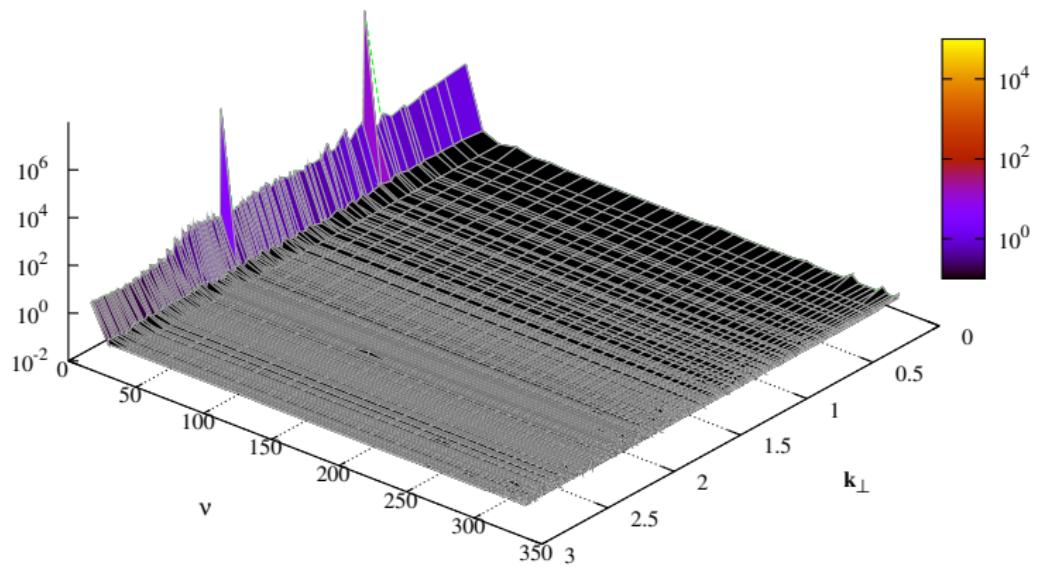


$$\text{Bjorken Law: } \partial_\tau \epsilon + \frac{\epsilon + P_L}{\tau} = 0$$

EVOLUTION OF $f_{\mathbf{v} \mathbf{k}_{\perp}}$



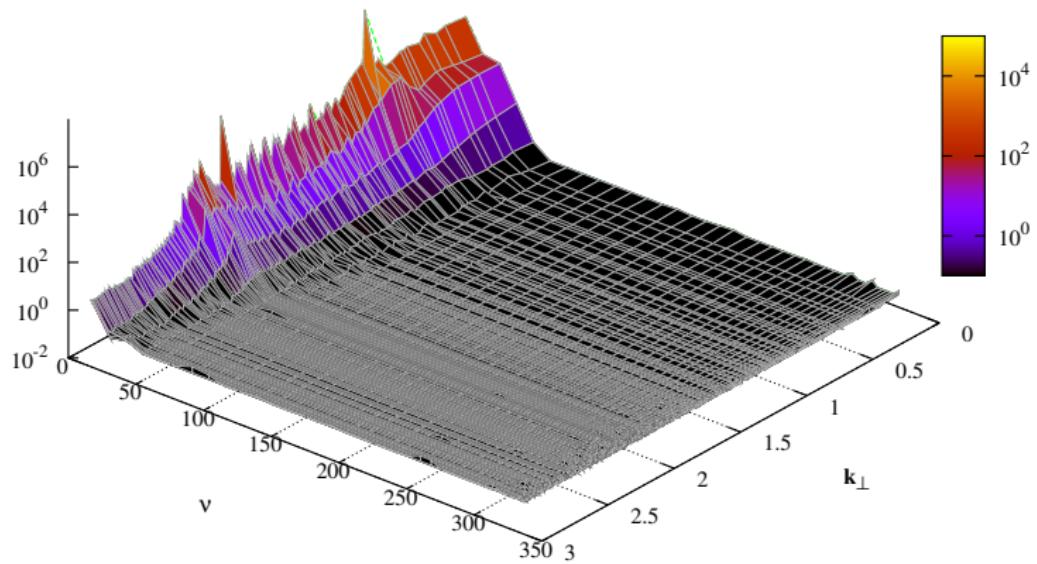
$\tau = 10$ [40 × 40 × 320]



