

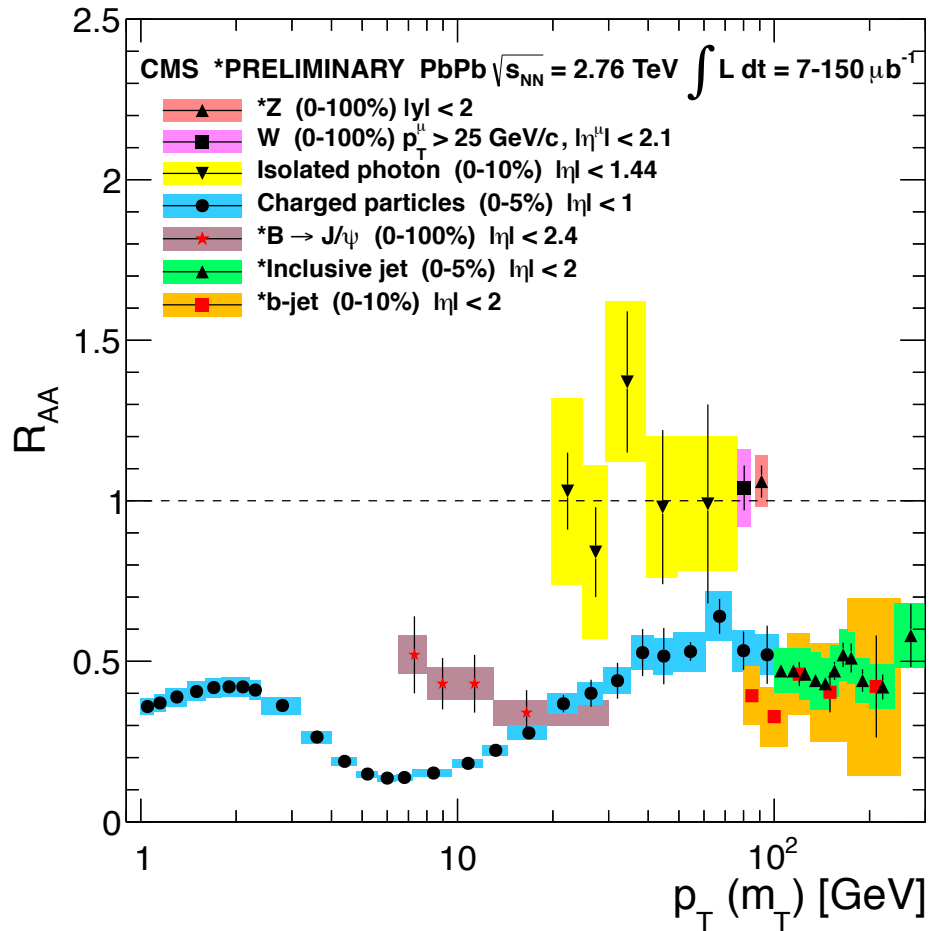
+

W
production
in pPb
collision

$$L = 34.6 \text{ nb}^{-1}$$

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Lamia Benhabib (CERN)

+ Z and W boson: medium blind



Z and W R_{AA} :

- Compatible with binary scaling ($R_{AA}=1$)
- Large difference between electroweak bosons and QCD probes
- A R_{AA} compatible with 1 indicates that there are no strong nuclear effect but:
 - the effect might be symmetric
 - The error does not allow to see small initial cold effects

PLB 710 (2012) 256
 PLB 715 (2012) 66
 CMS PAS HIN-13-004

EPJC 72 (2012) 1945
 JHEP 10 (2012) 087
 PRC 84 (2011) 024906
 CMS PAS HIN-12-003



+ Initial state observable

nPDF: probability density to find a parton with longitudinal momentum x at a resolution scale Q^2

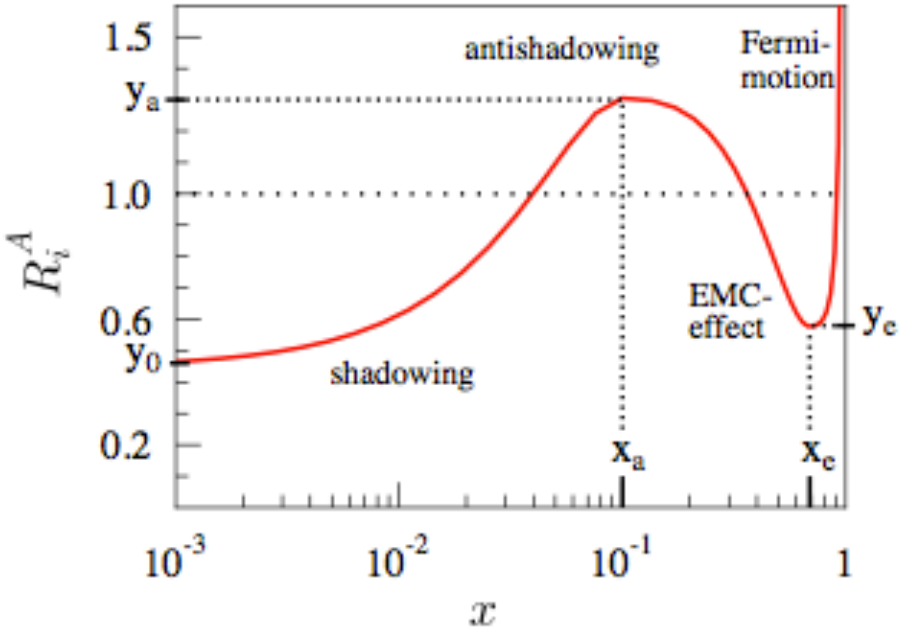
$$f^A = \frac{Z}{A} f^{\text{proton},A} + \frac{N}{A} f^{\text{neutron},A}$$

nPDF

Bound proton PDF

$$f^{\text{proton},A}(x, Q^2) \equiv R_f(x, Q^2) f^{\text{free proton}}(x, Q^2)$$

Nuclear modification factor



+ W in pPb: physics motivation

Aim: probe the nPdf

With: the W boson in collisions pPb

- The collision is asymmetric
- W boson is sensitive to quark flavor \longrightarrow **W⁺, W⁻**
- W boson are produced 10 times more than Z boson

How: we propose 4 asymmetries results

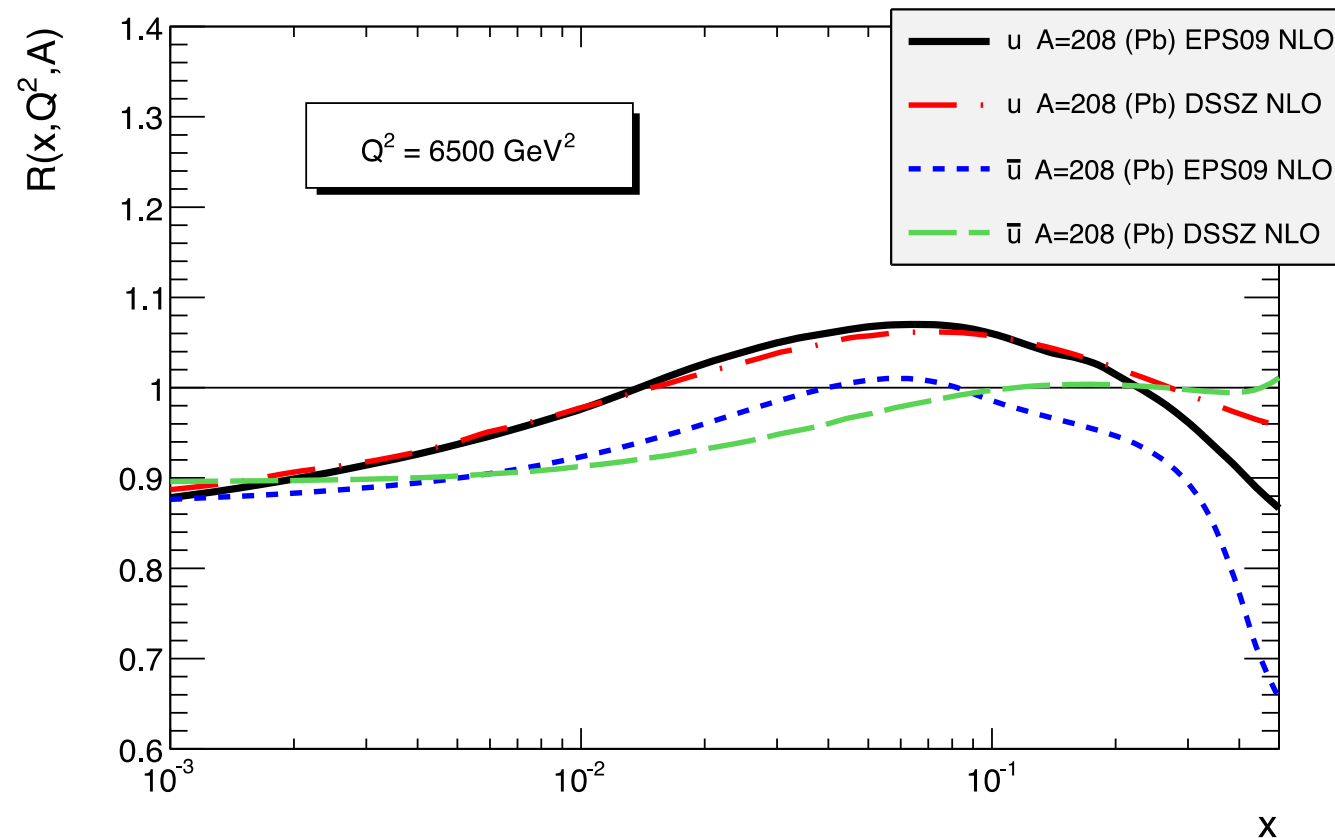
- charge asymmetry
- forward/backward W⁺
- forward/backward W⁻
- forward/backward W

Details:

- The first 3 are inspired by Salgado et al. JHEP03 (2011)071
- The last one is an asymmetry we proposed to replace A2 from their paper
- These 4 results are compared to two predicted PDFs set
 - CT10 (proton Pdf)
 - EPS09 (proton Pdf with nuclear effect)

+ W and nPDFs

- $Q^2 \sim 6500 \text{ GeV}^2$ ($\sim m_w^2$)
- The asymmetries will probe
 - antiquarks in the nucleus at small $x=[0.002;0.02]$ (**shadowing** region)
 - quarks at large $x=[0.02;0.3]$ (**anti-shadowing** region)



+ Analysis strategy

- **Single lepton selection**
- **MET fitting**
- **Single lepton efficiency**
- **Scales Factors on single lepton Tag and Probe**
- **Cross section**
- **Asymmetry**

$$N^{\text{corr}} = \frac{N^{\text{raw}}}{\epsilon_{\text{all}}^{\text{MC,W}} \times \text{SF}}$$

+ Analysis strategy

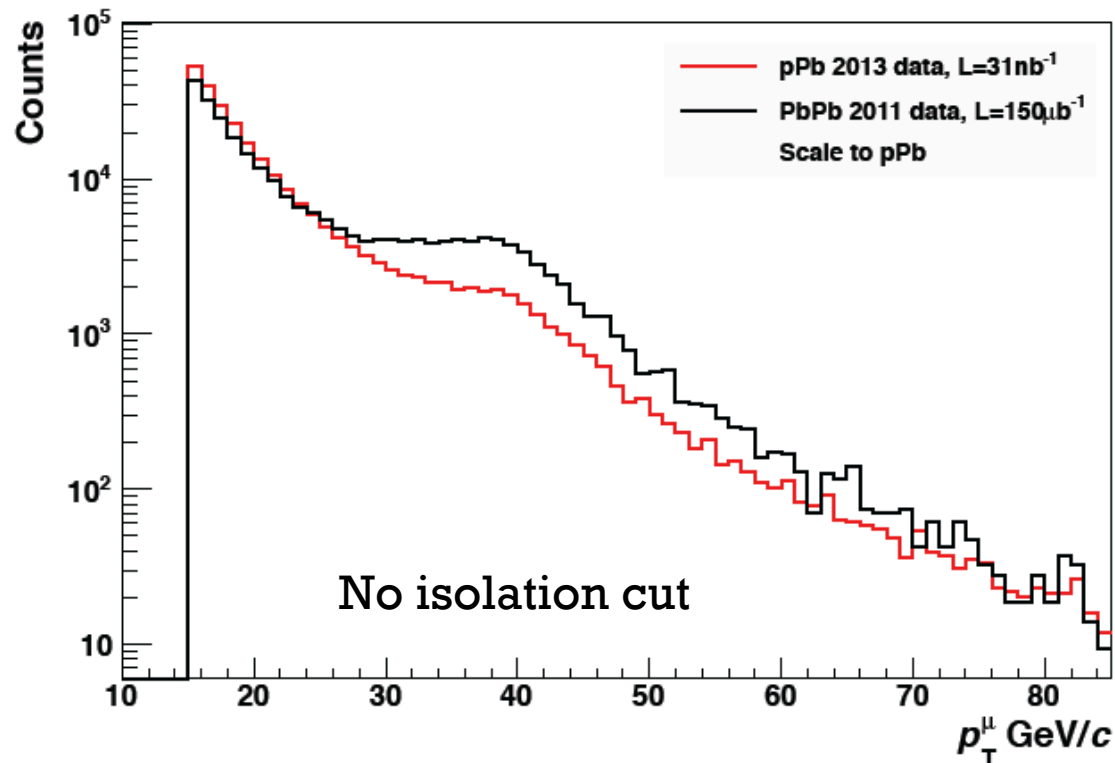
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+ Single lepton selection

