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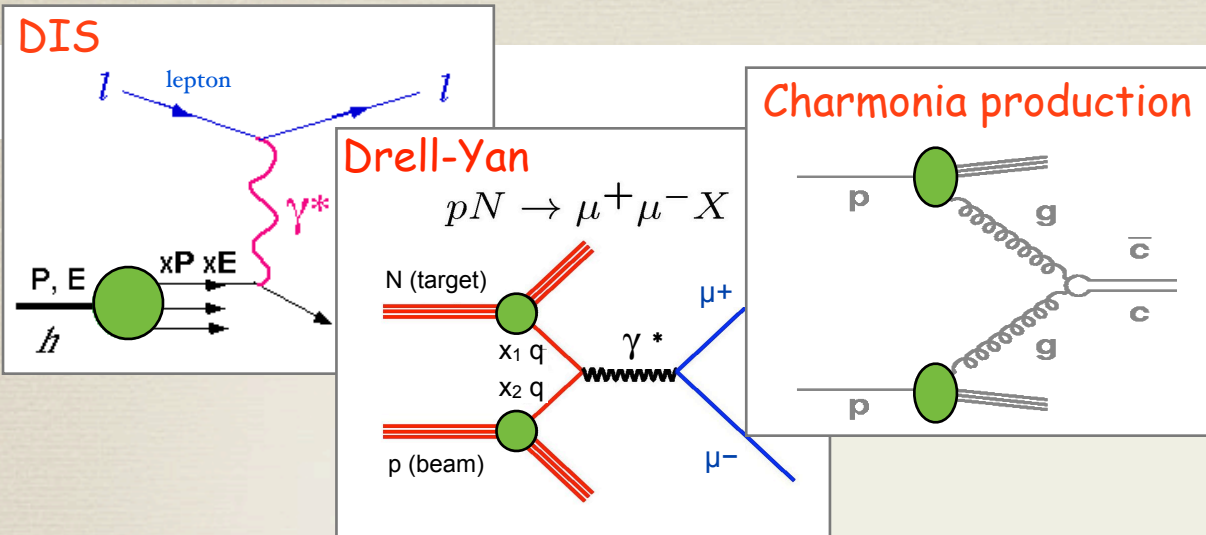
In collaboration with E. G. Ferreira,  
F. Fleuret, J. P. Lansberg & N. Matagne

# CNM EFFECTS ON QUARKONIUM PRODUCTION : FROM RHIC TO LHC

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CEA (Saclay) IRFU/SPhN

Rencontres Ions Lourds, July 26th, 2011

# Shadowing : a cold nuclear matter effect



Processes used to probe :

nucleon struct. f.

$$F_2 = \sum_i e_i^2 \cdot x f_i(x, Q^2)$$

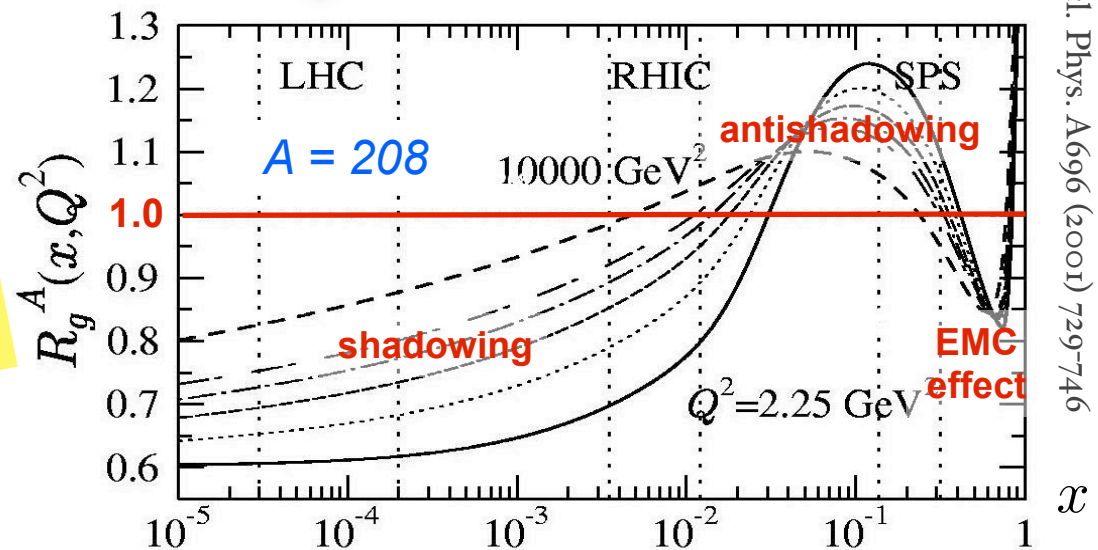
with  $f_i(x, Q^2) = \text{PDF}$  and  $i = q, \bar{q}, g$

nuclear struct. f. per nucleon

$$R_g^A = \frac{\text{g PDF} \in \text{bound nucleon}}{\text{g PDF} \in \text{free nucleon}}$$

(Anti-)shadowing :

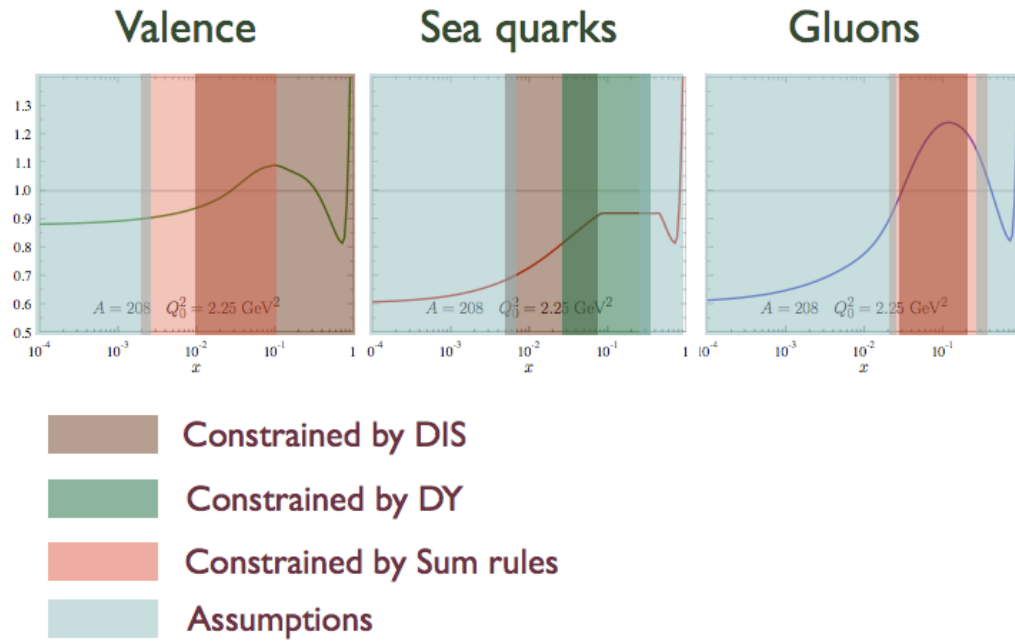
- initial-state effect "calibrated" in  $d(p)+A$
- (enhances) decreases  $\sigma^{pA}$  wrt  $\langle N_{\text{coll}} \rangle \sigma^{pp}$