EVOLUTION OF $f_{\mathbf{v} \mathbf{k}_{\perp}}$



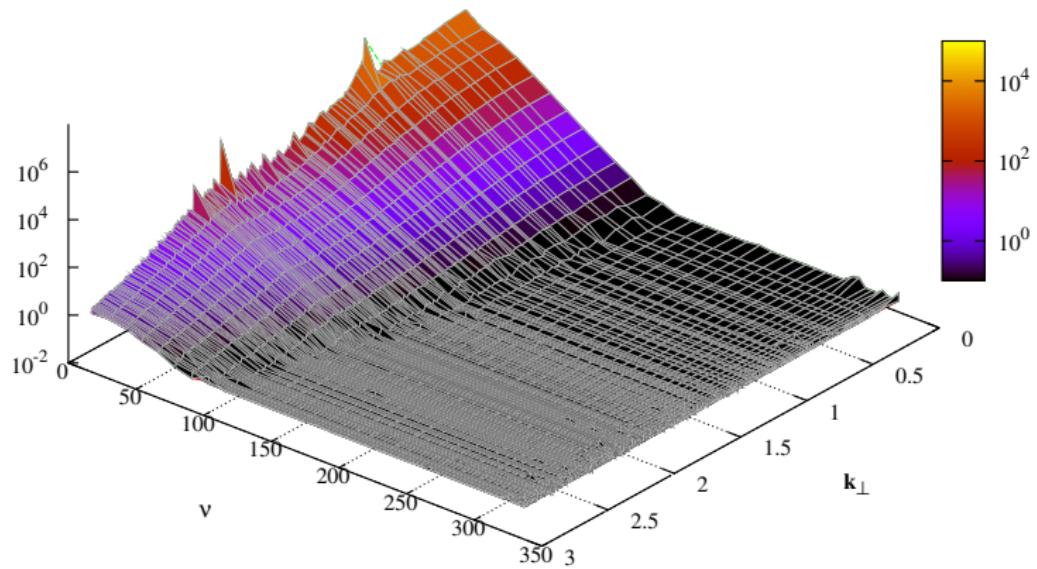
$\tau = 50$ [40 × 40 × 320]



EVOLUTION OF $f_{\mathbf{v} \mathbf{k}_{\perp}}$



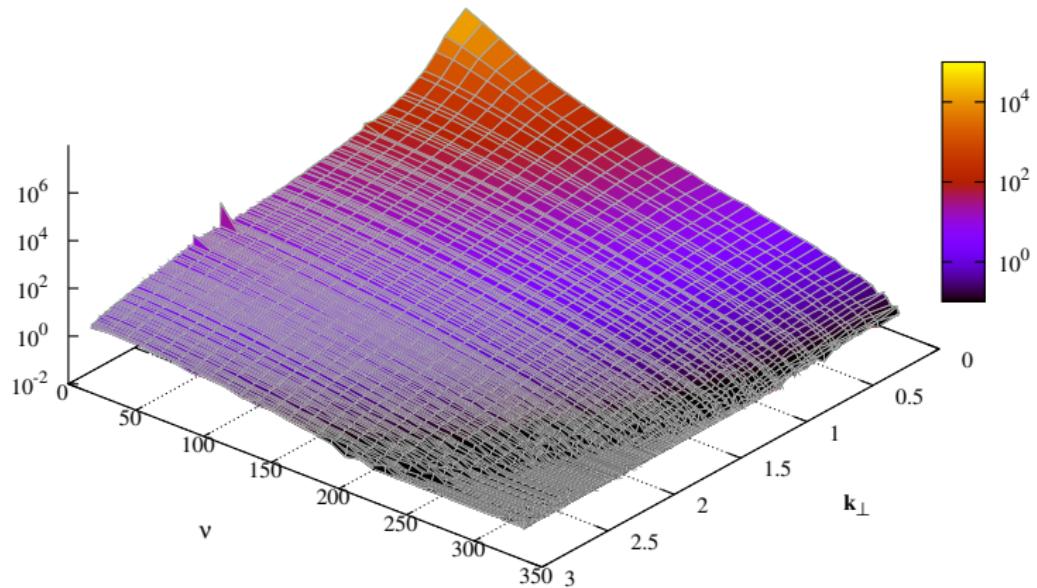
$\tau = 100$ [40 × 40 × 320]