1) Isolated lepton

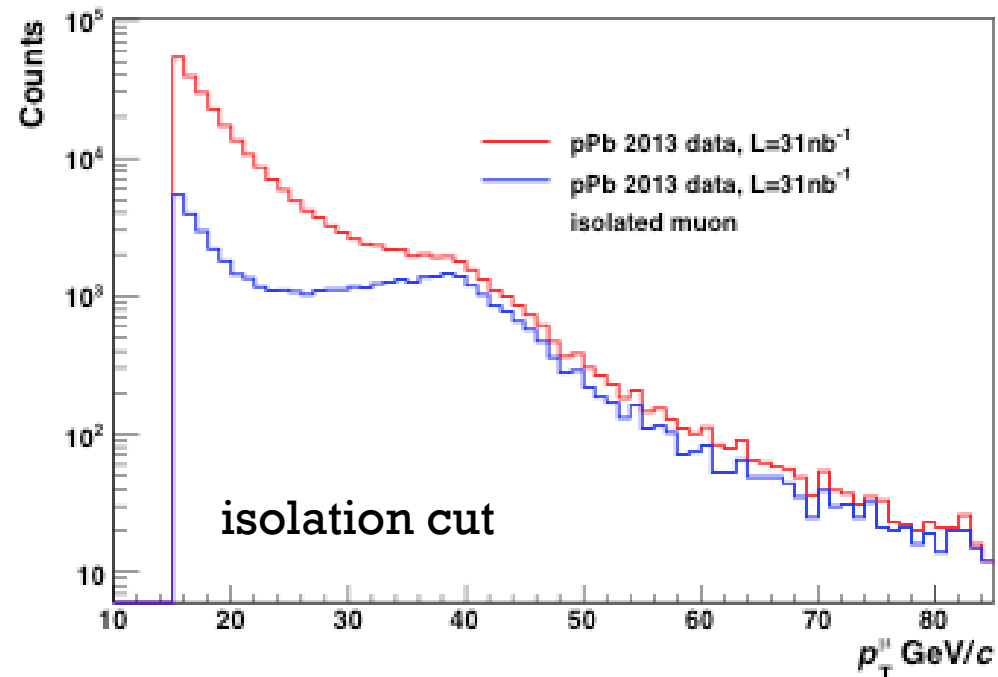
PbPb and pPb comparison



- Signal more visible in PbPb than in pPb
- Background suppression because of jet quenched?
- First use of isolation on lepton in Pb

+ Single lepton selection

1) Isolated lepton



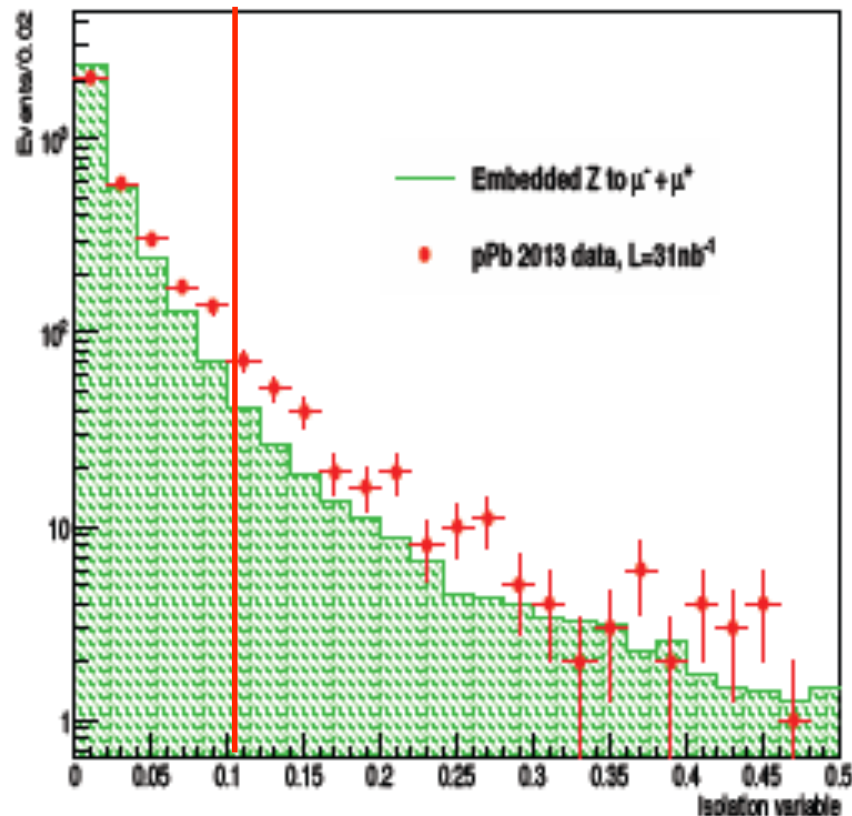
- Sum over the Particle Flow candidates inside a cone $\Delta R = 0.3$
With $\Delta R = \sqrt{(\eta_\mu - \eta_{\text{PF}})^2 + (\Phi_\mu - \Phi_{\text{PF}})^2}$
- A muon is considered as “isolated” if the sum inside the cone is lower than **10% of its p_T**

+ Single lepton selection

1) Isolated lepton

- **Caveat**

Z events from data and MC



- MC does not reproduce data perfectly. Several % of difference around 0.1.
- The centrality weight ensures global multiplicity agreement but not local one. More details about multiplicity in back up.
- The difference between the isolation efficiency in data and MC will be corrected and controlled with Scale Factors based on TnP

+ Single lepton selection

3) Z veto

High quality lepton = trigger+
selection cut + isolation

Muon:

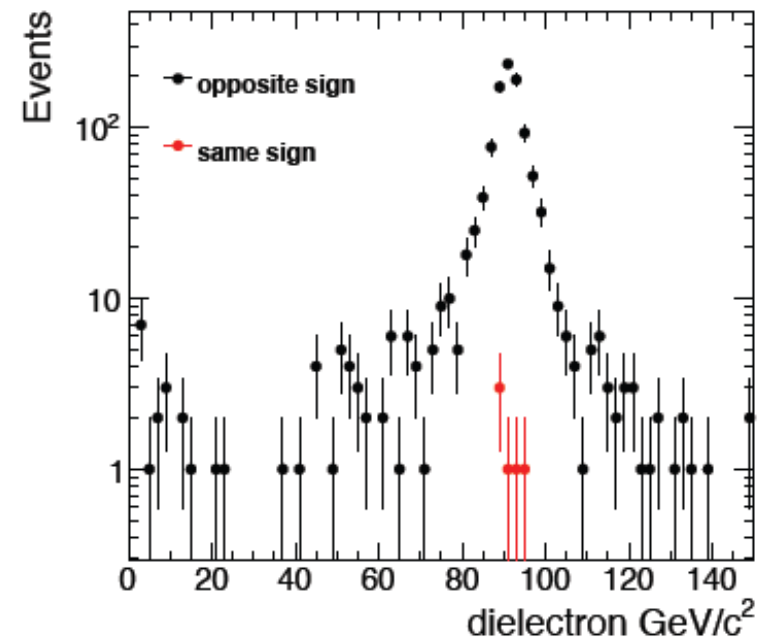
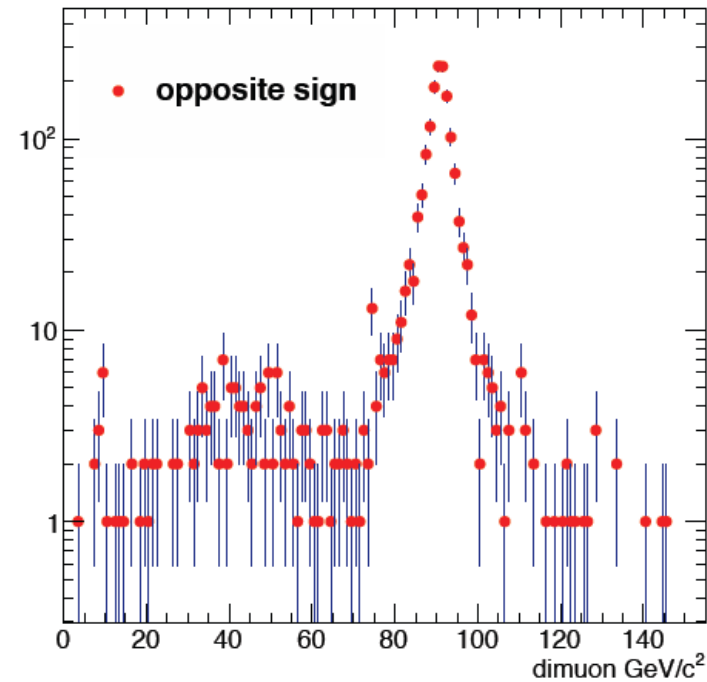
Event with 2 high quality leptons,
with $l_1 > 25$ GeV and $l_2 > 15$ GeV
are removed.

No invariant mass cut, since no
same sign pair can be found in the
mass spectra.

Electron:

Event with 2 high quality leptons,
with $l_1 > 25$ GeV and $l_2 > 10$ GeV
are removed.

No invariant mass cut, since only 6
same sign pair.



+ Indirect neutrino reconstruction

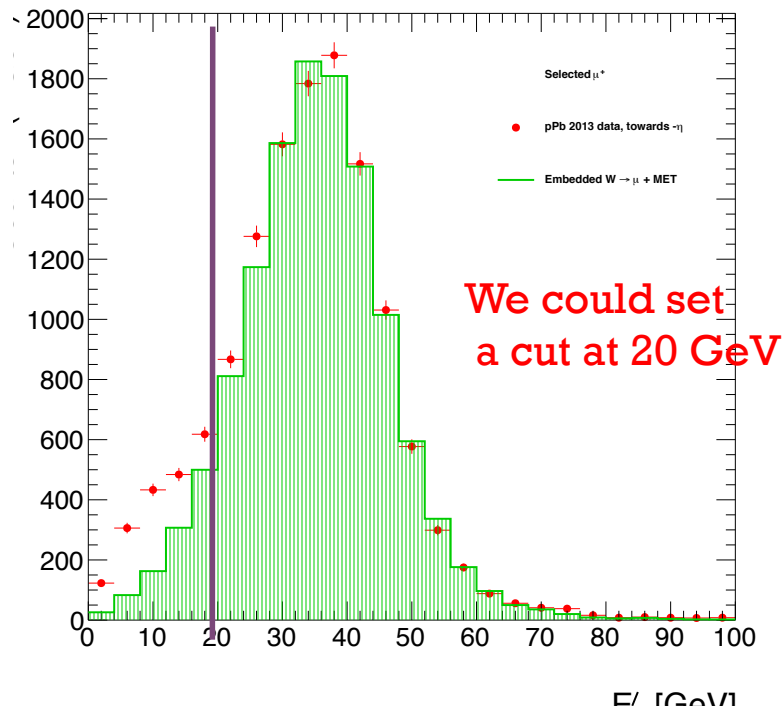
Unbalanced energy in the transverse plane

$$\cancel{E}_{Tx} = -\sum_{\text{particles}} E_{T_x}^{\vec{}} = -\sum_{\text{particles}} E_{Tx} \cos(\phi)$$

$$\cancel{E}_{Ty} = -\sum_{\text{particles}} E_{T_y}^{\vec{}} = -\sum_{\text{particles}} E_{Ty} \sin(\phi)$$

$$\cancel{E}_T = -\sum_{\text{particles}} E_T^{\vec{}} = \sqrt{\cancel{E}_{Tx}^2 + \cancel{E}_{Ty}^2}$$

Particles = Particles Flow candidates: muon, electron, photon, charged hadrons and neutral hadrons



+ Indirect neutrino reconstruction

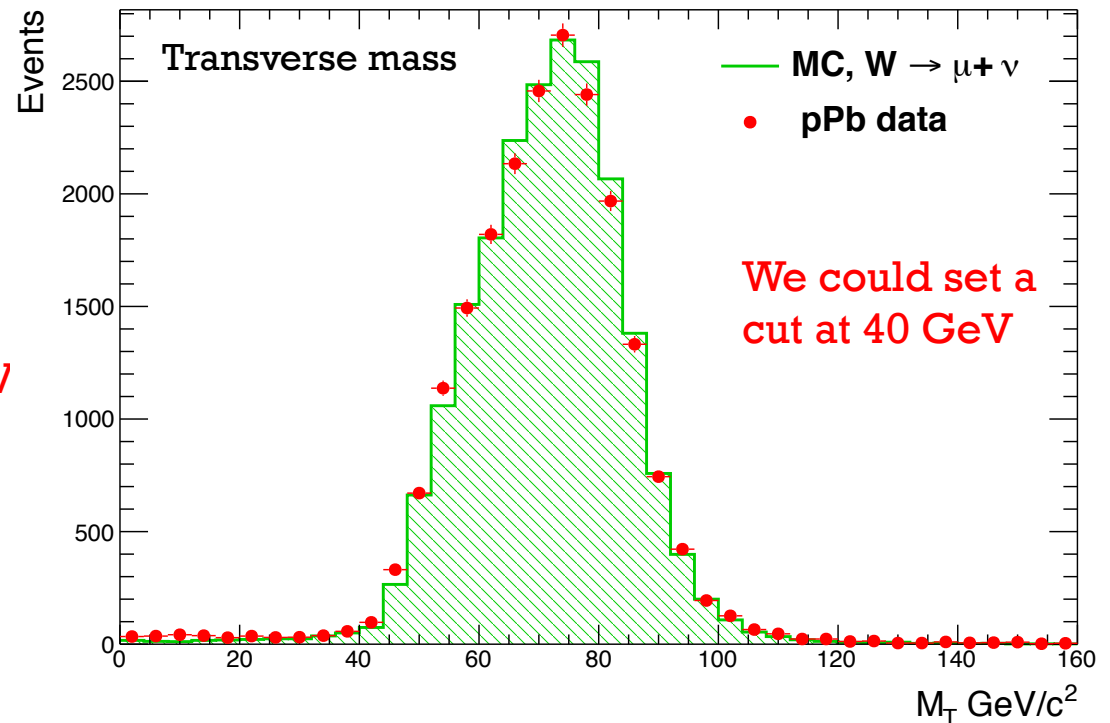
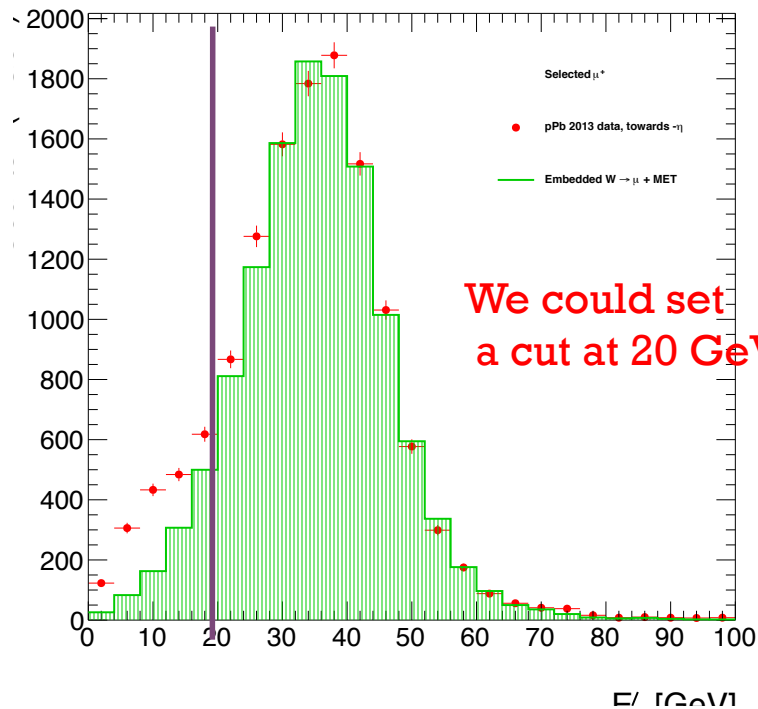
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$$\cancel{E}_T = -\sum_{\text{particles}} E_T^{\vec{T}} = \sqrt{\cancel{E}_{Tx}^2 + \cancel{E}_{Ty}^2}$$