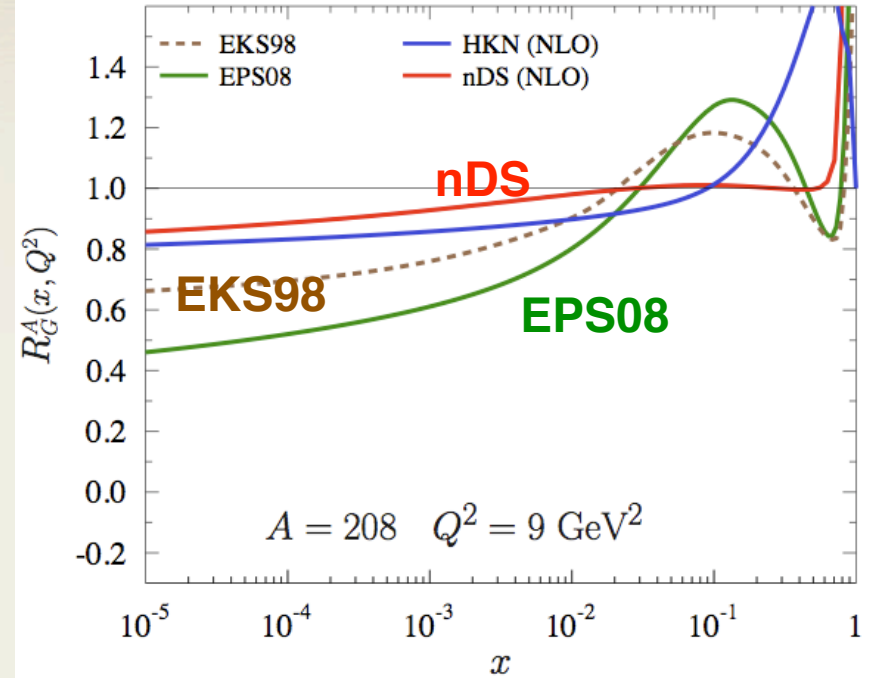
Nucl. Phys. A696 (2001) 729-746

# (Some) nPDFs available on the market

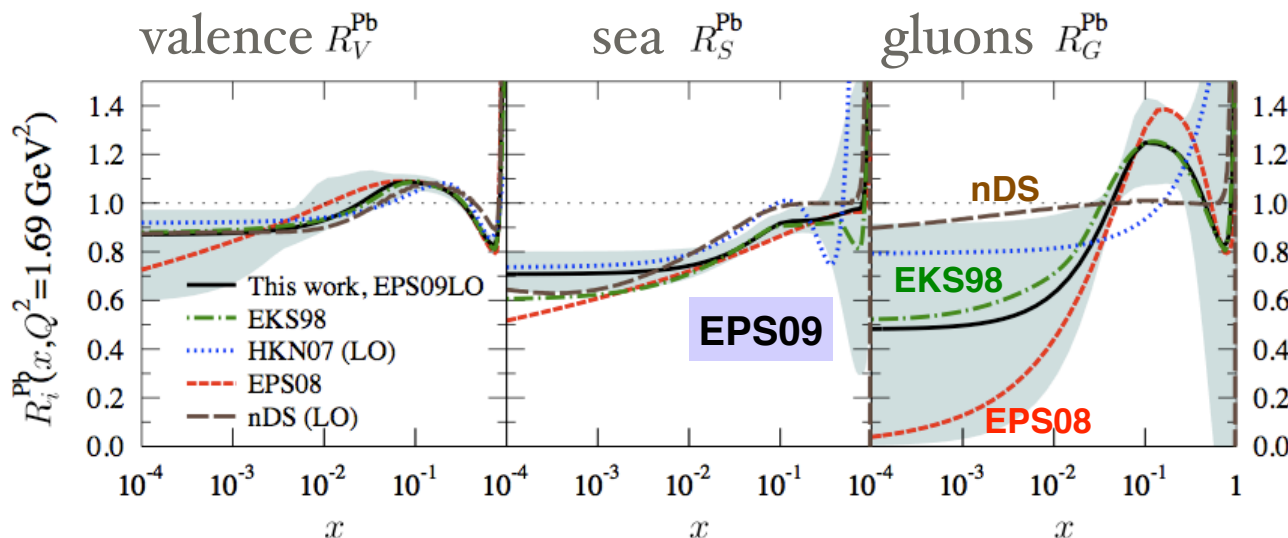
## Approximate ranges and constraints in EKS98



## Ratios for gluons and Pb nuclei



C. Salgado, ECT Trento July 2008



nDS(g), EKS98 and EPS08 span the uncertainty on nPDF

Eskola, Paukkunen, Salgado, JHEP 0904:065 (2009)

# Shadowing computation


- in A+B: quarkonia production cross-section modified by a shadowing correction factor :

$$\mathcal{F}(R_g^A(x_1, Q^2)) \times \mathcal{F}(R_g^B(x_2, Q^2))$$

- 4-mom conservation relates  $(x_1, x_2)$  to quarkonia  $(y, p_T)$
- production models (CEM, NRQCD, CSM ...) in p+p gives quarkonia thanks to various processes, each with:
  - a given phase-space in  $(x_1, x_2, y, p_T)$
  - a given weight (differential cross-section) for each point in this phase-space


☑ different production models may result in quite different shadowings

# How prod. models can differ ?

 intrinsic scheme  
 $2 \rightarrow 1$  process

$$g + g \rightarrow c\bar{c} \text{ or } b\bar{b}$$


$$x_{1,2} = \frac{m}{\sqrt{s_{NN}}} e^{\pm y}$$

 Handy : unequivocal  
correspondence

$$(x_1, x_2) \Leftrightarrow (y, p_T)$$


 Quarkonia  $p_T$  comes from  
initial partons

 e.g. CEM LO

 extrinsic scheme  
 $2 \rightarrow 2$  process

$$g + g \rightarrow \{J/\psi, \Upsilon\} + g$$

more degrees of freedom in  
the kinematics :

 several  $(x_1, x_2) \Leftarrow (y, p_T)$

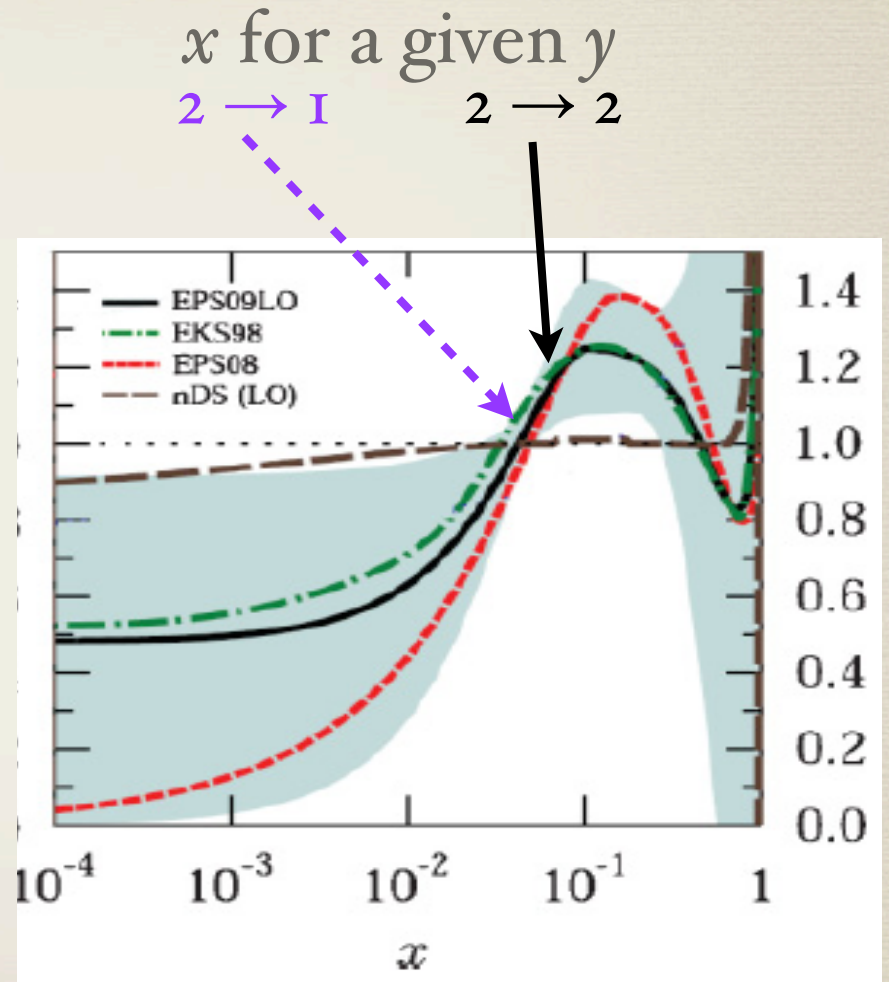
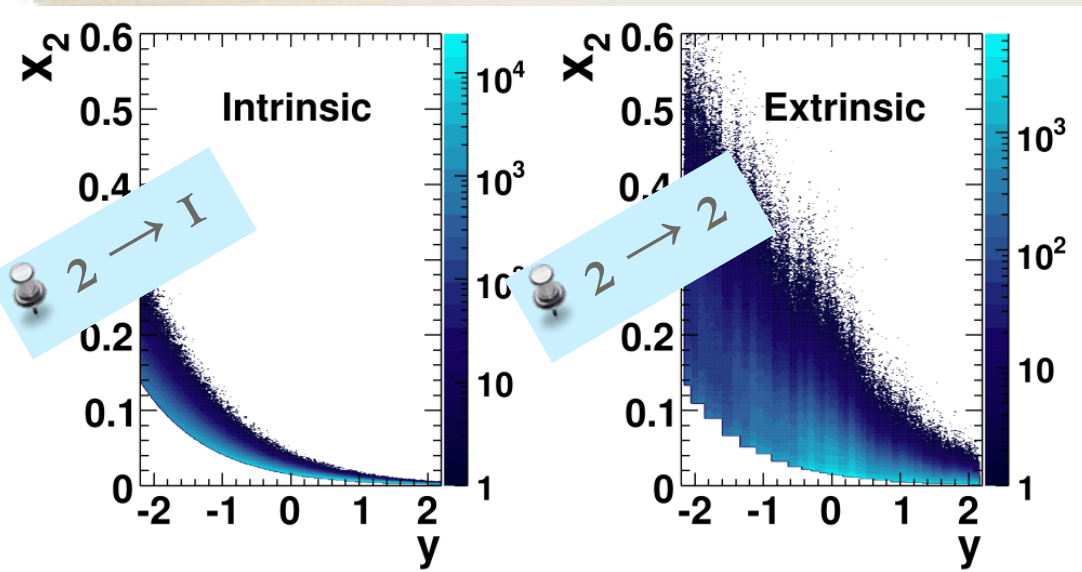
$$y, p_T, x_1 \implies x_2 = \frac{x_1 m_T \sqrt{s} e^{-y} - M^2}{\sqrt{s} (\sqrt{s} x_1 - m_T e^y)}$$

 Quarkonia  $p_T$  is balanced  
by the outgoing gluon

 e.g. CSM LO, COM LO

 Use reasonably good models in p+p to  
compute CNM effects in p+A, A+A

# CNM effects at RHIC : $J/\psi$ in dAu



for a given  $y$ ,  $\langle x \rangle$  is larger in the  $2 \rightarrow 2$  process

# CNM effects at RHIC : $J/\psi$ in dAu

E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A. R.  
 PRC 81 (2010) 064911