EVOLUTION OF $f_{\mathbf{v} \mathbf{k}_\perp}$



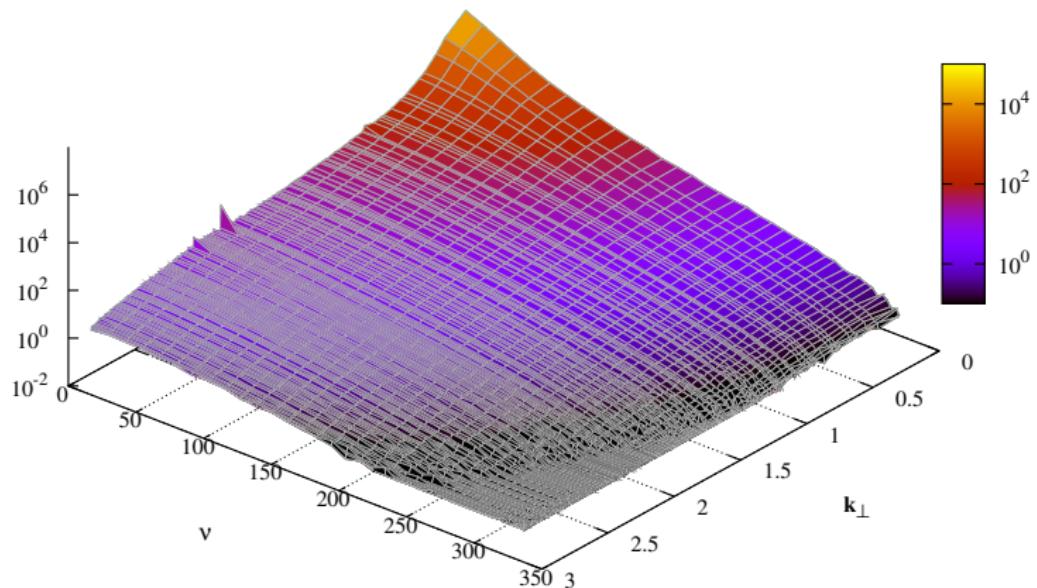
$\tau = 300$ [40 × 40 × 320]



EVOLUTION OF $f_{\nu k_\perp}$



$\tau = 300$ [40 × 40 × 320]



$$\nu \propto \tau^{2/3}$$

③ EXPANDING VOLUME

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Equation of state: $\epsilon = 2P_L + P_T$

IDEAL HYDRO

Isotropic system

$$T_{\text{ideal}}^{\mu\nu} = \epsilon u^\mu u^\nu - P(g^{\mu\nu} - u^\mu u^\nu)$$

Equation of state: $\epsilon = 2P_L + P_T$

IDEAL HYDRO

Isotropic system

$$T_{\text{ideal}}^{\mu\nu} = \epsilon u^\mu u^\nu - P(g^{\mu\nu} - u^\mu u^\nu)$$

VISCOUS HYDRO

Anisotropic system

$$T^{\mu\nu} = T_{\text{ideal}}^{\mu\nu} + \eta \pi^{\mu\nu}$$

In our case

$$P_T = \frac{\epsilon}{3} + \frac{2\eta}{3\tau} \quad P_L = \frac{\epsilon}{3} - \frac{4\eta}{3\tau}$$

BJORKEN's Law (coming from $\partial_\mu T^{\mu\nu} = 0$):

$$\partial_\tau \epsilon + \frac{\epsilon + P_L}{\tau} = 0 \rightarrow \partial_\tau \epsilon + \frac{4}{3} \frac{\epsilon}{\tau} = \frac{4}{3} \frac{\eta}{\tau^2}$$