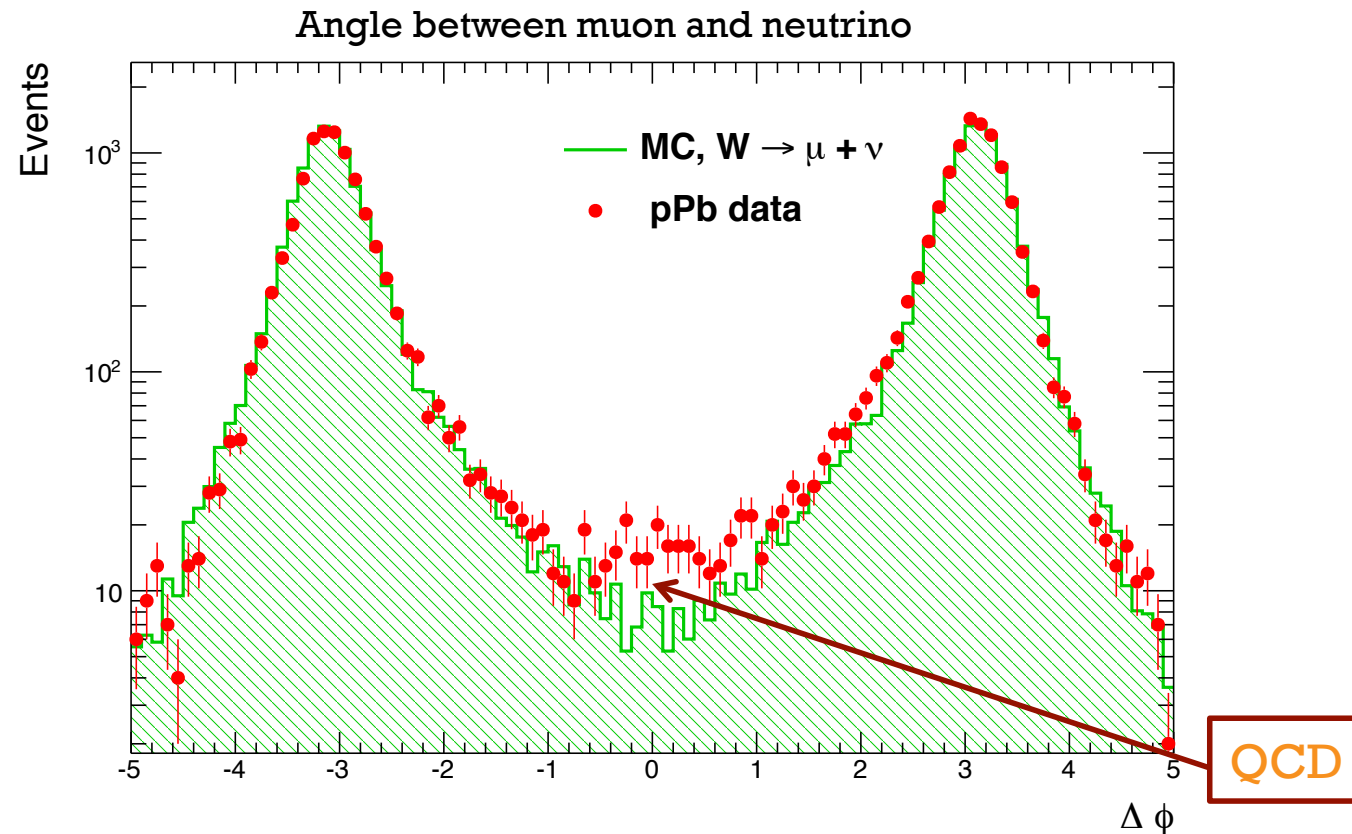
Particles = Particles Flow candidates: muon, electron, photon, charged hadrons and neutral hadrons



+ Extracting the W yields with cuts?

From the MET and m_T cuts, we get a sample with only 3-6% background left:

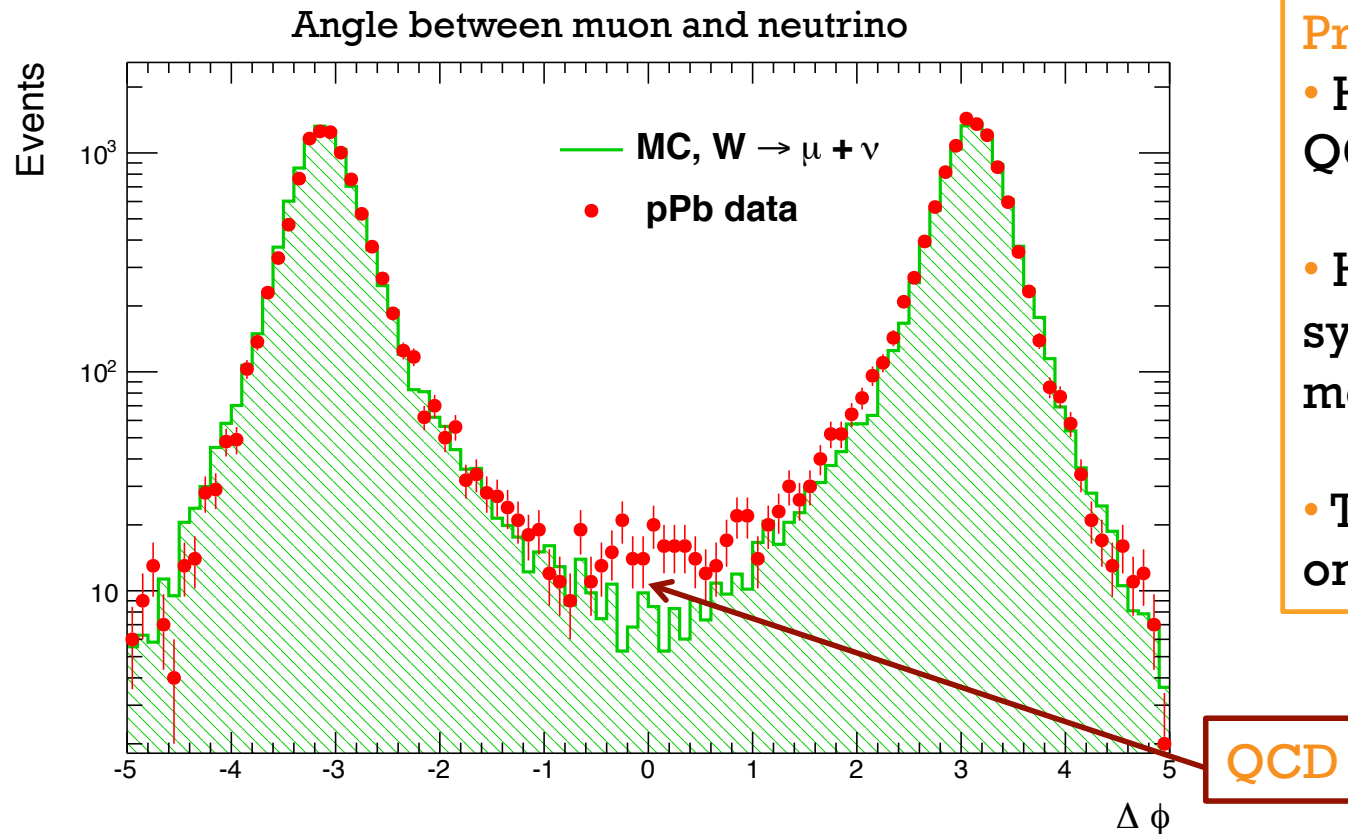
- QCD with heavy quarks inside jets decaying into $\mu + \nu$
- $Z \rightarrow \mu^+ + \mu^-$ with one of the two muons passing the analysis cut and the other outside the acceptance
- $W \rightarrow \tau + \nu \rightarrow \mu + \nu$



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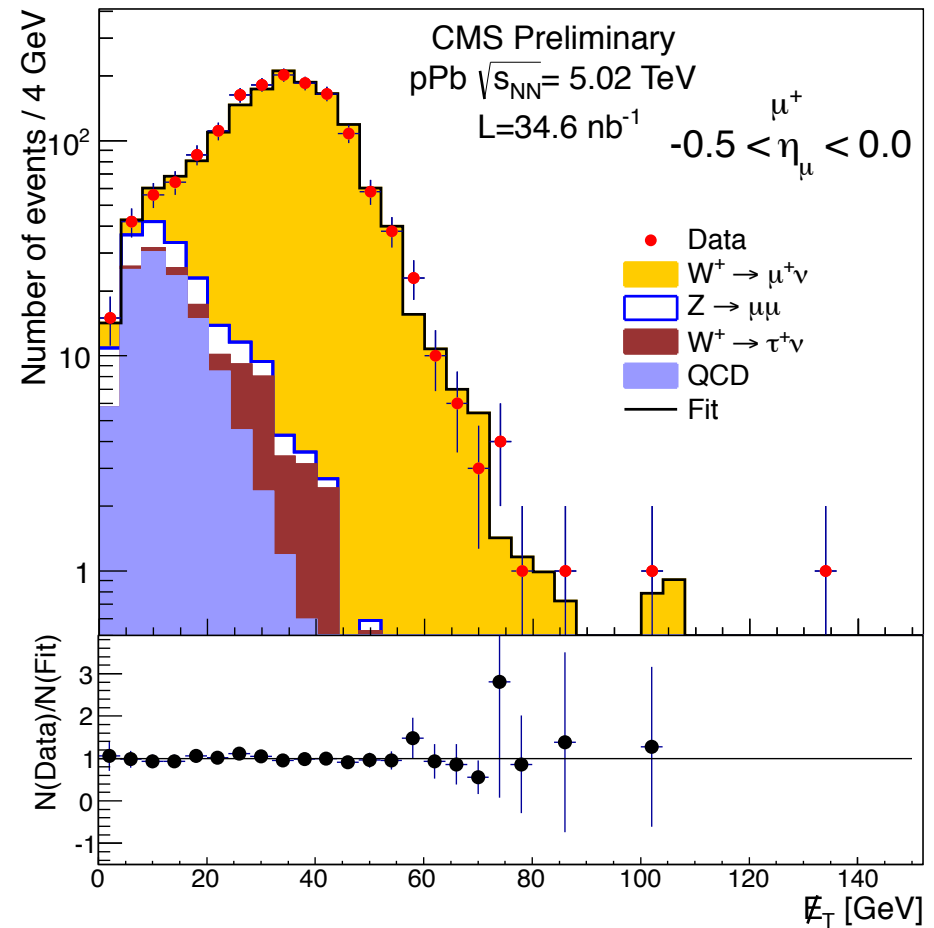
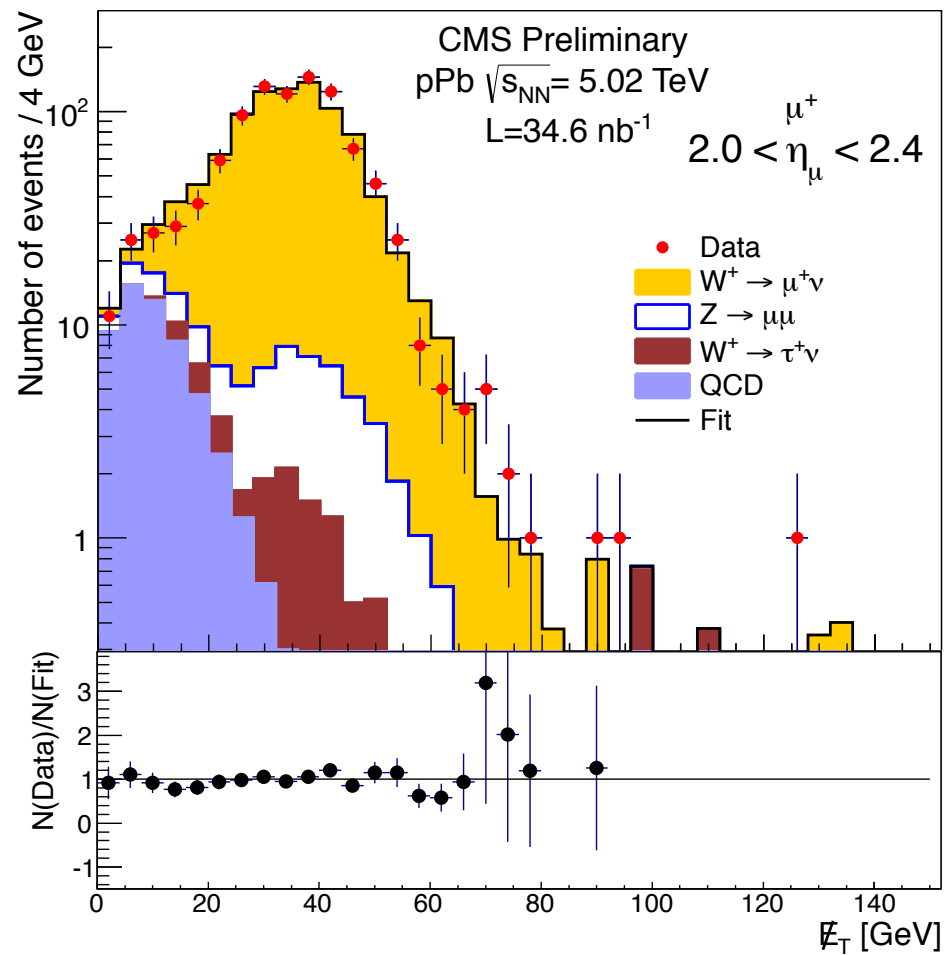
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Problems:

- How to estimate the QCD contamination?
- How to establish a systematic on this method?
- Tree cuts that depends on the MET resolution...