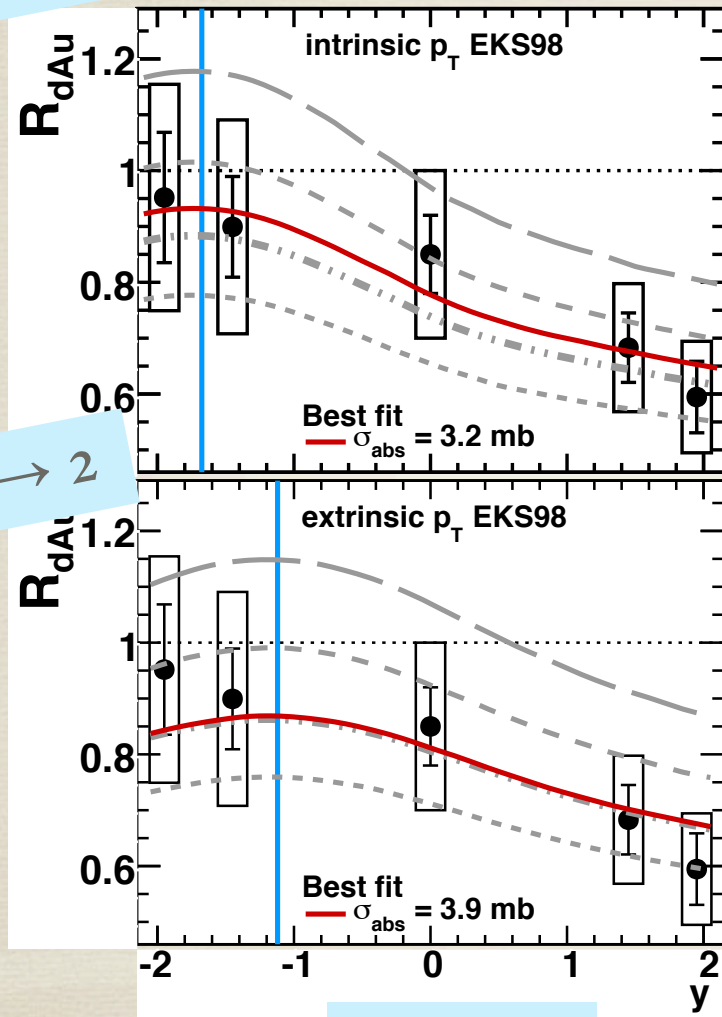
$R_{dAu}$  vs rapidity

$2 \rightarrow 1$

$g + g \rightarrow c\bar{c}$

$2 \rightarrow 2$

$g + g \rightarrow J/\psi + g$



EKS98

for a given  $y$ ,  $x$  is larger in the  $2 \rightarrow 2$  process

larger absorption cross-section needed

antishadowing peak shifted to larger  $y$

systematic effects seen for all nPDFs

# CNM effects at RHIC : $J/\psi$ in dAu

E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A. R.  
 PRC 81 (2010) 064911

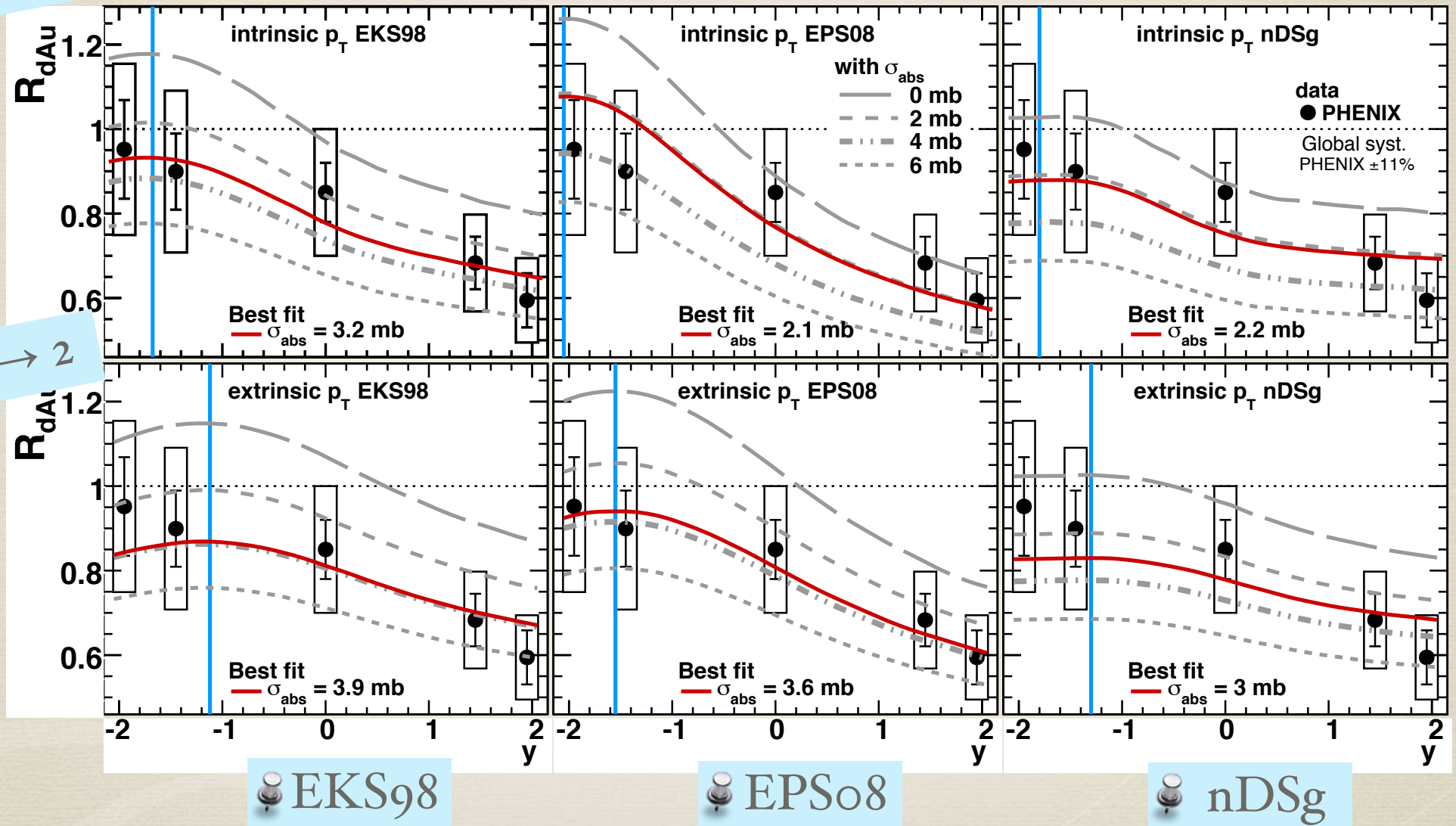
$R_{dAu}$  vs rapidity

$2 \rightarrow 1$

$g + g \rightarrow c\bar{c}$

$2 \rightarrow 2$

$g + g \rightarrow J/\psi + g$



EKS98

EPS08

nDSg



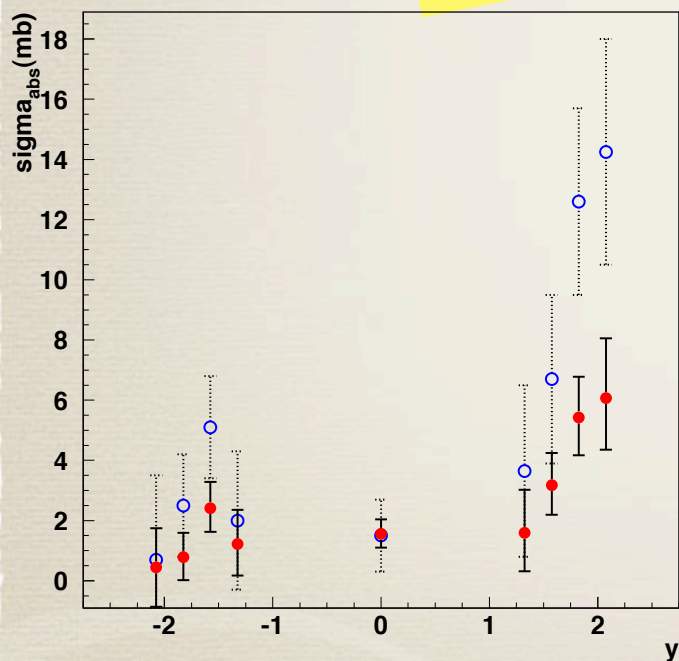
# CEM vs s-channel cut CSM as prod. model :

## $\sigma_{\text{abs}}(y)$ from Rcp in dAu @ 200 GeV

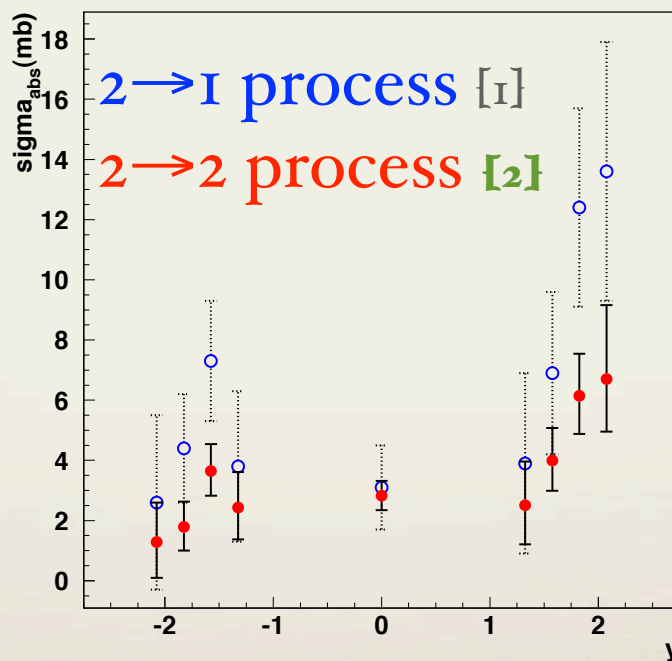
**$\sigma_{\text{abs}}(y)$  much flatter for the  $2 \rightarrow 2$  process**