BJORKEN's Law (coming from $\partial_\mu T^{\mu\nu} = 0$):

$$\partial_\tau \epsilon + \frac{\epsilon + P_L}{\tau} = 0 \rightarrow \partial_\tau \epsilon + \frac{4}{3} \frac{\epsilon}{\tau} = \frac{4}{3} \frac{\eta}{\tau^2}$$

assuming $\eta = \frac{\eta_0}{\tau}$ and STEFAN-BOLTZMANN entropy $s \approx \epsilon^{\frac{3}{4}}$

$$\partial_\tau \epsilon + \frac{4}{3} \frac{\epsilon}{\tau} = \frac{4}{3} \underbrace{\frac{\eta}{s}}_{\text{cte}} \frac{\epsilon^{\frac{3}{4}}}{\tau^2}$$

BJORKEN's Law (coming from $\partial_\mu T^{\mu\nu} = 0$):

$$\partial_\tau \epsilon + \frac{\epsilon + P_L}{\tau} = 0 \rightarrow \partial_\tau \epsilon + \frac{4}{3} \frac{\epsilon}{\tau} = \frac{4}{3} \frac{\eta}{\tau^2}$$

assuming $\eta = \frac{\eta_0}{\tau}$ and STEFAN-BOLTZMANN entropy $s \approx \epsilon^{\frac{3}{4}}$

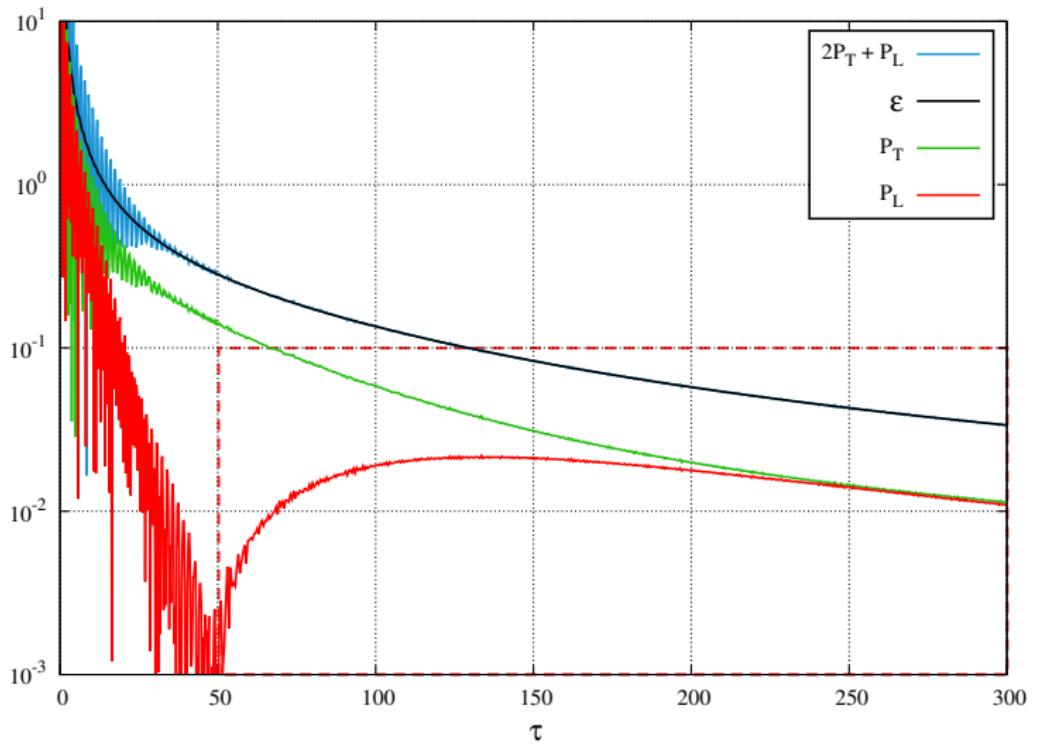
$$\partial_\tau \epsilon + \frac{4}{3} \frac{\epsilon}{\tau} = \frac{4}{3} \underbrace{\frac{\eta}{s}}_{\text{cte}} \frac{\epsilon^{\frac{3}{4}}}{\tau^2}$$