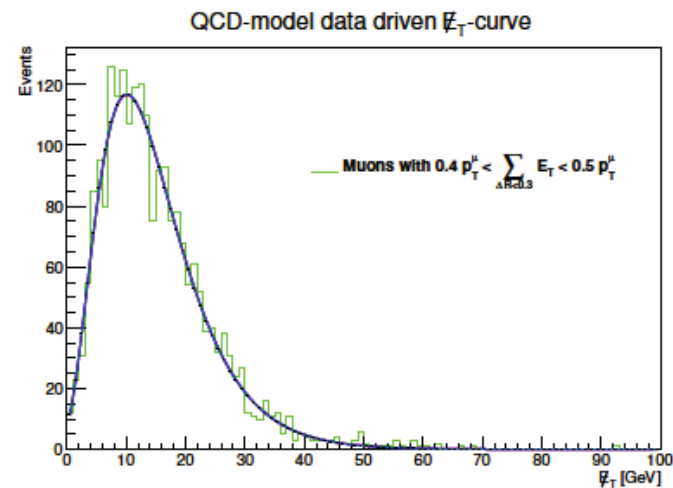
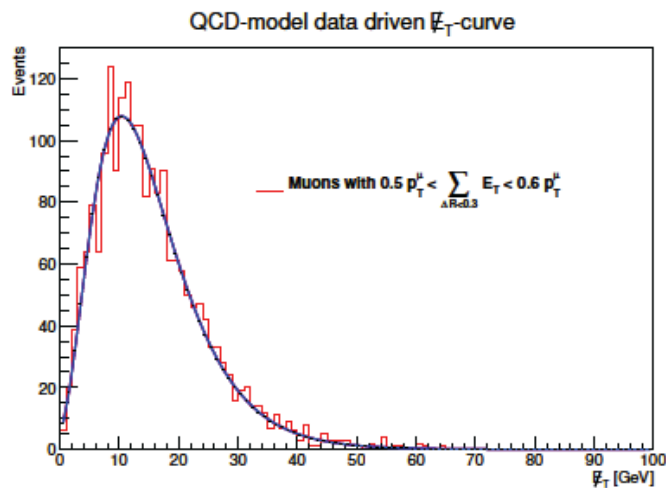
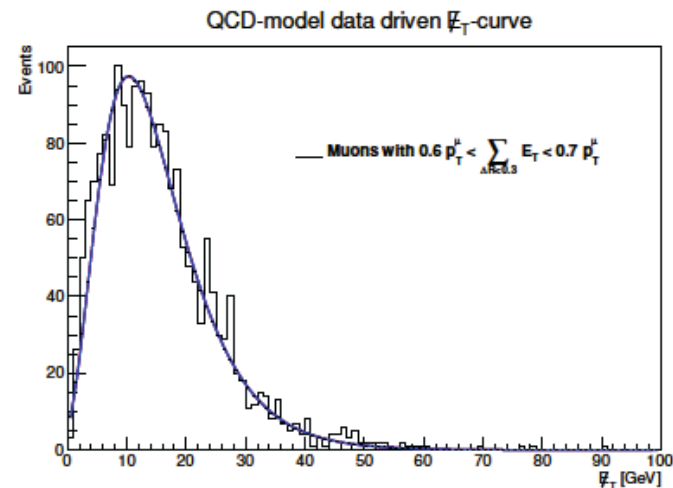
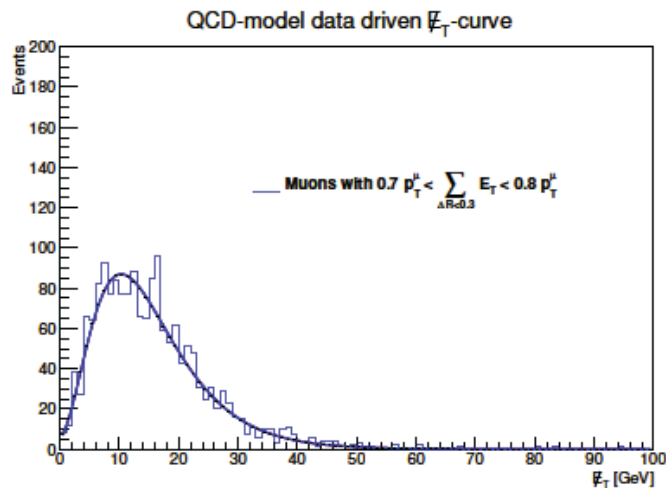
+ MET templates



+ MET fitting

Muon QCD

- Low QCD contamination and low non-isolated muons
- 4 sets of $(\beta, \alpha X_0)$ on only one η bin $[-2.4, 2.4]$



+ MET fitting

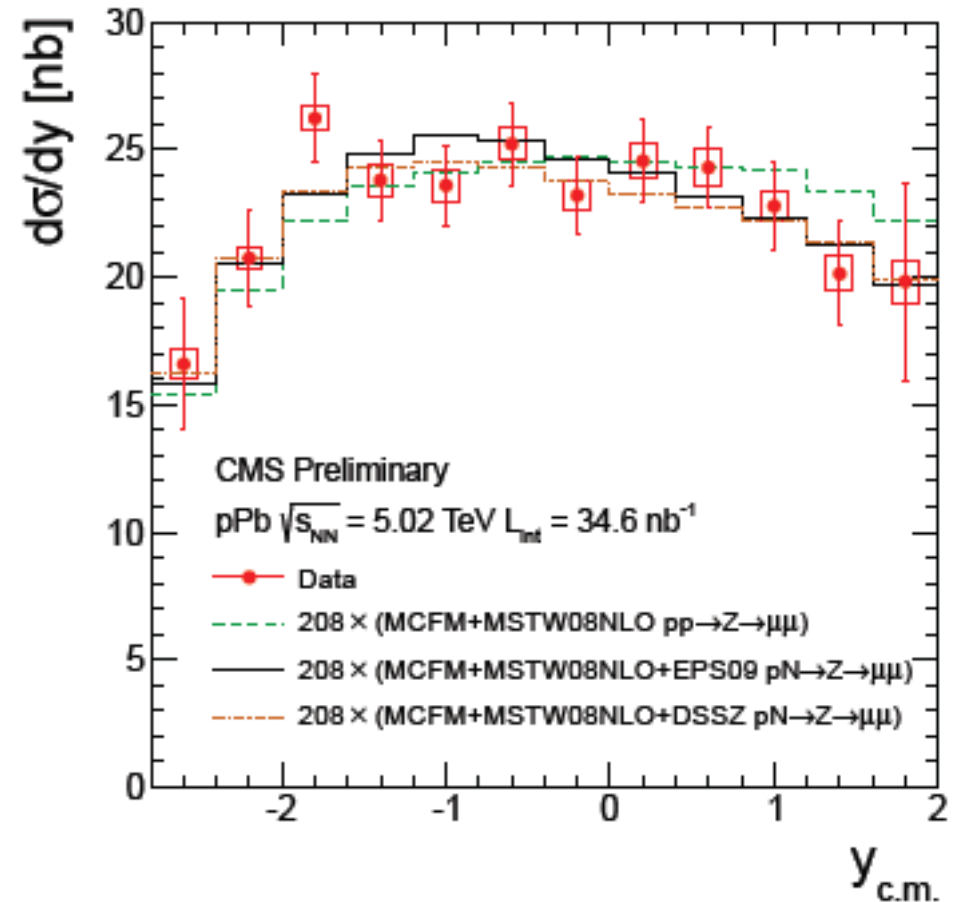
4) Systematic

QCD systematic:

- From 0.3 to 2 % with muon channel
- From 0.8 to 3.8 % with electron channel

Electroweak systematic:

- The W signal over background (Z and W τ) is fixed to the theoretical cross sections in pp collisions (k1 and k2)
- We need to take into account the effects on nPDFs not taken into account in k1 and k2
- Looking at Z analysis results we set a limit at 20% (conservative): between 1-3%.

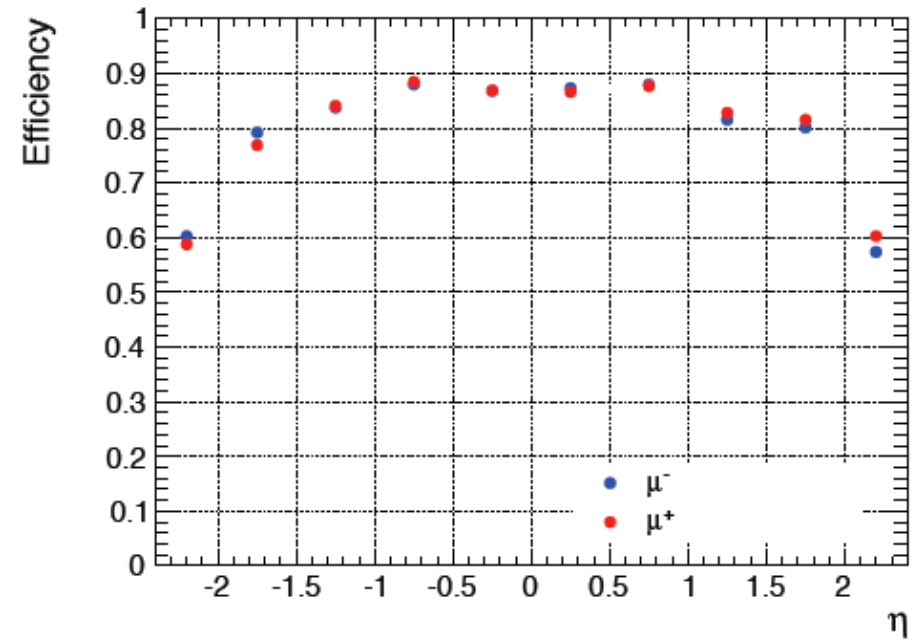
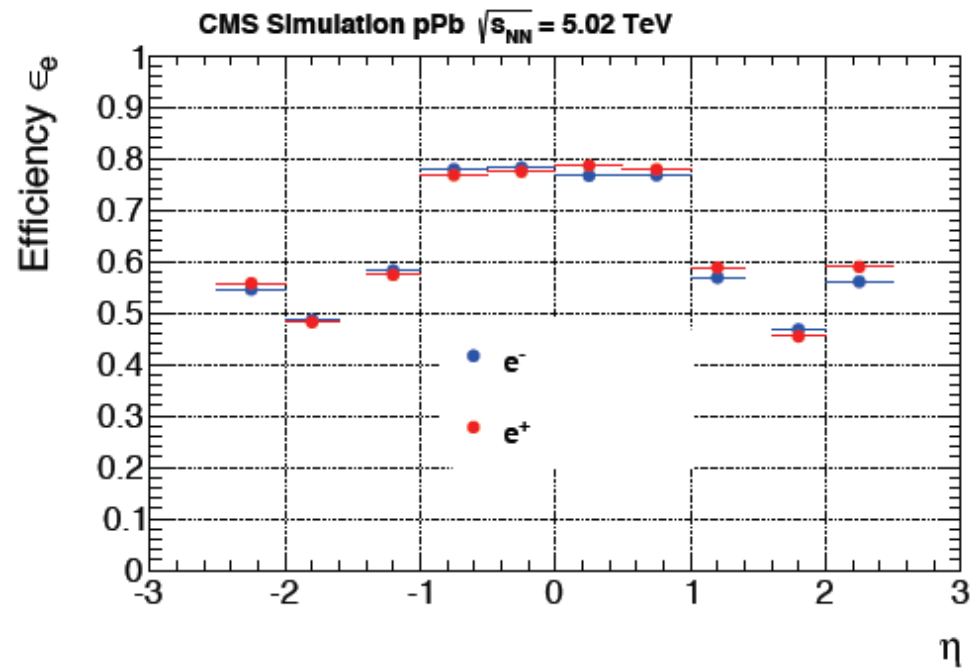


+ Analysis strategy

- Single lepton selection
- MET fitting
- **Single lepton efficiency**
- Scales Factors on single lepton Tag and Probe
- Cross section
- Asymmetry

$$N^{\text{corr}} = \frac{N^{\text{raw}}}{\epsilon_{\text{all}}^{\text{MC,W}} \times \text{SF}}$$

+ Single lepton efficiency



+ Analysis strategy

- Single lepton selection
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+ Scales Factors on single lepton Tag and Probe