[1] A. D. Frawley, INT, Seattle USA, June 2009

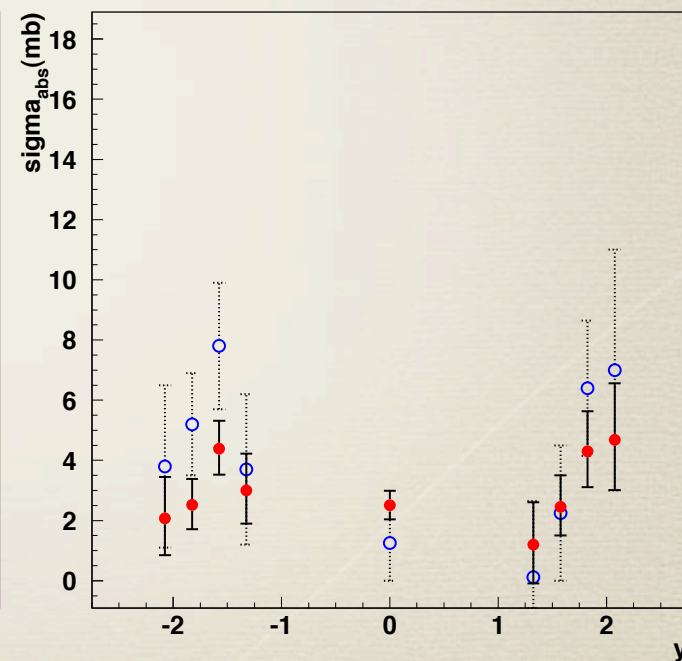
[2] E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A. R., PRC 81 (2010) 064911



nDSg



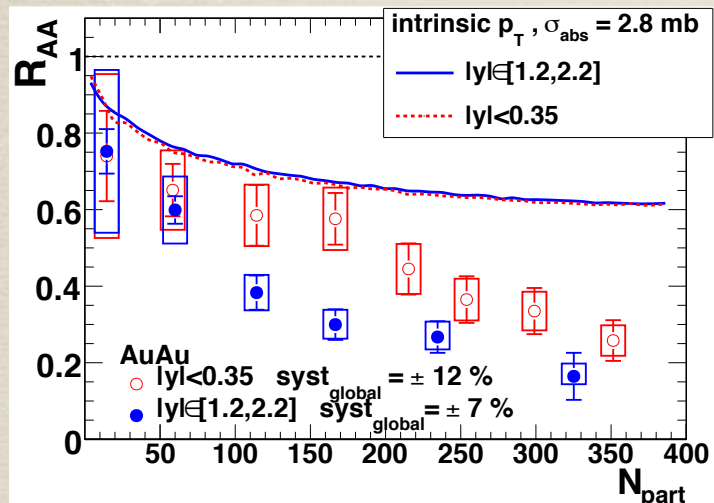
EKS98



EPS08

$2 \rightarrow 1$  process : mid-y & fwd-y

200 GeV

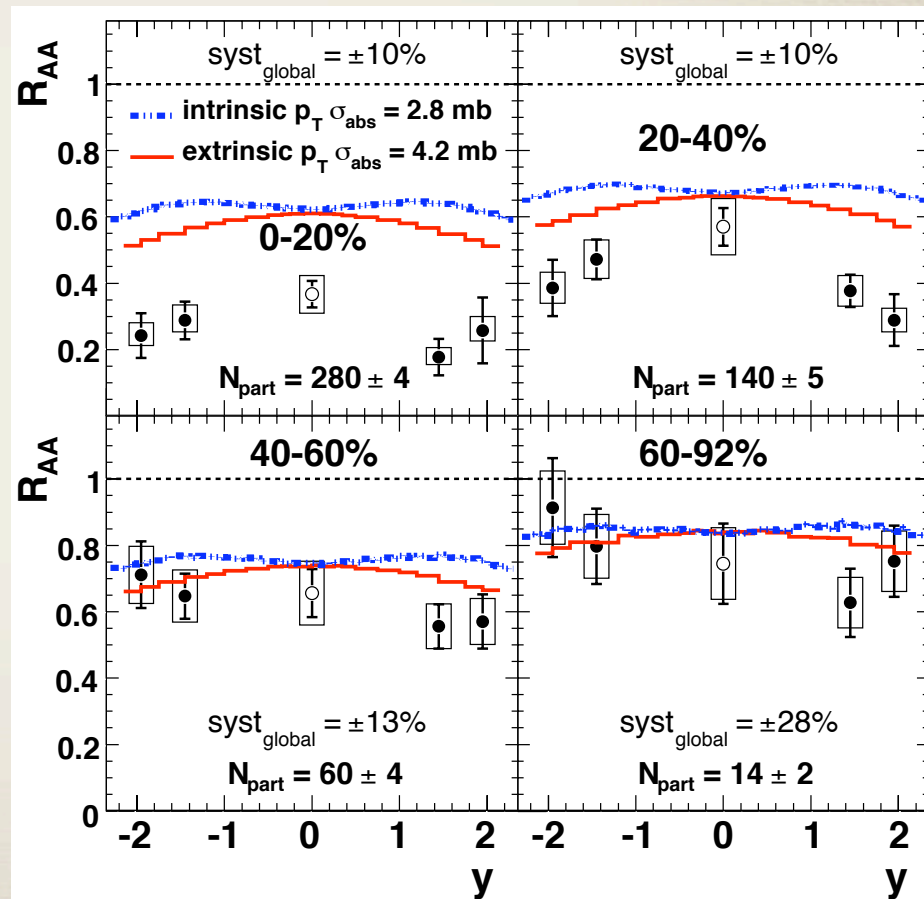
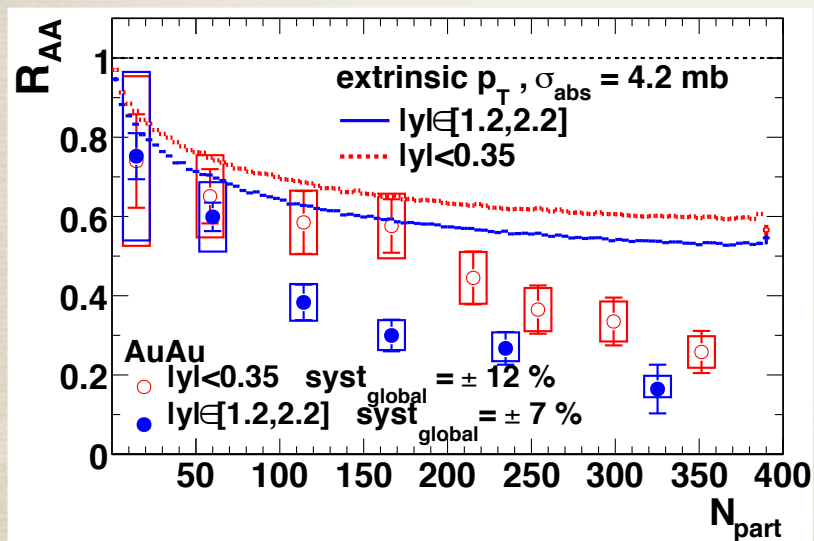