At a given time, knowing ϵ, P_T, P_L and assuming

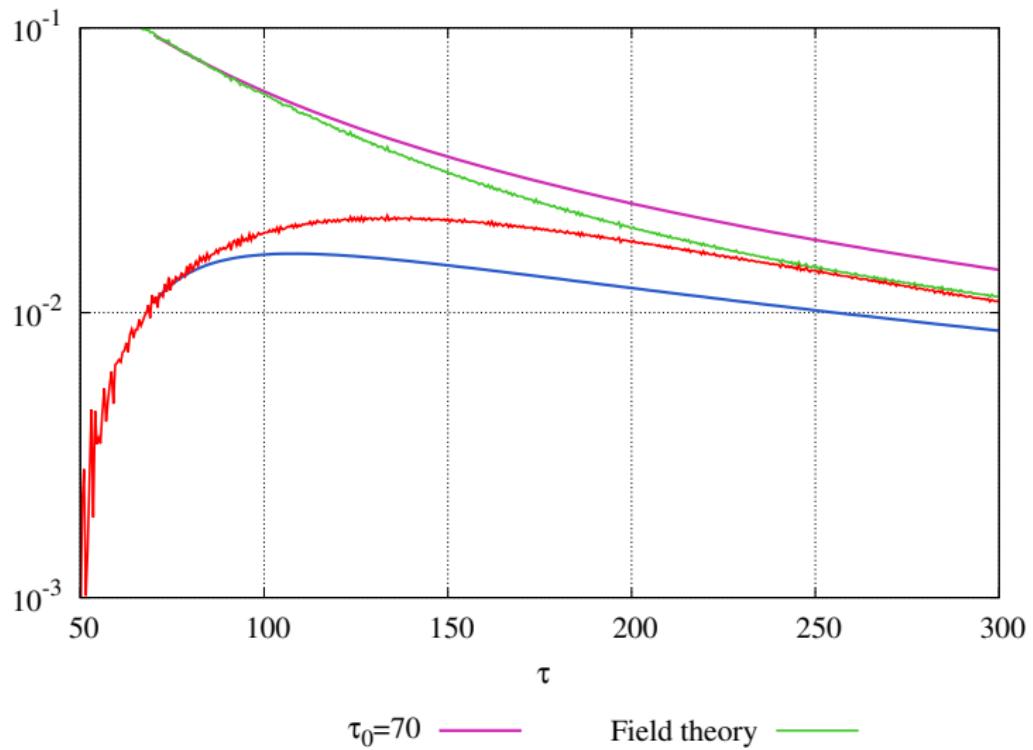
- an EOS
- STEFAN-BOLTZMANN entropy
- $\eta = \frac{\eta_0}{\tau}$
- $\frac{\eta}{s} = \text{cte}$

gives a very simple hydro model.