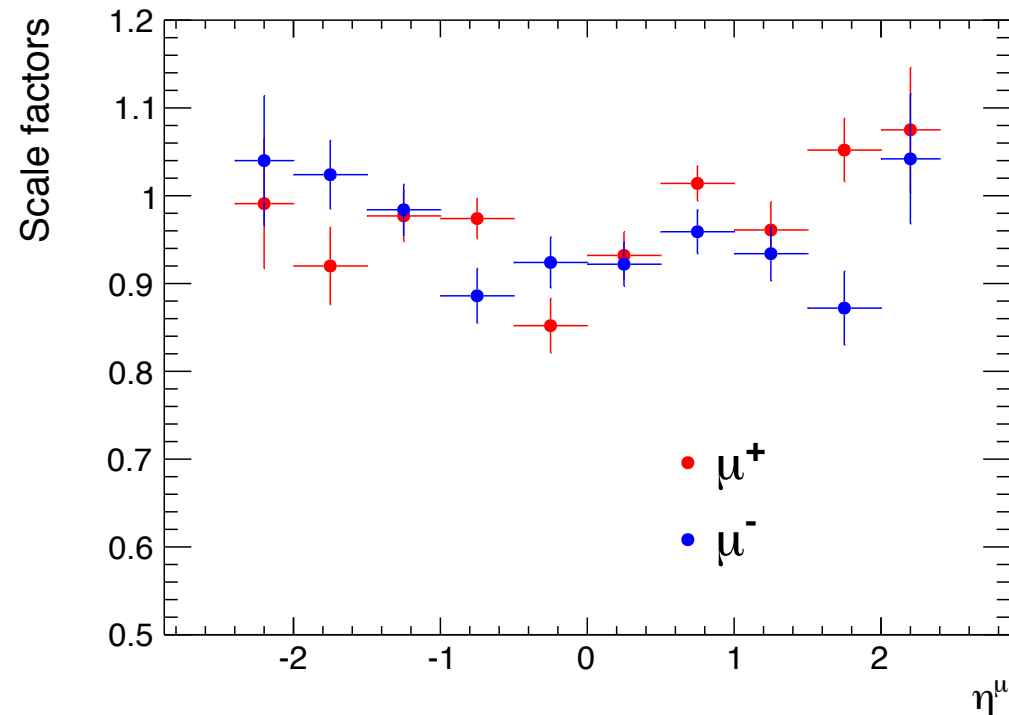
- To correct the MC efficiency from eventual bias we use the Tag and Probe method on Z events
- We define 3 sub-efficiency ϵ_{Reco} , $\epsilon_{\text{Id+Iso}}$, ϵ_{trig}
- SF are define for each η bin

$$\text{SF}(\eta) = \frac{(\epsilon_{\text{Reco}} * \epsilon_{\text{Id+Iso}} * \epsilon_{\text{trig}})_{\text{Data}}}{(\epsilon_{\text{Reco}} * \epsilon_{\text{Id+Iso}} * \epsilon_{\text{trig}})_{\text{MC}}}$$

+ Scales Factors on single lepton Tag and Probe

3) Muon SF

- The TnP measurements are dominated by the limited data statistics. The main systematic is then the sum of the tree TnP data errors: around 3%.
- In addition the difference between the Efficiency true and the product of the three MC TnP, is taken as systematic (MC closure test). At maximum 0.9%.



+ Analysis strategy

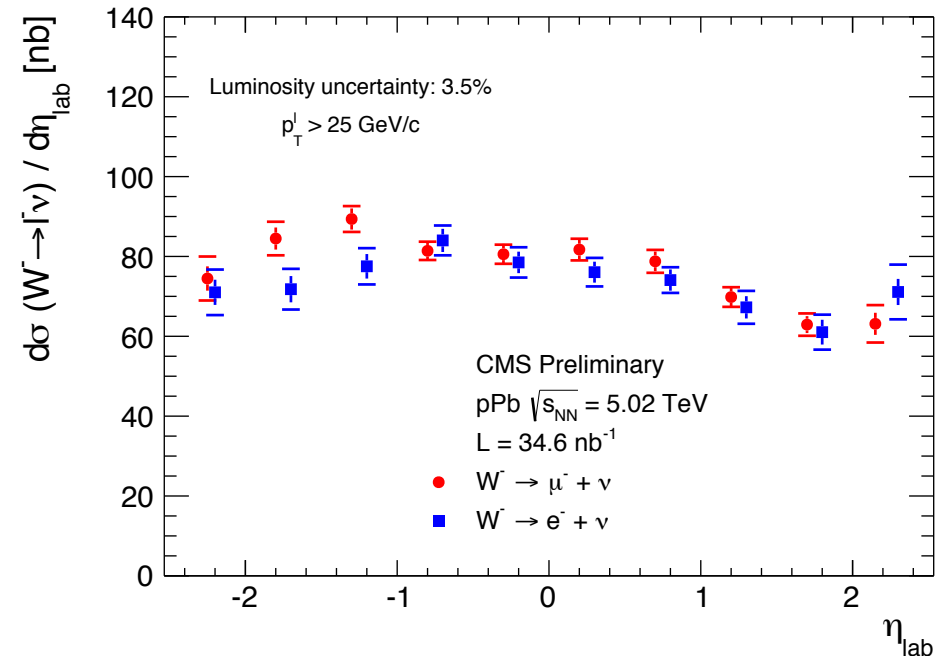
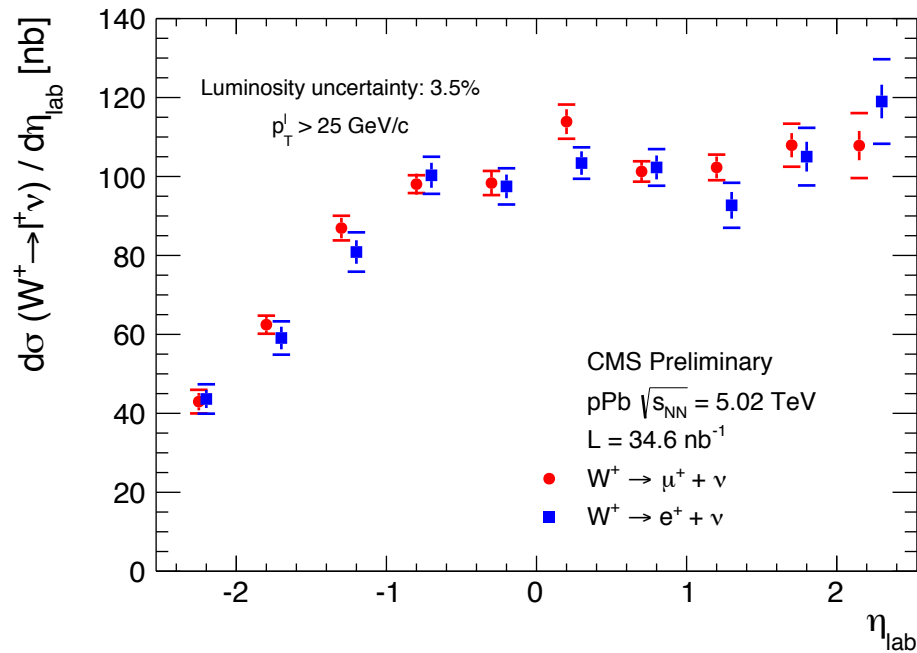
- Single lepton selection
- MET fitting
- Single lepton efficiency
- Scales Factors on single lepton Tag and Probe
- **Cross section**
- Asymmetry

$$N^{\text{corr}} = \frac{N^{\text{raw}}}{\epsilon_{\text{all}}^{\text{MC,W}} \times \text{SF}}$$

+ Cross section



HIN-13-007

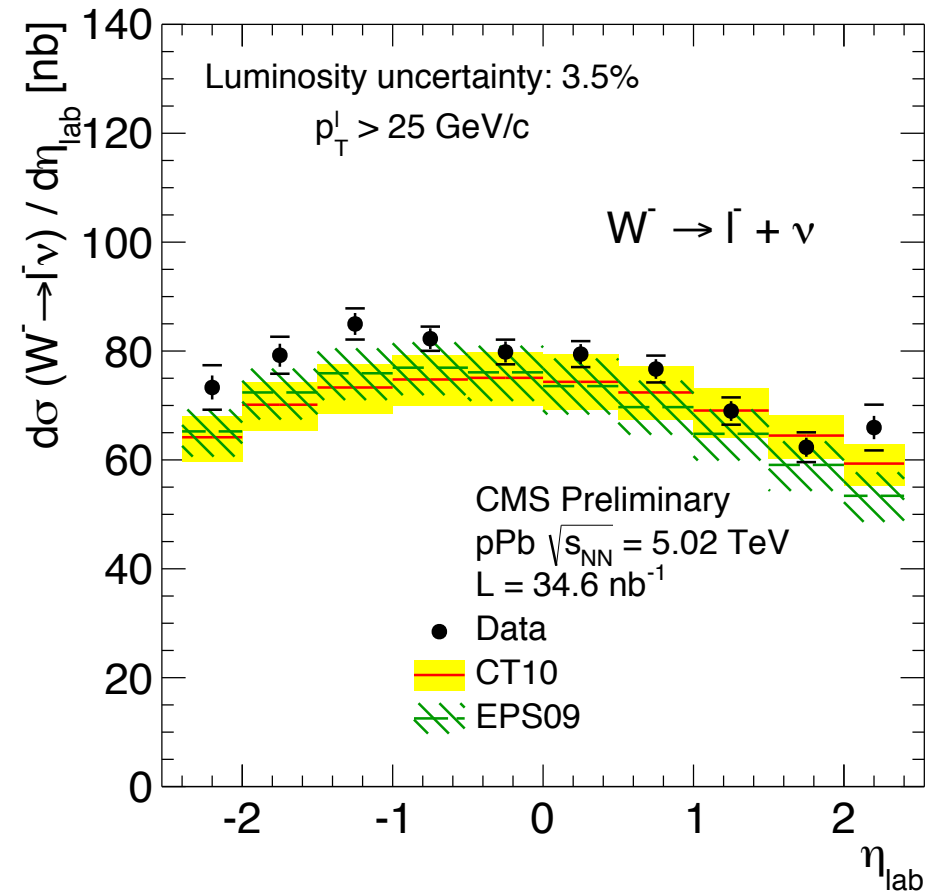
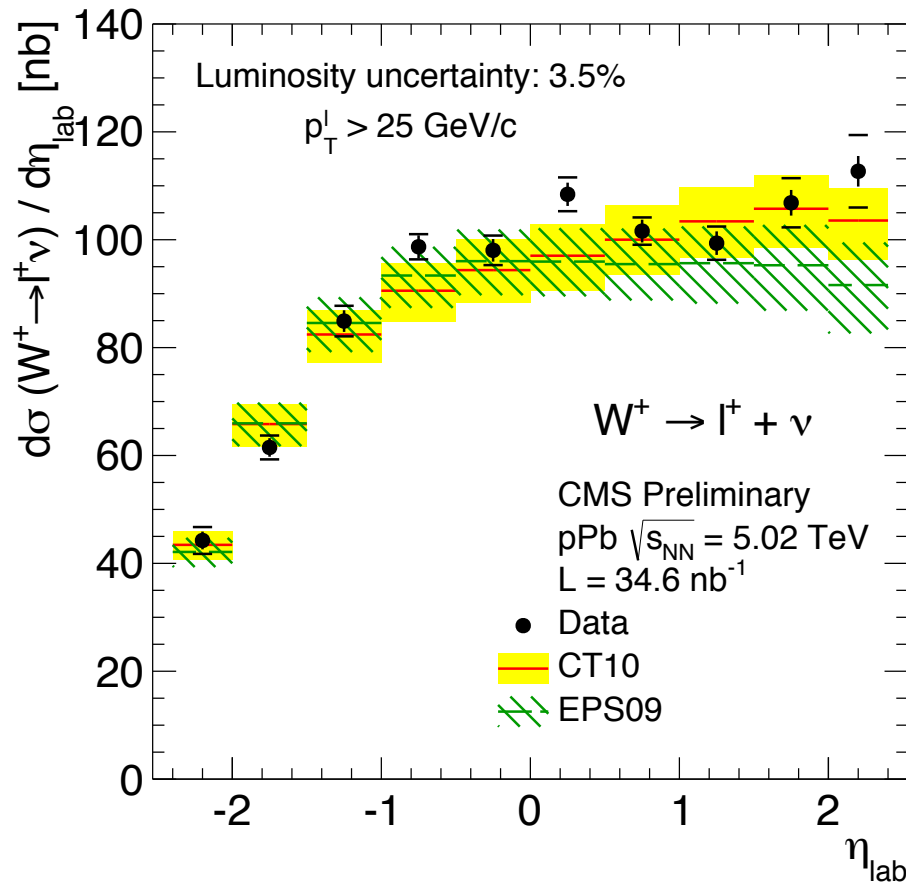


- Cross section tables are provided in the PAS
- W^+ are found mainly in the proton region fragmentation because of u valence quark coming from boosted proton

+ Cross section



HIN-13-007



- Results are combined according to the standard blue method
- There are no correlations bin to bin
- The two PDFs set have been provided by C.Salgados and H.