$2 \rightarrow 2$  process :  
less amount of  
recombination needed

$2 \rightarrow 1$  vs  $2 \rightarrow 2$  process

using EKS98

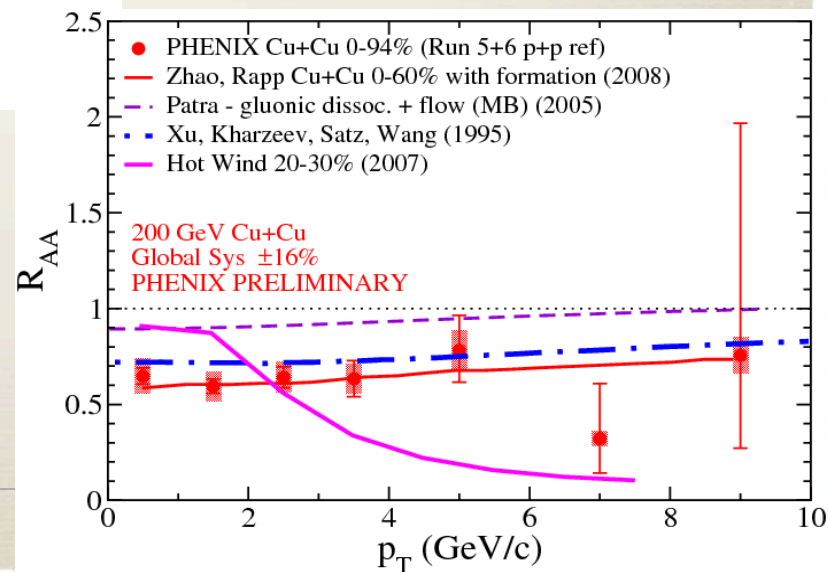
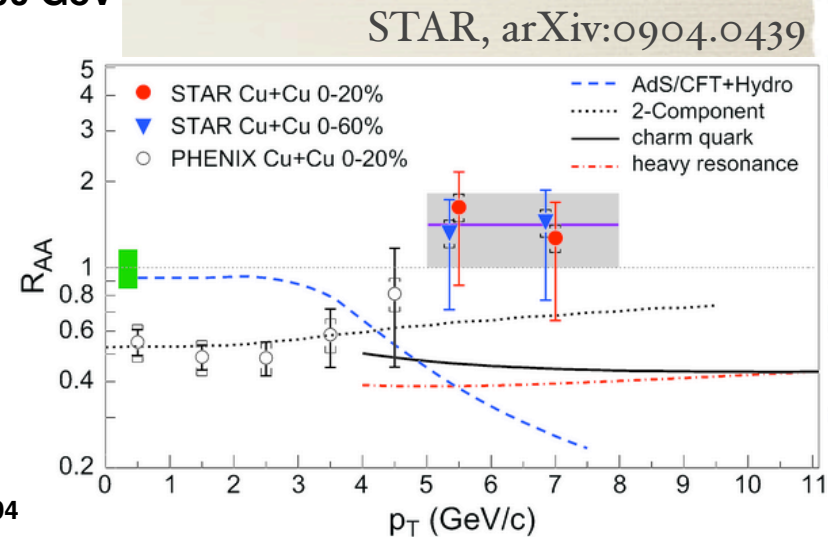
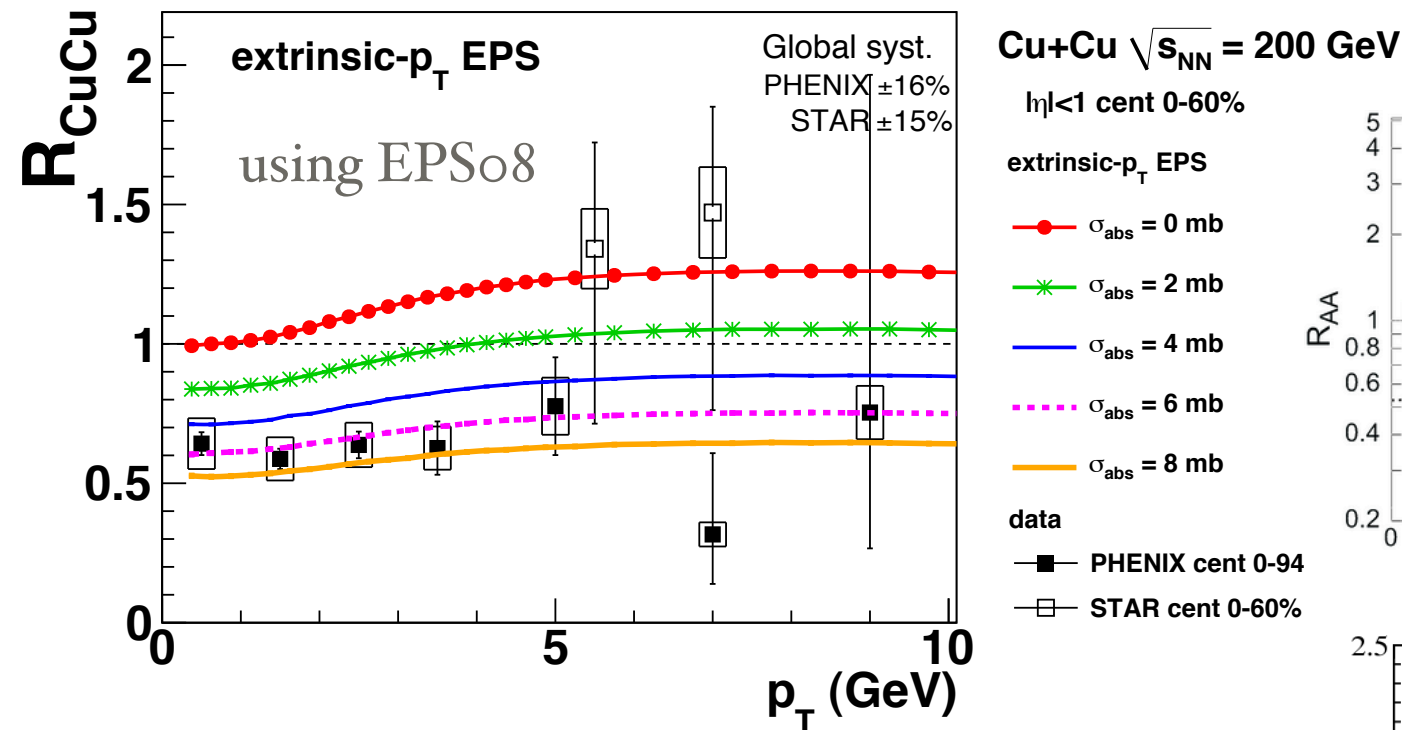
$2 \rightarrow 2$  process : mid-y & fwd-y



# CNM effects at RHIC : $J/\psi$ $R_{\text{CuCu}}$ vs $p_T$

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg and A. R.  
PRC 81 (2010) 064911

$2 \rightarrow 2$  process



$2 \rightarrow 2$  CNM effects with  $2 \rightarrow 2$  process as input also in the game

# CNM effects at RHIC : $\Upsilon$ in dAu

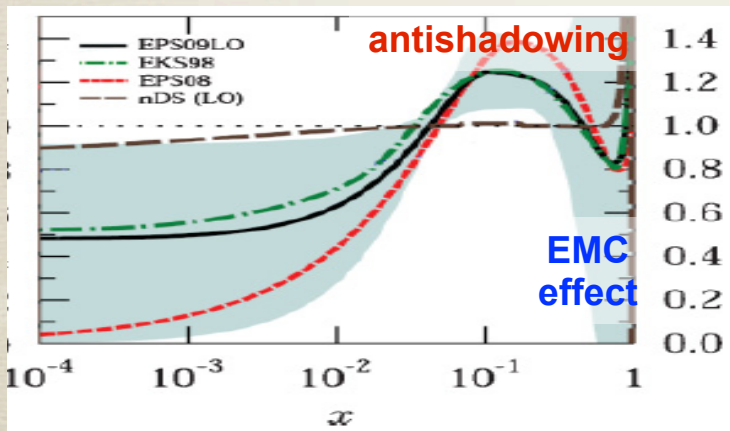
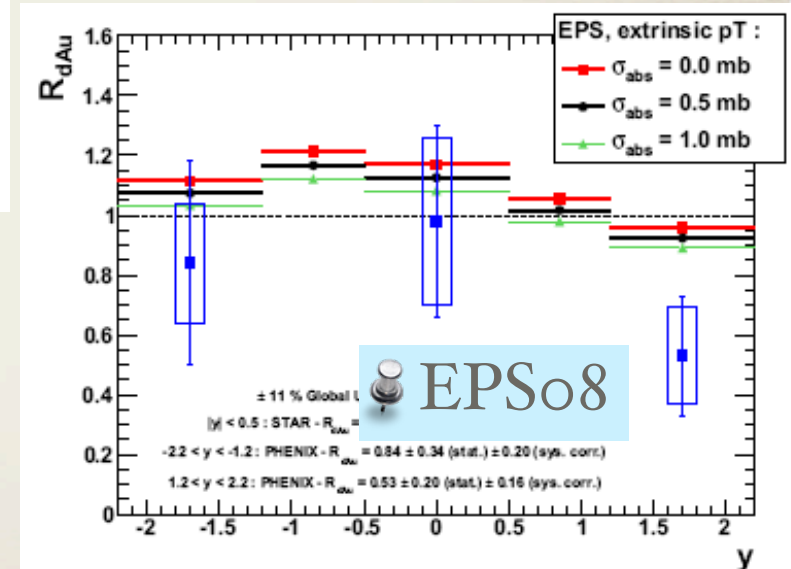
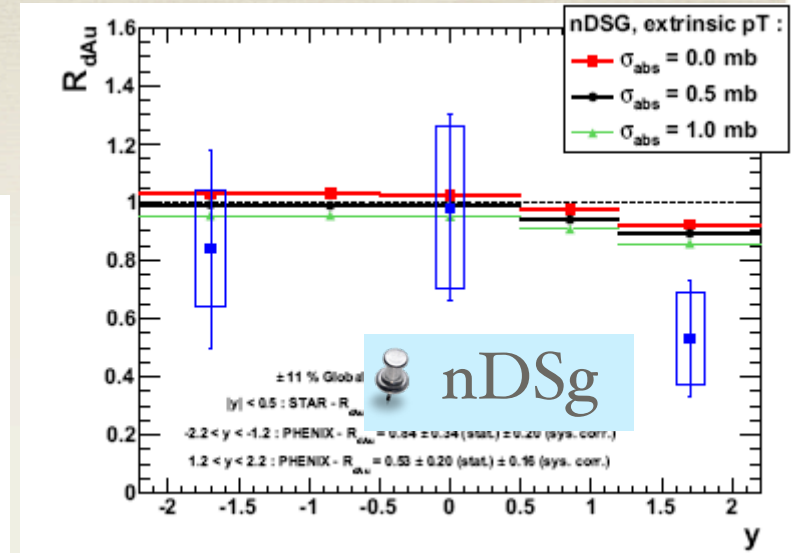
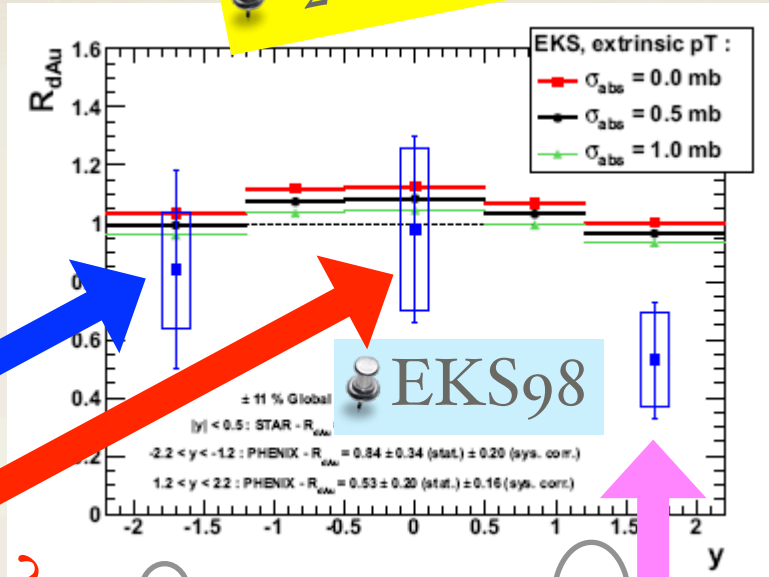
E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne and A. R.  
in preparation

**2 → 2 process**

EMC effect ?

antishadowing ?

energy loss ?