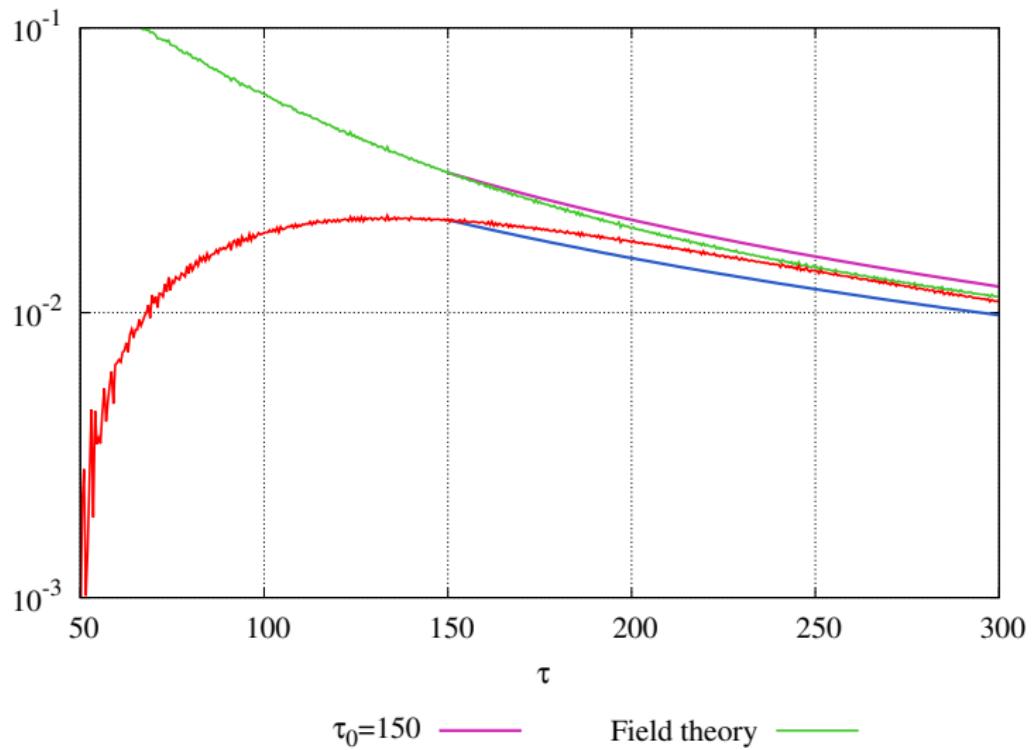
COMPARISON WITH HYDRO



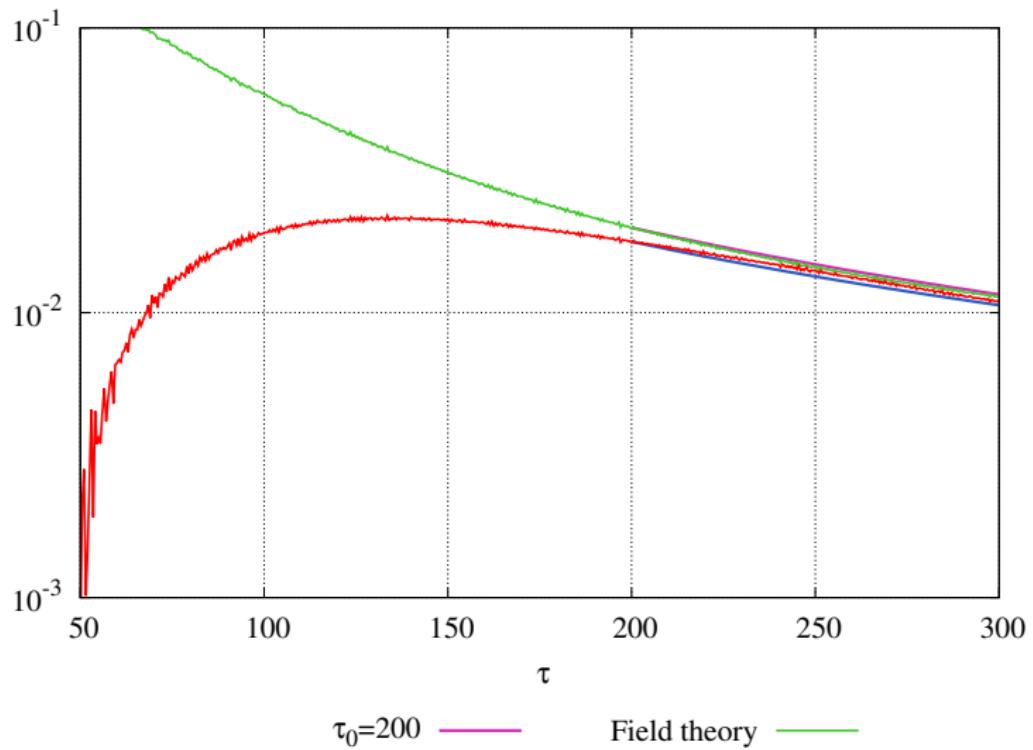
COMPARISON WITH HYDRO



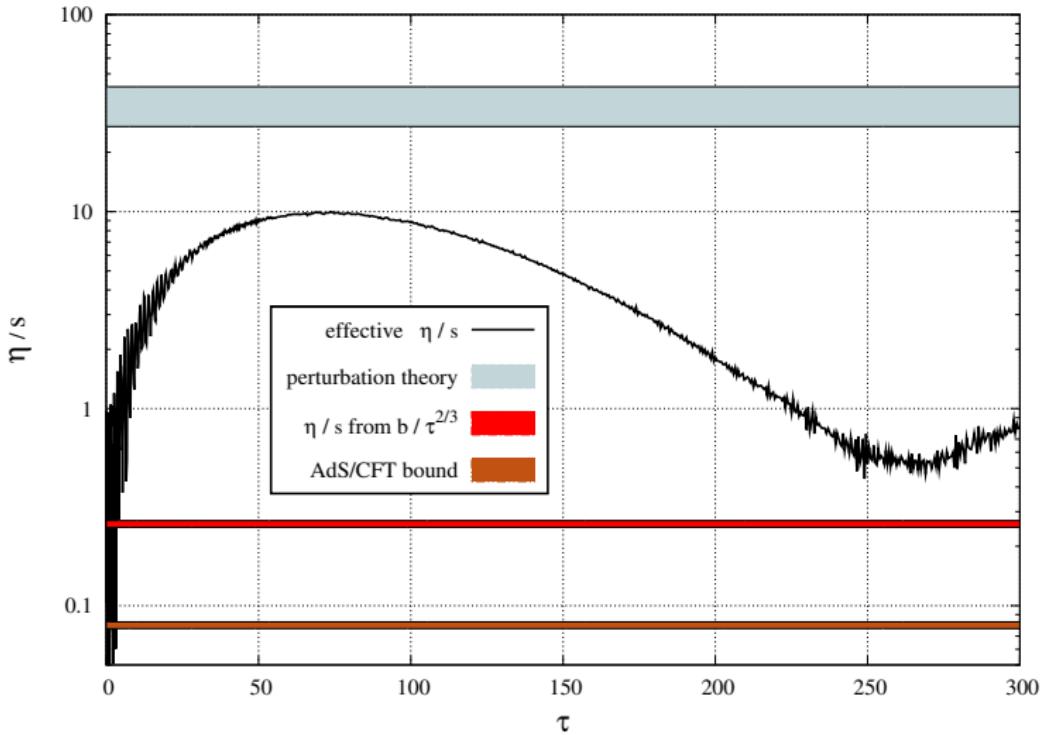
COMPARISON WITH HYDRO



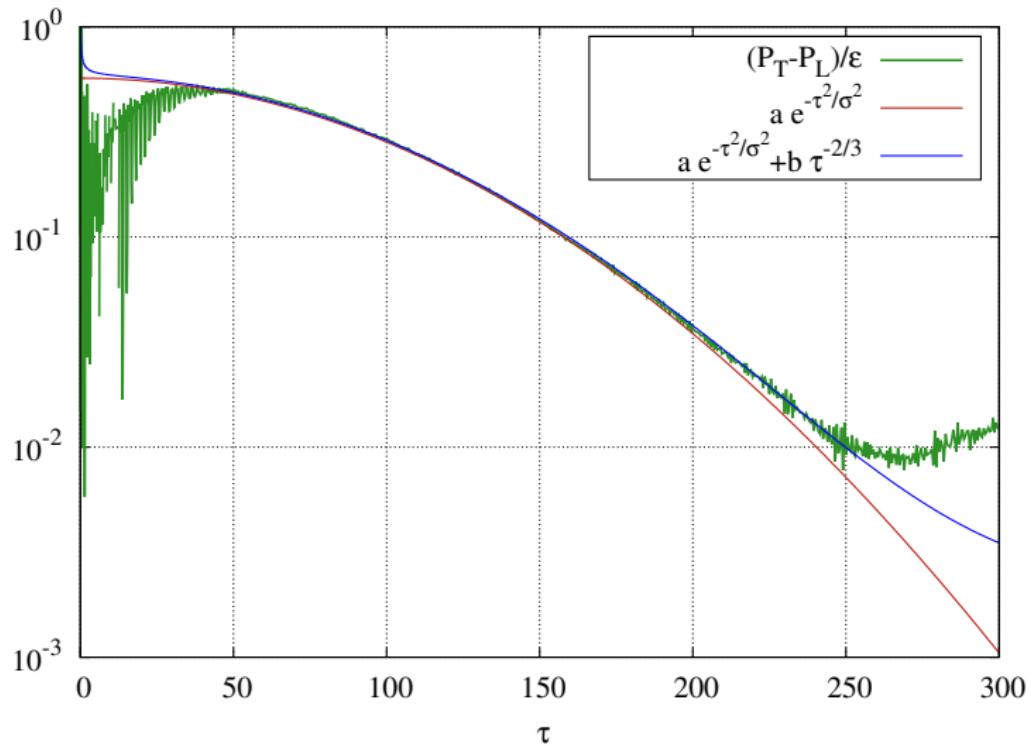
COMPARISON WITH HYDRO



COMPARISON WITH HYDRO



COMPARISON WITH HYDRO



GAUGE CASE
[DUSLING, GELIS, VENUGOPALAN (2011)]



- EOM for the field → Yang-Mills Equations known

GAUGE CASE
[DUSLING, GELIS, VENUGOPALAN (2011)]



- EOM for the field → Yang-Mills Equations known
- Linearized EOM for the small fluctuations → known

GAUGE CASE
[DUSLING, GELIS, VENUGOPALAN (2011)]



- EOM for the field → Yang-Mills Equations known
- Linearized EOM for the small fluctuations → known
- Final form of the solution → known

GAUGE CASE
[DUSLING, GELIS, VENUGOPALAN (2011)]