+ Analysis strategy

- Single lepton selection
- MET fitting
- Single lepton efficiency
- Scales Factors on single lepton Tag and Probe
- Cross section
- **Asymmetry : N^{corr} / L**

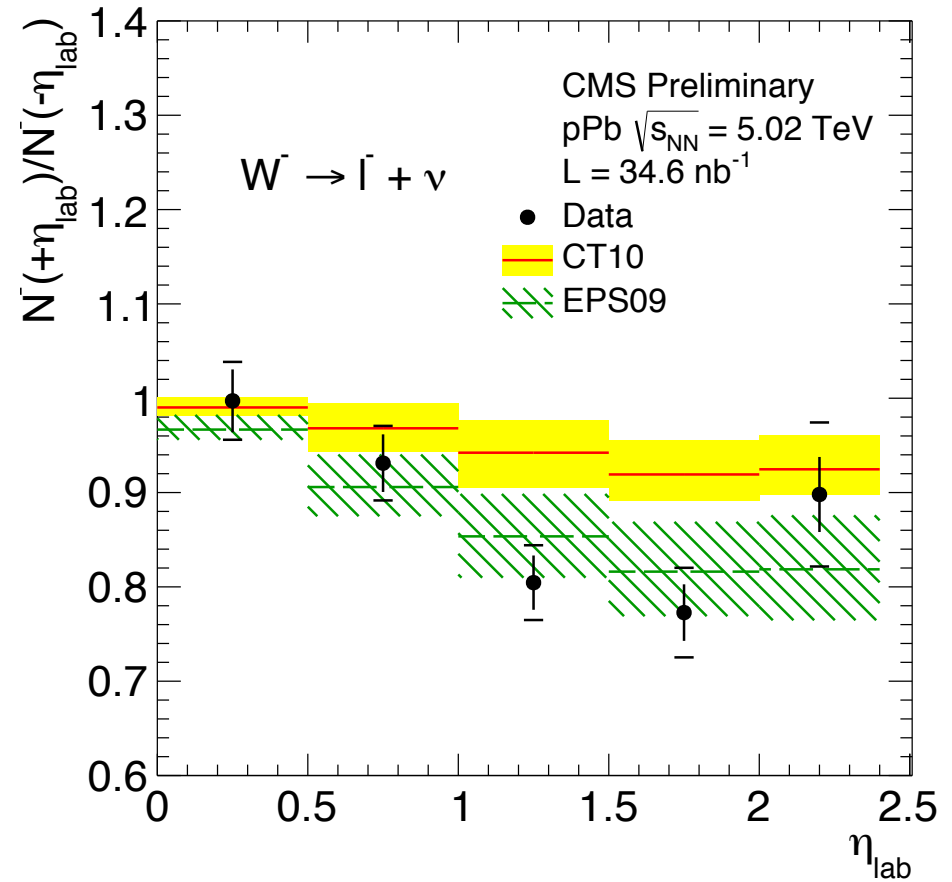
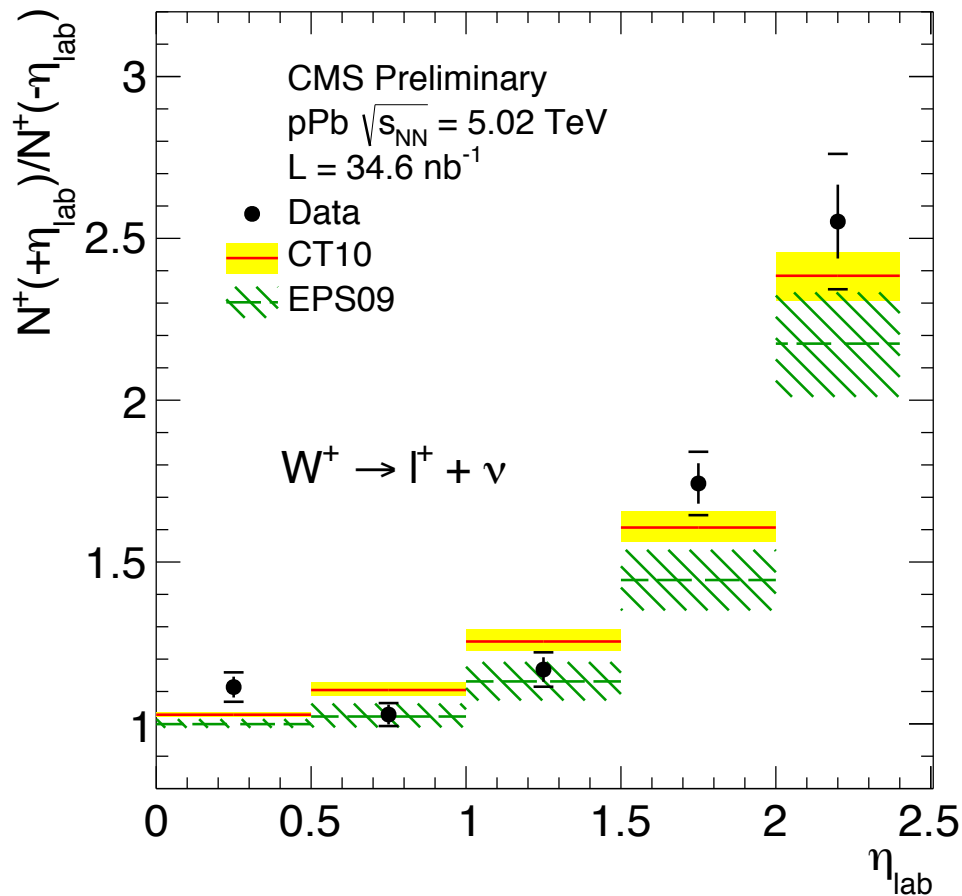
$$N^{\text{corr}} = \frac{N^{\text{raw}}}{\epsilon_{\text{all}}^{\text{MC,W}} \times \text{SF}}$$

+ Asymmetry



1) Forward/Backward W^+ and W^-

HIN-13-007



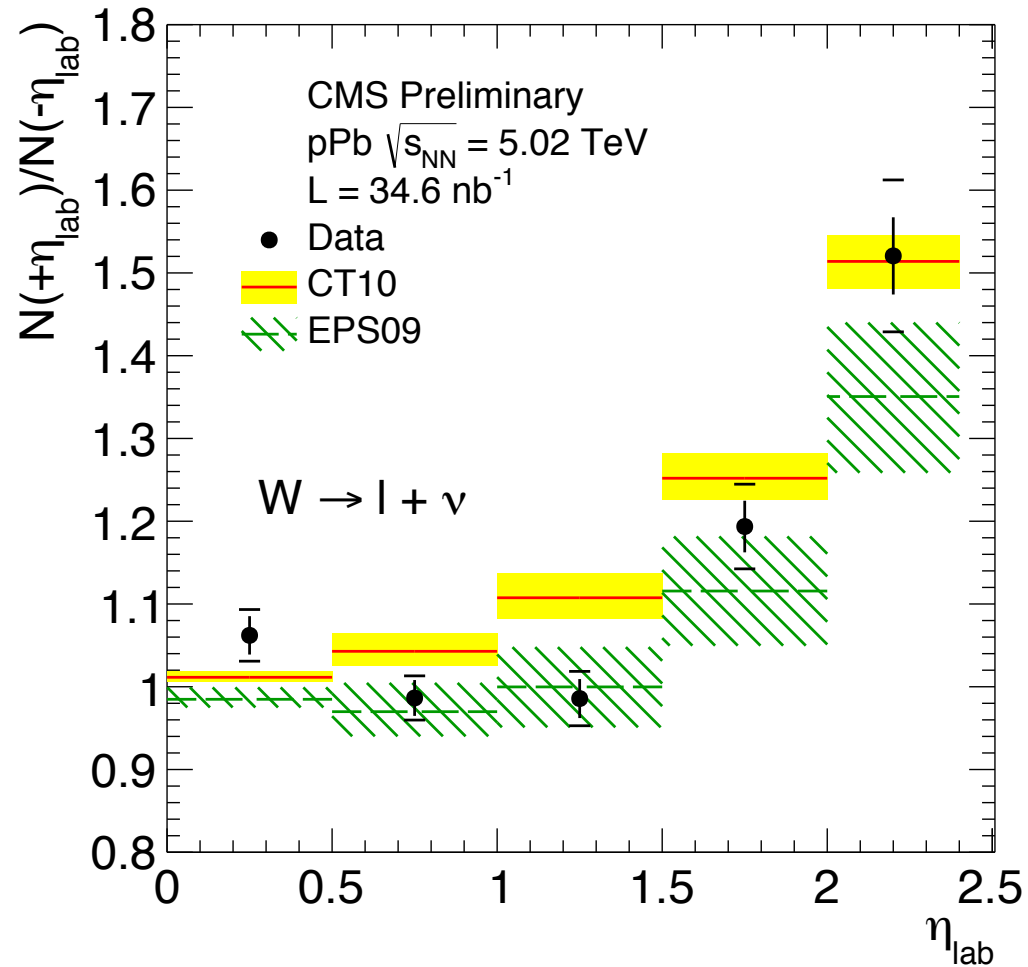
At High rapidity we probe the valence quark at high X (proton) compare to the sea quark at low X (lead)

+ Asymmetry



2) Forward/backward W

HIN-13-007



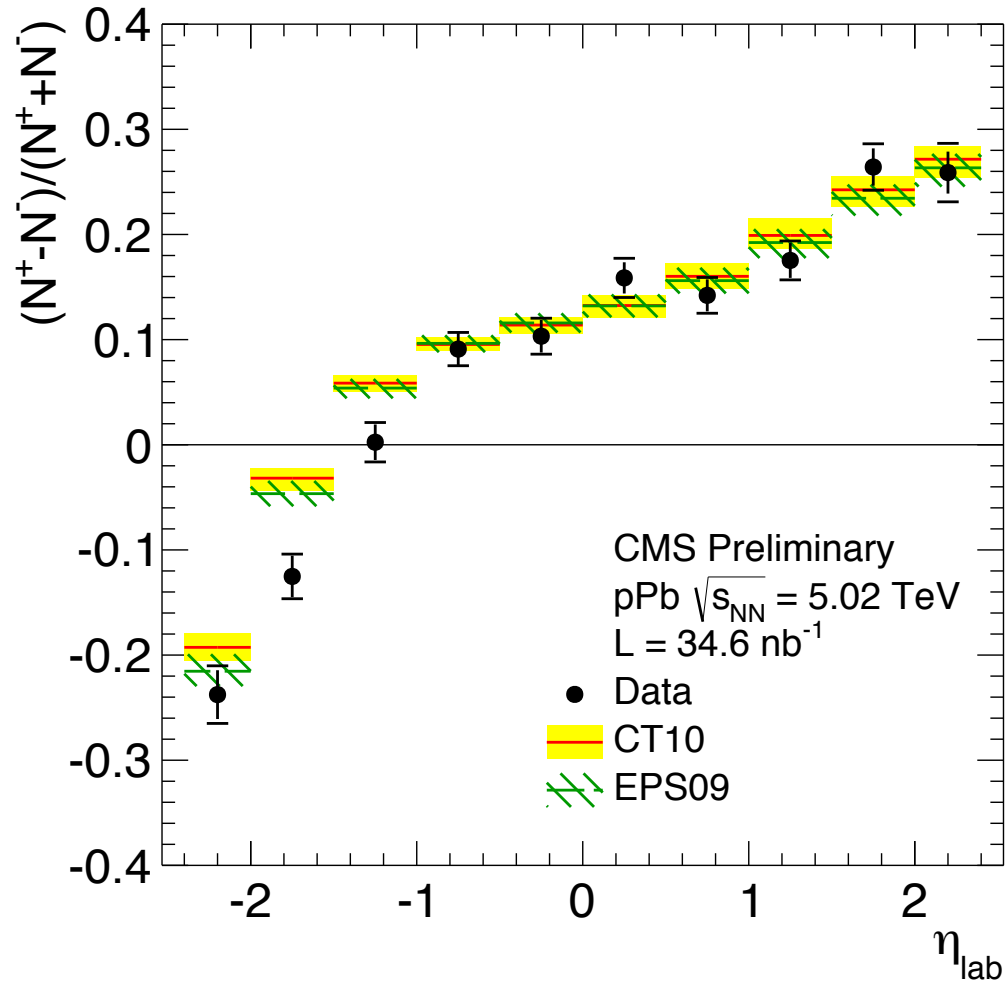
+ Asymmetry

1) Charge asymmetry



HIN-13-007

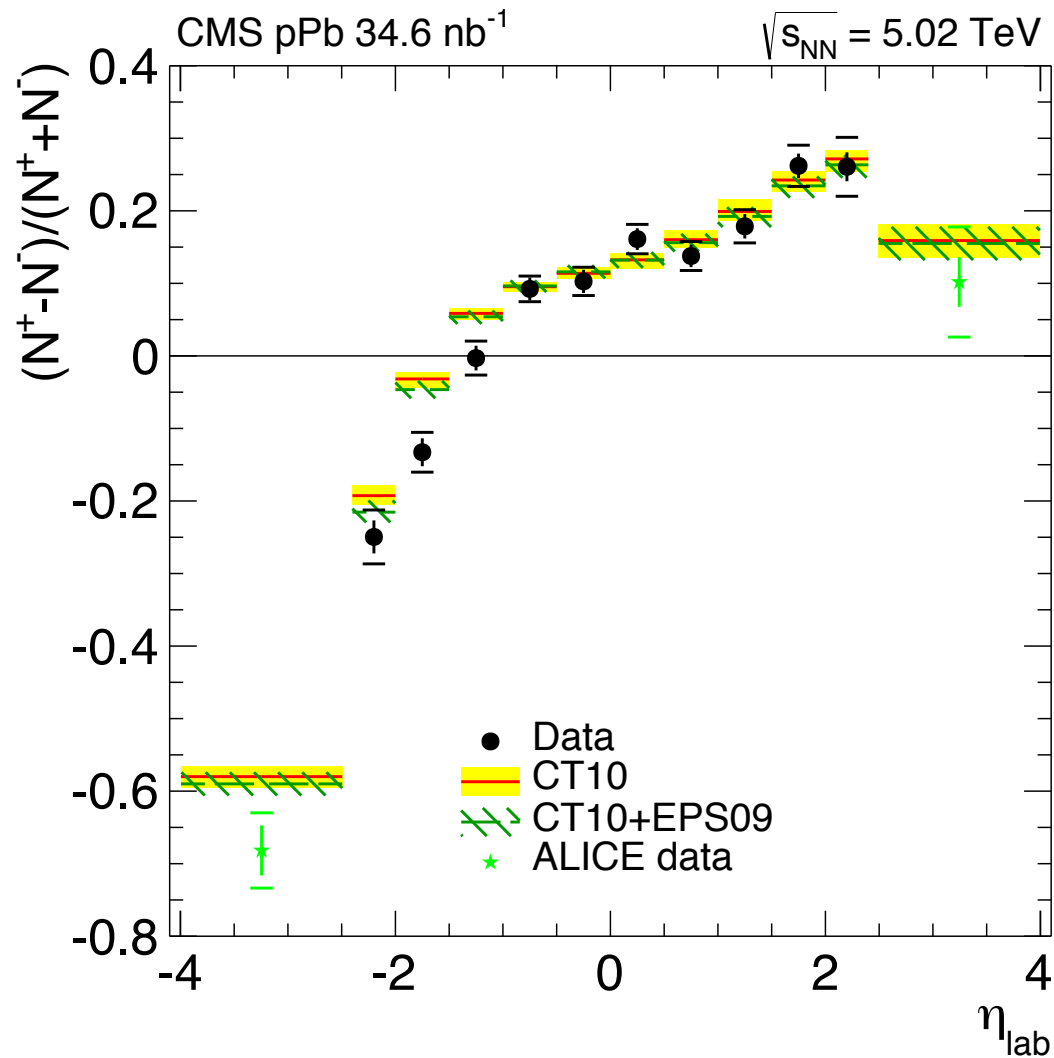
30



$$\frac{N(\mu^+) - N(\mu^-)}{N(\mu^+) + N(\mu^-)}$$

- Isospin effect: difference between the W- and W+ because of their sensitivity to quarks contents.
- The charge asymmetry is supposed to be independent of the nuclear modifications of partons densities, excepted if $R(u) \neq R(d)$