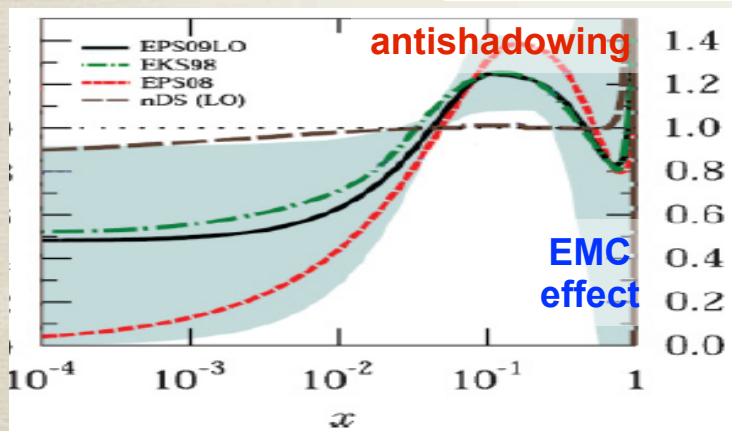
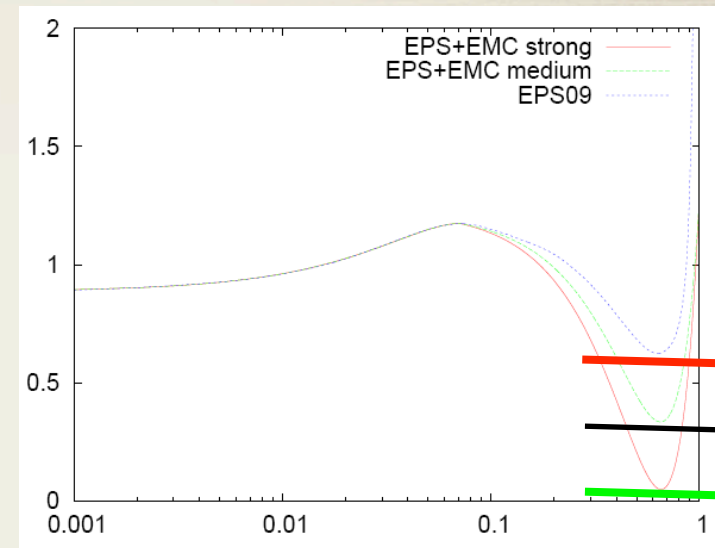
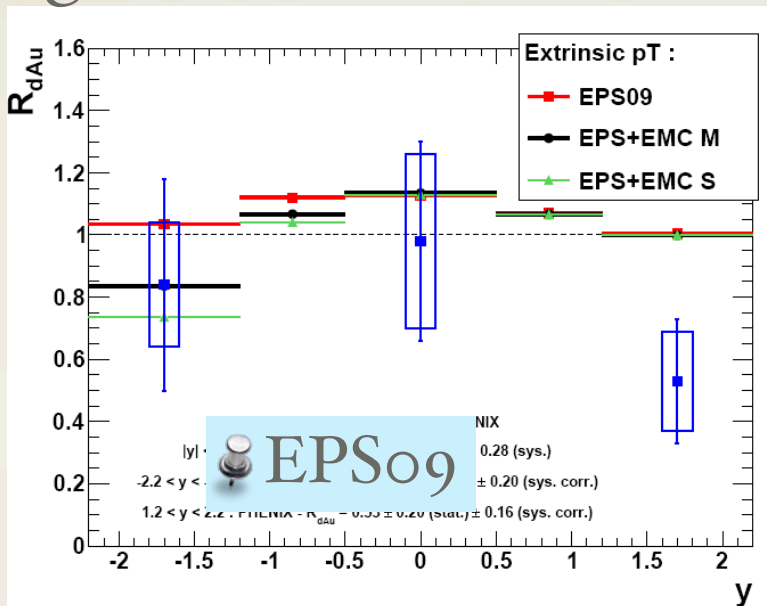
Data:  
STAR Preliminary, Nucl. Phys. A830 (2009) 235c  
PHENIX Preliminary, H. Pereira Da Costa, talk at  
The Rencontres de Moriond, March 2010

# CNM effects at RHIC : $\gamma$ in dAu

**2 → 2 process**

Let us try to increase the suppression of  $g(x)$  in the EMC region, keeping momentum conservation :  $\int xg(x) dx = Cte$

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne  
and A. R., in preparation

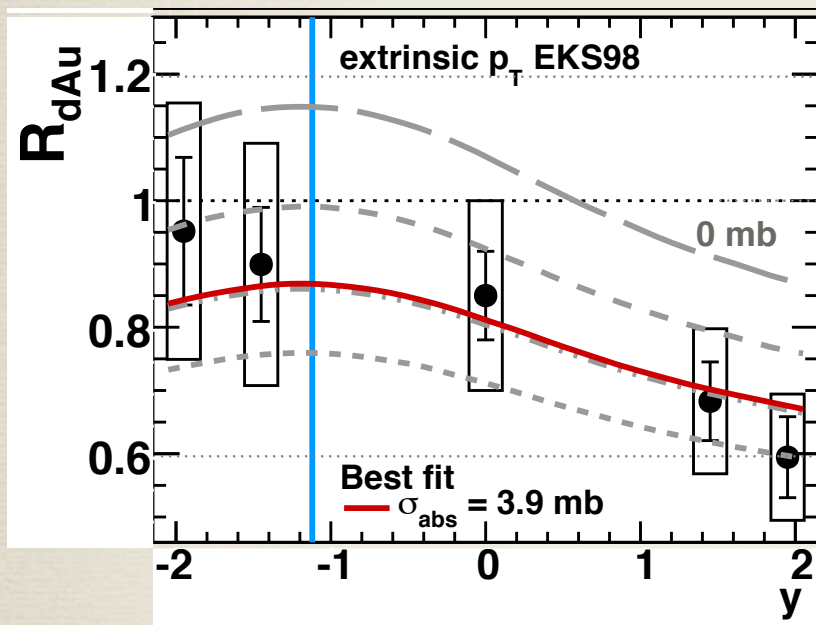


Data:  
 STAR Preliminary, Nucl. Phys. A830 (2009) 235c  
 PHENIX Preliminary, H. Pereira Da Costa, talk at  
 The Rencontres de Moriond, March 2010

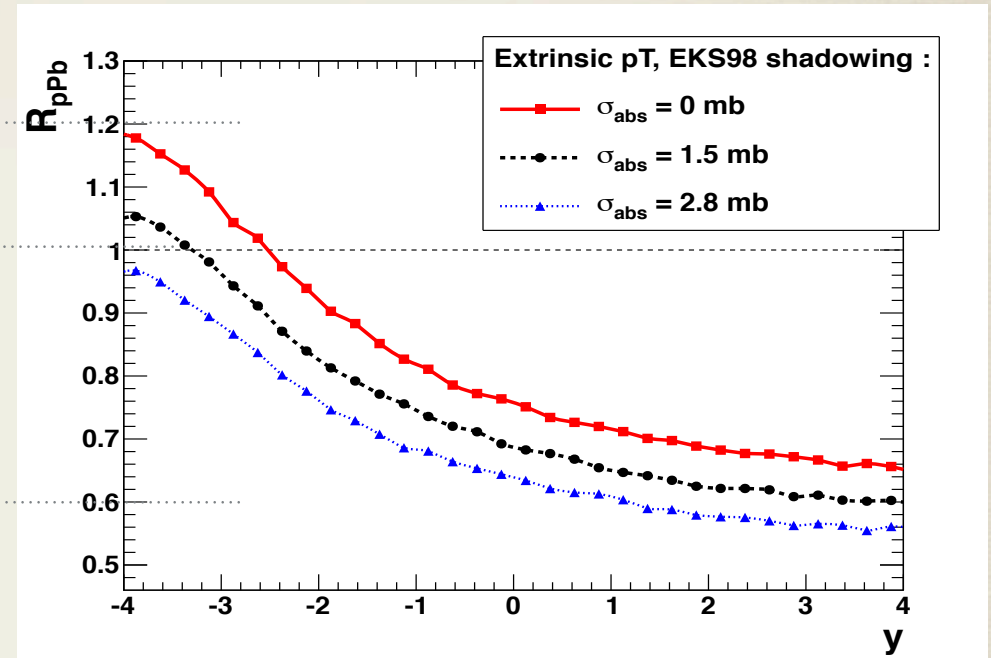
# CNM effects from RHIC to LHC

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne and A.  
R., Nucl. Phys. A 855 (2011) 327-330