- EOM for the field → Yang-Mills Equations known
- Linearized EOM for the small fluctuations → known
- Final form of the solution → known
- initial A_μ^a

$$A_\mu^a(\tau_0, \eta, \mathbf{x}_\perp) = \mathcal{A}_\mu^a(\mathbf{x}_\perp) + \text{Re} \int_k c_k a_{\mu k}^a(\tau_0, \eta, \mathbf{x}_\perp)$$

GAUGE CASE
[DUSLING, GELIS, VENUGOPALAN (2011)]



- EOM for the field → Yang-Mills Equations known
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- Final form of the solution → known
- initial A_μ^a

$$A_\mu^a(\tau_0, \eta, \mathbf{x}_\perp) = \mathcal{A}_\mu^a(\mathbf{x}_\perp) + \text{Re} \int_k c_k a_{\mu k}^a(\tau_0, \eta, \mathbf{x}_\perp)$$

- Lattice implementation → Work in progress

Decoherence

Decoherence



Equation of state $\epsilon = 3P$ (isotropy by construction)

Decoherence



Equation of state $\epsilon = 3P$ (isotropy by construction)



BOSE-EINSTEIN condensate

Decoherence



Equation of state $\epsilon = 3P$ (isotropy by construction)



BOSE-EINSTEIN condensate



$$f_{\mathbf{k}} \propto \frac{T}{\omega_{\mathbf{k}}^{-m}} - \frac{1}{2}$$

Decoherence

Decoherence



Equation of state $\epsilon = 2P_T + P_L$

Decoherence



Equation of state $\epsilon = 2P_T + P_L$



Instabilities [BERGES, BOGUSLAVSKI, SCHLICHTING (2012)]

Decoherence



Equation of state $\epsilon = 2P_T + P_L$



Instabilities [BERGES, BOGUSLAVSKI, SCHLICHTING (2012)]



Isotropization

Decoherence



Equation of state $\epsilon = 2P_T + P_L$



Instabilities [BERGES, BOGUSLAVSKI, SCHLICHTING (2012)]



Isotropization



Thermal distribution



BOSE-EINSTEIN condensate?

Thank You!