+ Bonus: ALICE + CMS points



+ Back up

+ MET fitting

4) Systematic

QCD:

- Change function: asymmetric lognormal, 4 parameters
- Ideally we would need an other set of data... Production of toy experiment

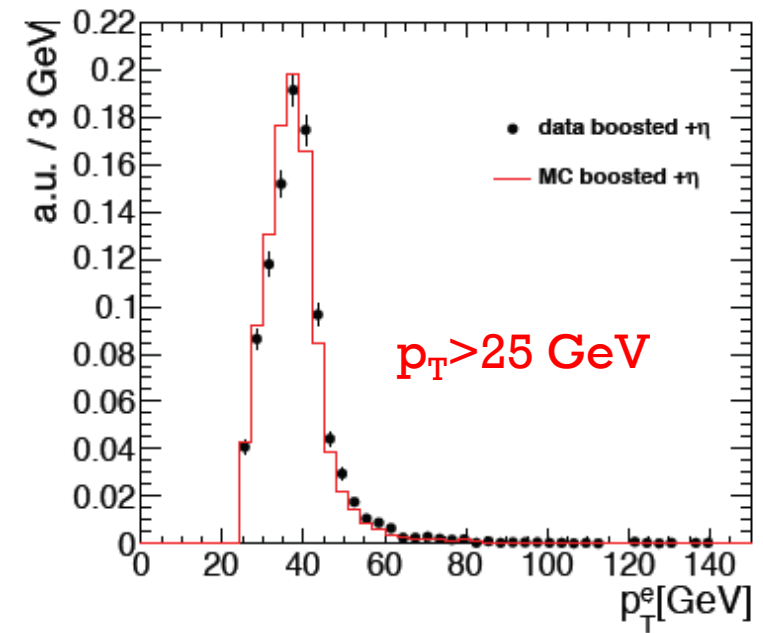
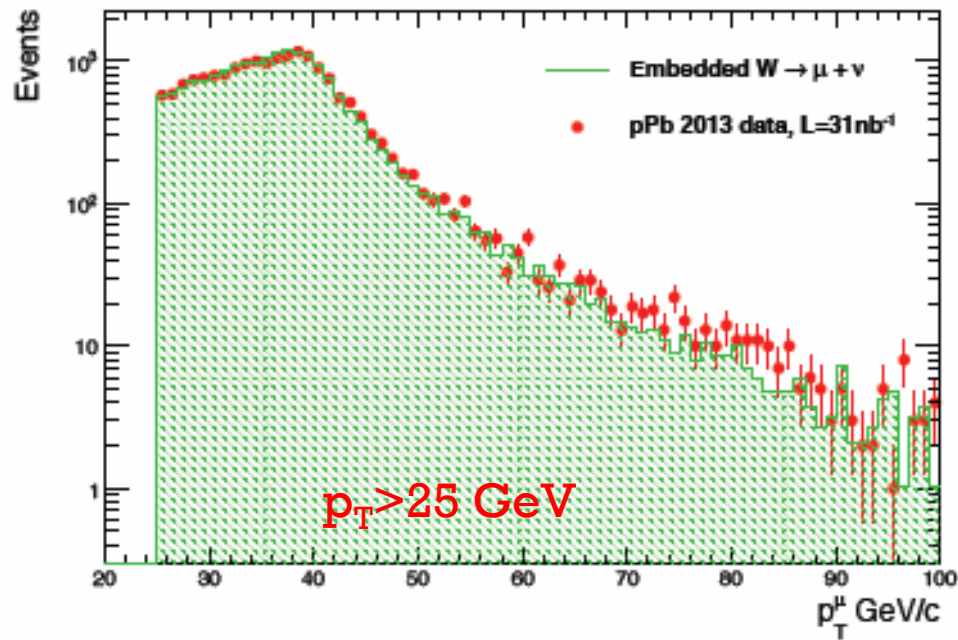
- 1) Production of a data sample where the QCD varies according to the **asymmetric lognormal**

$$f(x) = C / \sqrt{2\pi\sigma_L\sigma_R} \exp(-0.5((\log(x) - x_0)/\sigma_L)^2) \quad \text{if } \log(x) < x_0$$
$$f(x) = C / \sqrt{2\pi\sigma_L\sigma_R} \exp(-0.5((\log(x) - x_0)/\sigma_R)^2) \quad \text{if } \log(x) > x_0$$

- 2) The rest of the pseudo data: $W + Z + W \tau$ id estimated by subtracting the obtained QCD to the selected data.
- 3) **1000 pseudo experiment per lepton charge and η bin**
- 4) Each pseudo data set is fitted with the analysis function
- 5) **The yields difference is taken as systematic** coming from QCD modeling

+ Single lepton selection

2) High p_T lepton

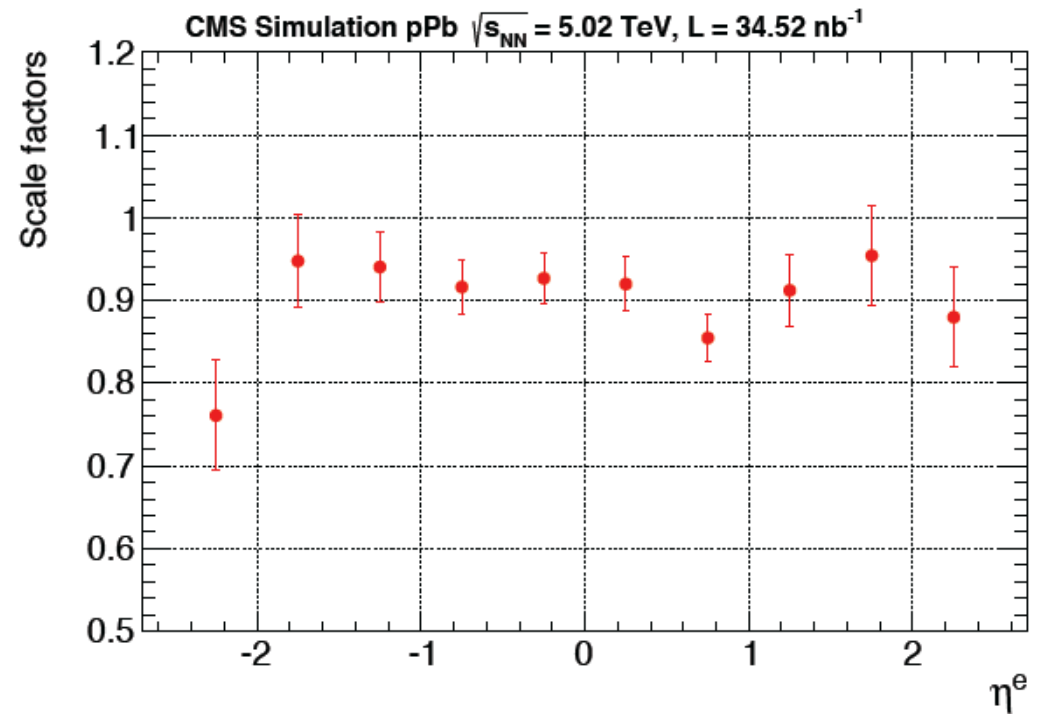


+ Scales Factors on single lepton Tag and Probe

35

3) Electron total SF

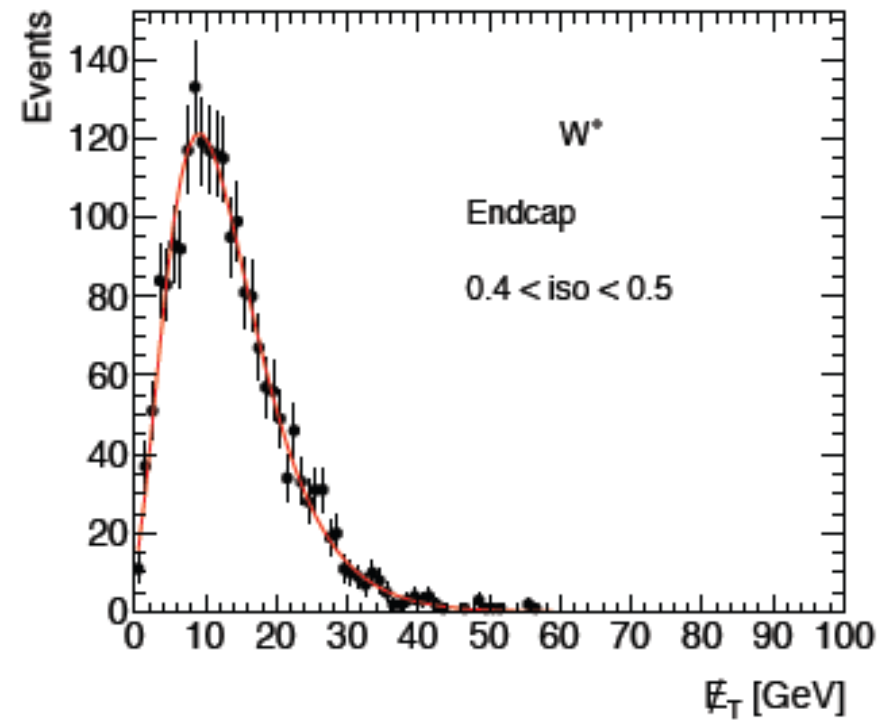
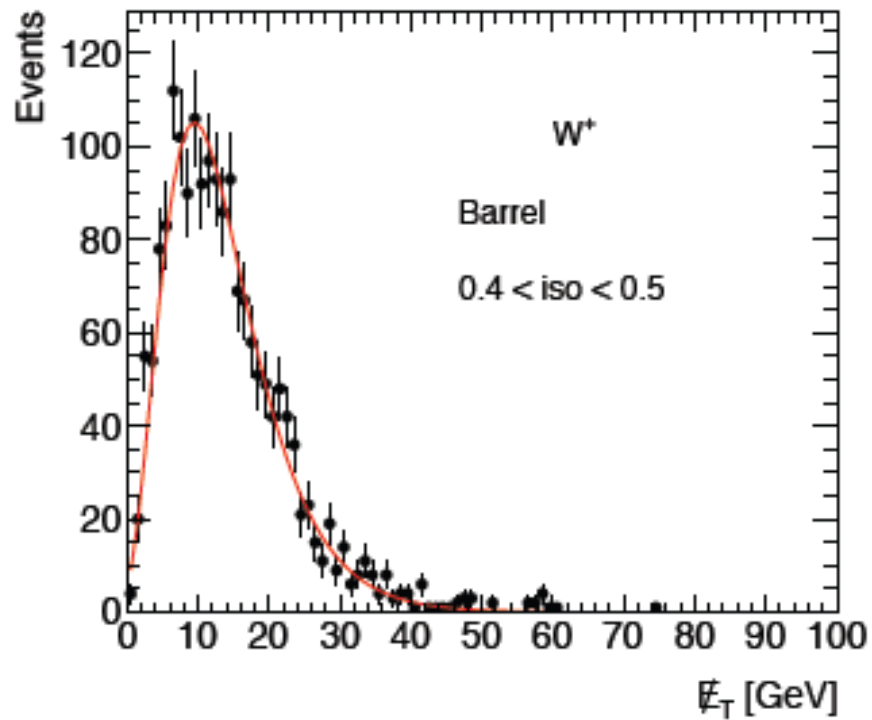
- The TnP measurements are dominated by the limited data statistics. The main systematic is then the sum of the tree TnP data errors
- In addition the difference between counting method and fit is taken as systematic
- For charge asymmetry these contributions will cancel since they are the same for both charge in a given η bin