  $2 \rightarrow 2$  process



RHIC dAu @ 200 GeV



LHC pPb @ 2.76 TeV

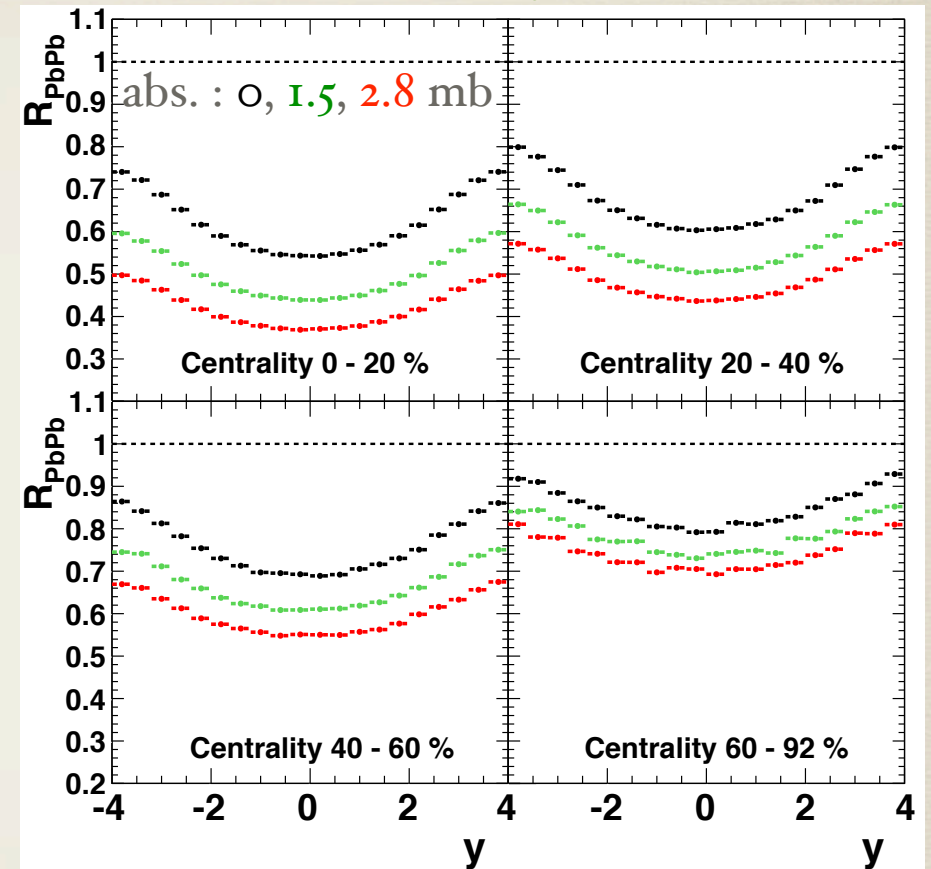
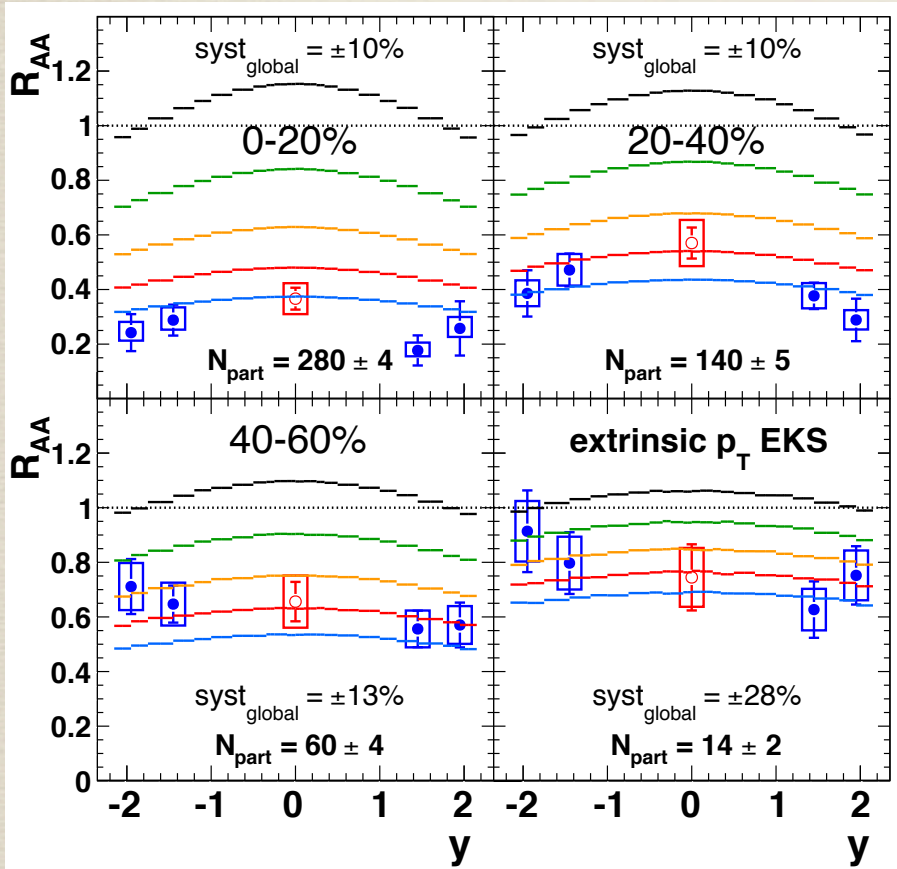
LHC : much smaller  $x$ , antshadowing peak at much lower  $y$

# CNM effects from RHIC to LHC

  $2 \rightarrow 2$  process

absorption : 0, 2, 4, 6, 8 mb

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne and A.  
R., Nucl. Phys. A 855 (2011) 327-330



RHIC AuAu @ 200 GeV

LHC PbPb @ 2.76 TeV

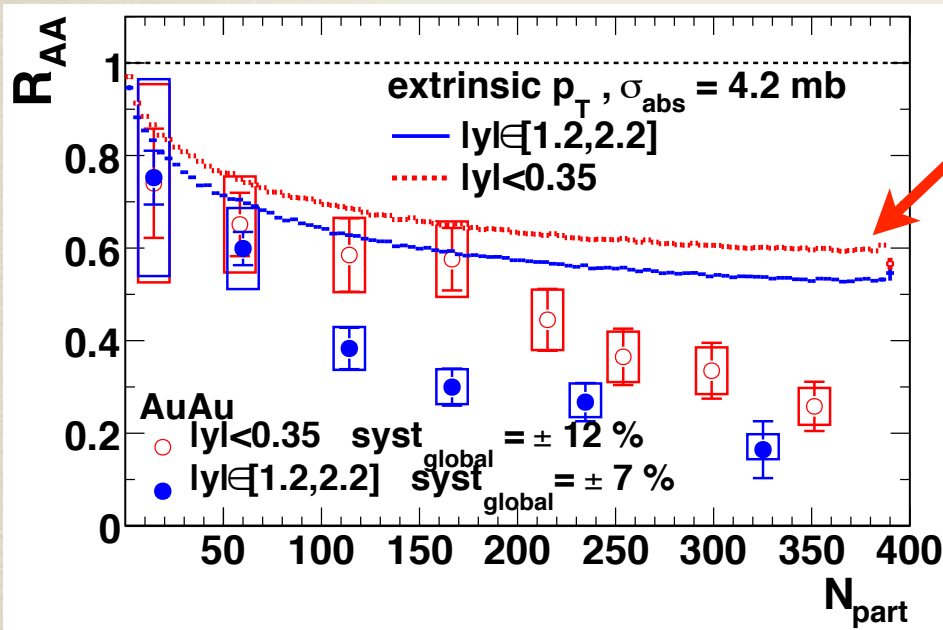
LHC : opposite shape for the  $y$  dependence per centrality bin

# CNM effects from RHIC to LHC

  $2 \rightarrow 2$  process

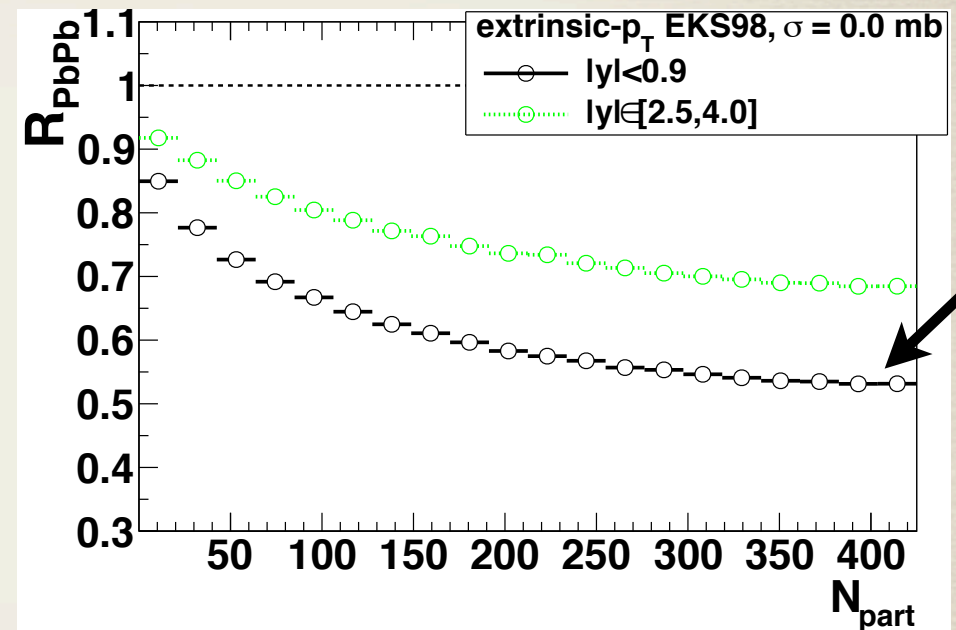
E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne and A.  
R., Nucl. Phys. A 855 (2011) 327-330

mid-y & fwd-y



RHIC AuAu @ 200 GeV

mid-y & fwd-y



LHC PbPb @ 2.76 TeV

LHC : CNM effects go in the opposite way as should go the recombination, with more suppression at mid-y than at fwd-y.