+ MET fitting

2) Electron QCD

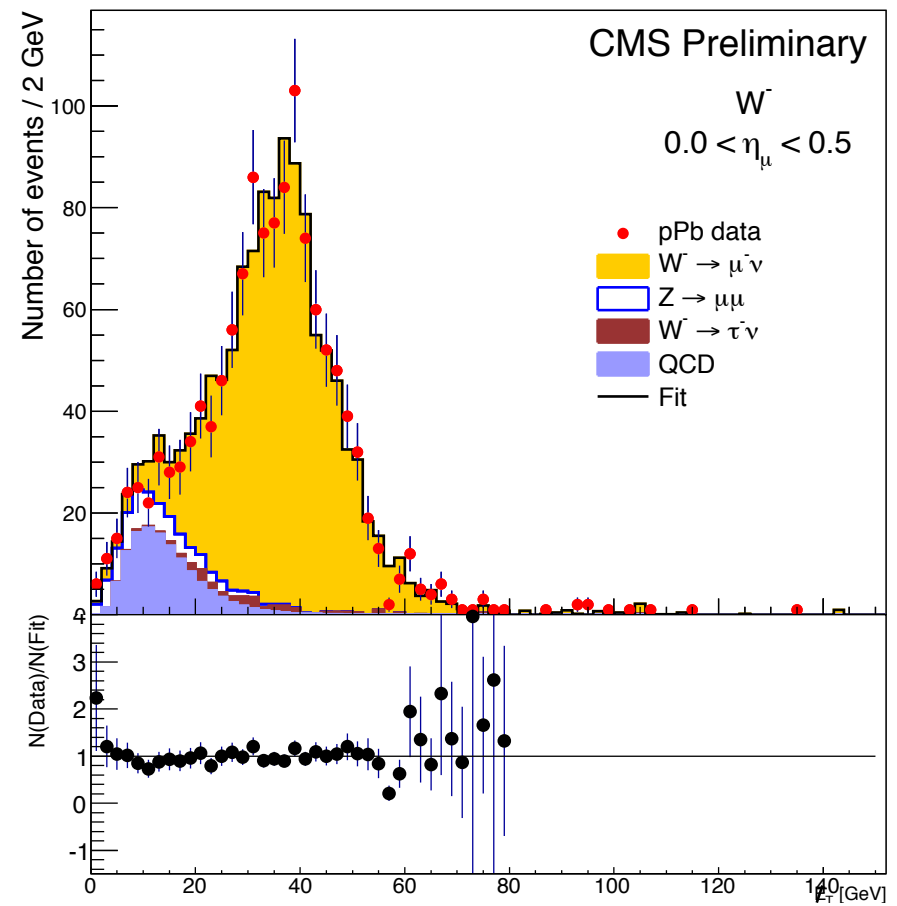
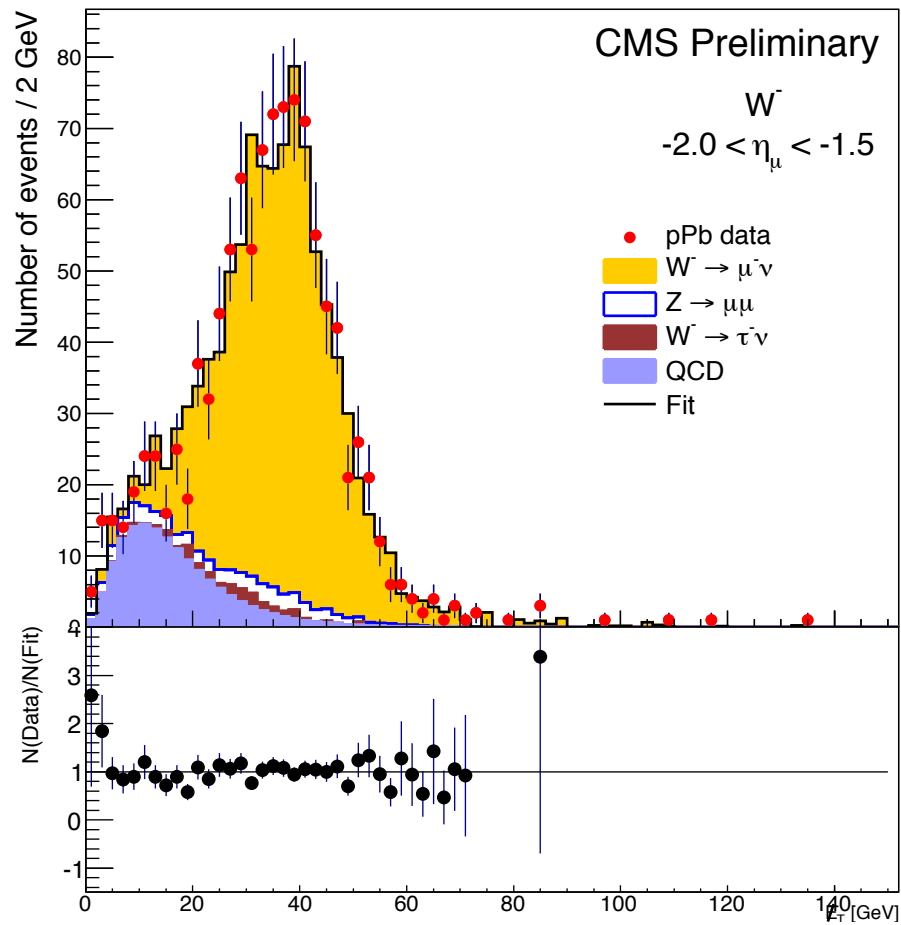
- Large QCD contamination in end caps
- 4 sets of $(\beta, \alpha X_0)$ for barrel region and 4 others for the end caps region
- We also looked per charge



+ MET fitting

3) Templates

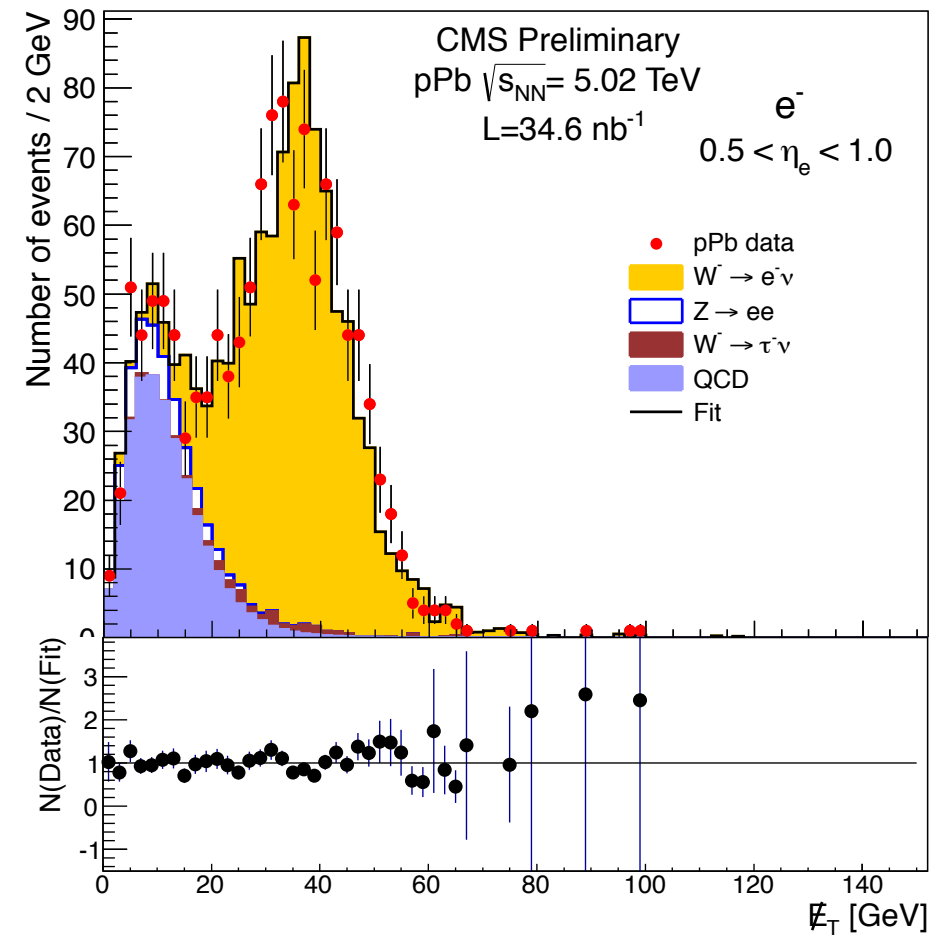
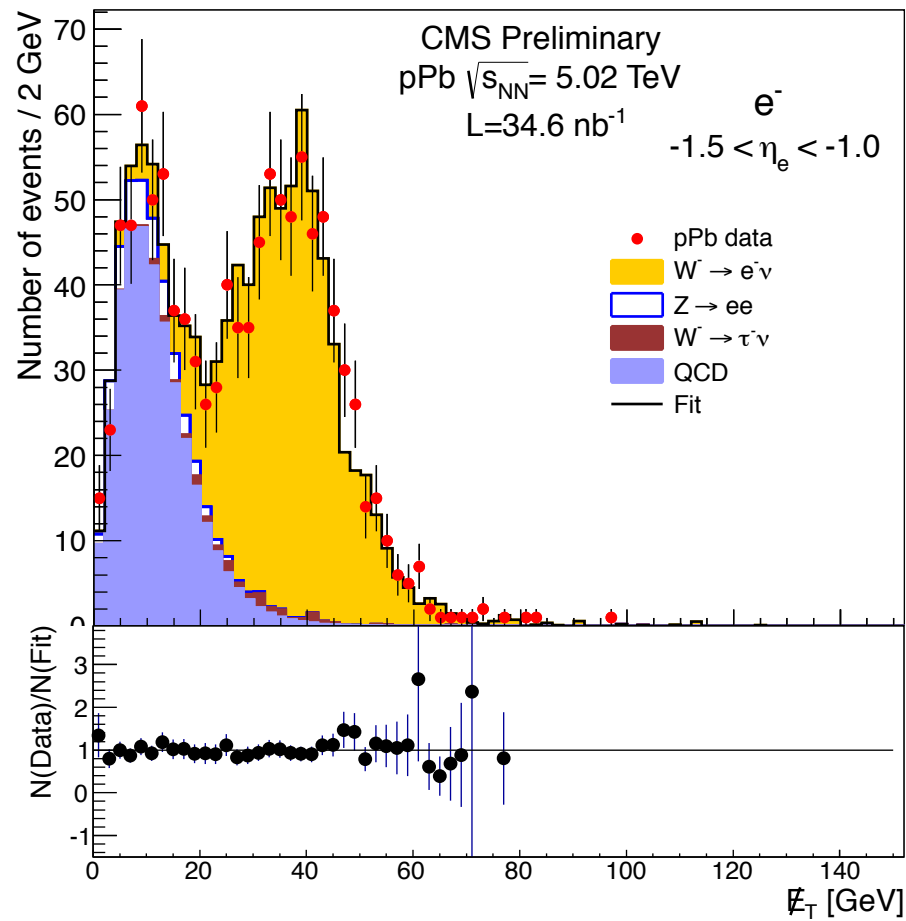
Muon



+ MET fitting

3) Templates

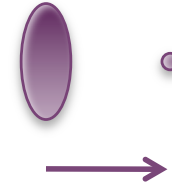
Electron



+ Data taking condition

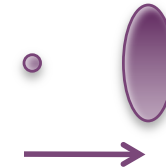
- **One Pbp run:** lead going toward $+\eta$ at $L_{\text{int}} = 20.7 \text{ nb}^{-1}$

- [210 498] to [211 256]



- **One pPb run:** proton going toward $+\eta$ at $L_{\text{int}} = 14 \text{ nb}^{-1}$

- [211 313] to [211 631]



- For the W analysis the results need to be in the same boost configuration

- **Second run is flipped in η** (the one with less statistic)

+ MET fitting

1) W, Z and W_τ

■ Shape:

- These two embedded simulation are passed over the analysis cut and the selected lepton fill MET histograms: T_W , T_Z and $T_{W\tau}$

■ Normalization:

- Correction for acceptance and efficiency
 - $\text{Acc}_{W_1} * \text{Eff}_{W_1} = N_{\text{selected}}/N_{\text{generated}} : A_W * \epsilon_W$
 - $\text{Acc}_Z * \text{Eff}_Z(\text{mimicking } W) = N_{\text{selected}}/N_{\text{generated}} : A_Z * \epsilon_Z, A_{W\tau} * \epsilon_{W\tau}$
- The W is normalized to the W theoretical cross section σ_W
- The Z and W_τ contamination are normalized to the W signal
 - The ratio of the theoretical Z and W cross sections $K_1 = \sigma_Z / \sigma_W$
 - The W_τ branching ratio K_2

$$N(E_T) = \{ \sigma_W \times [A_W \cdot \epsilon_W \cdot T_W(E_T) + K_1 \cdot A_Z \cdot \epsilon_Z \cdot T_Z(E_T) + K_2 \cdot A_{W\tau} \cdot \epsilon_{W\tau} \cdot T_{W\tau}(E_T) + \mathcal{F}_{\text{QCD}}(E_T)] \} \times \mathcal{L}$$

+ MET fitting

2) QCD

- The function used to fit the MET has 3 parameters: β , α and X_0

$$F(x) = (x+X_0)^\alpha \exp(-\beta \sqrt{x+X_0})$$

- β will be fixed thanks to data driven method. α and X_0 let free but cross checked with values from data driven method.
- These 3 parameters are extracted from data:
 - The shape of MET for events having non isolated lepton is fitted with this function
 - 4 different area of no isolation are chosen: [0.4,0.5], [0.5,0.6], [0.6,0.7] and [0.7,0.8], each one giving one (β, α, X_0) set
 - The 4 (β, α, X_0) set are plotted versus the isolation and the fit extrapolation is taken as $\beta_s, \alpha_s, X_{0s}$ in the signal region (iso<0.1)

$$N(E_T) = \{\sigma_W \times [\mathcal{A}_W \cdot \epsilon_W \cdot \mathcal{T}_W(E_T) + K1 \cdot \mathcal{A}_Z \cdot \epsilon_Z \cdot \mathcal{T}_Z(E_T) + K2 \cdot \mathcal{A}_{W\tau} \cdot \epsilon_{W\tau} \cdot \mathcal{T}_{W\tau}(E_T) + \mathcal{F}_{QCD}(E_T)]\} \times \mathcal{L}$$