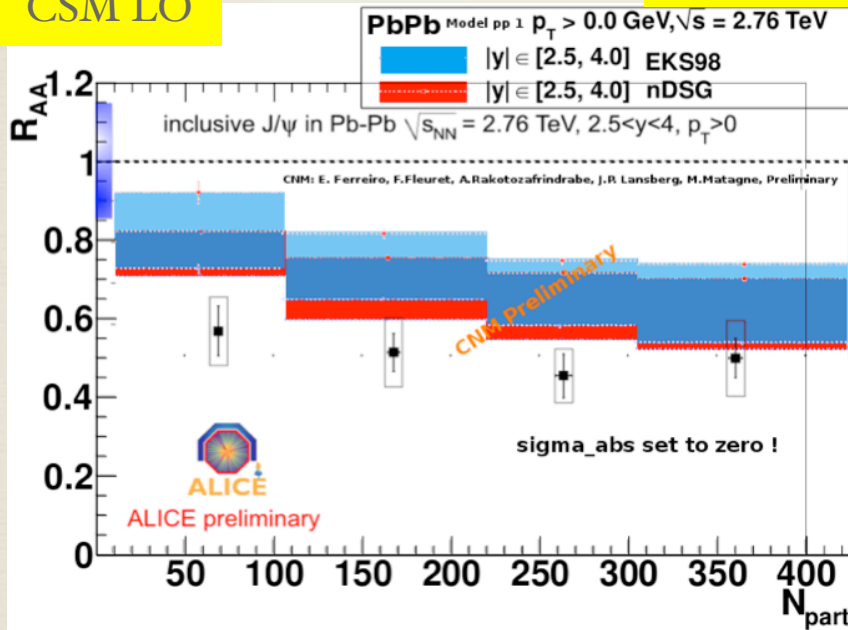
# CNM effects at LHC : $J/\psi$ in PbPb

Two models (prompt/direct  $J/\psi$ ) + nPDF and  $\mu_F$  uncertainties vs inclusive  $J/\psi$   $R_{AA}$  :

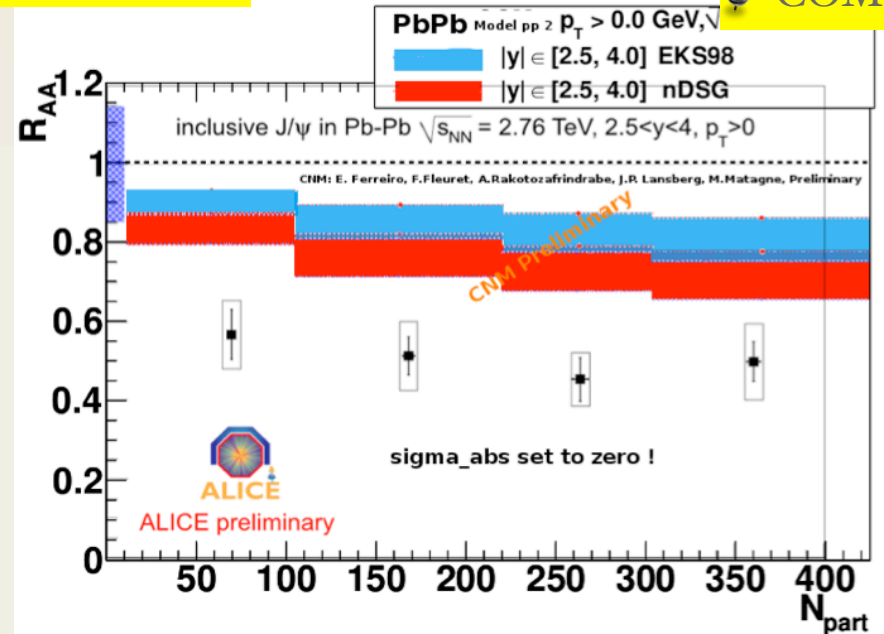
$2 \rightarrow 2$  process

E. G. Ferreira, F. Fleuret,  
J. P. Lansberg, N. Matagne  
and A. R., in preparation

CSM LO

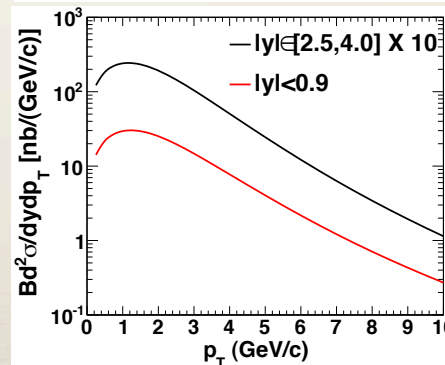
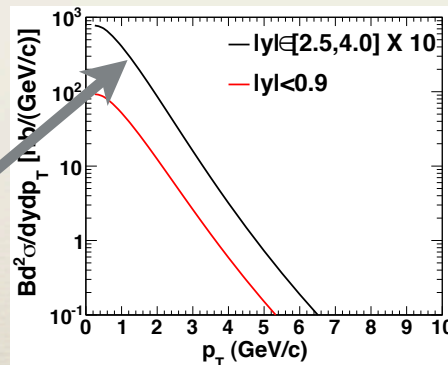


COM LO



$p_T$  spectra in p+p

closer to  $2 \rightarrow 1$  process



# CNM effects at LHC : $J/\psi$ in PbPb

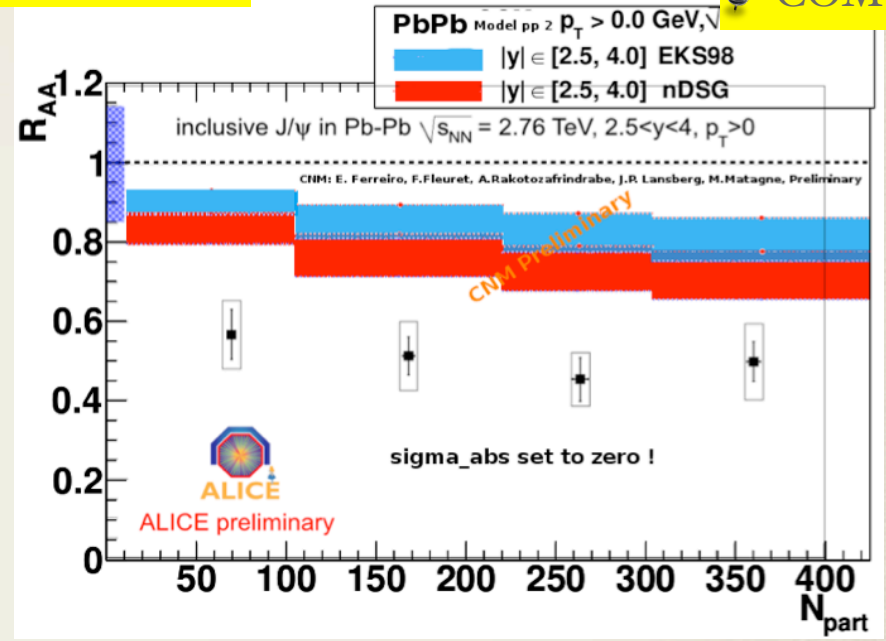
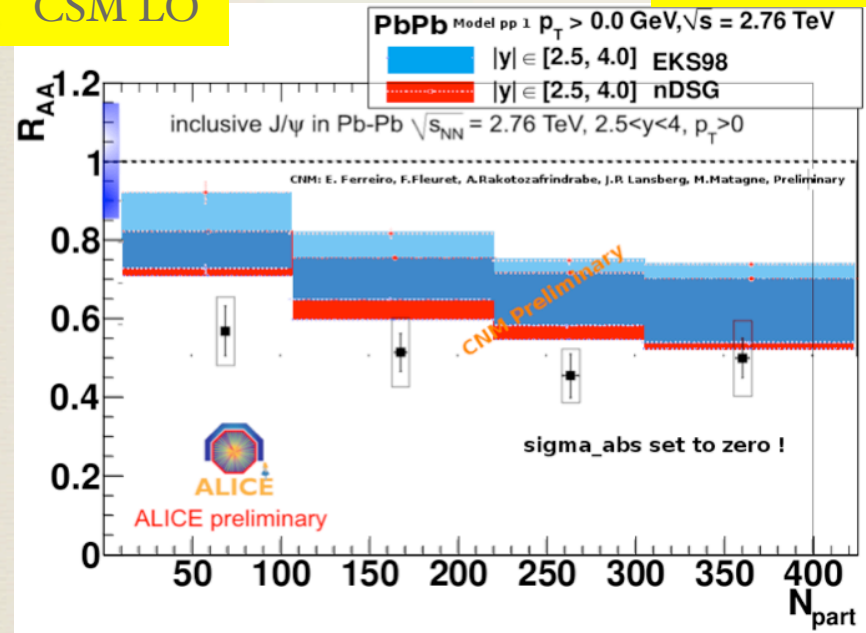
Two models (prompt/direct  $J/\psi$ ) + nPDF and  $\mu_F$  uncertainties vs inclusive  $J/\psi$   $R_{AA}$  :

$2 \rightarrow 2$  process

CSM LO

COM LO

E. G. Ferreira, F. Fleuret,  
J. P. Lansberg, N. Matagne  
and A. R., in preparation



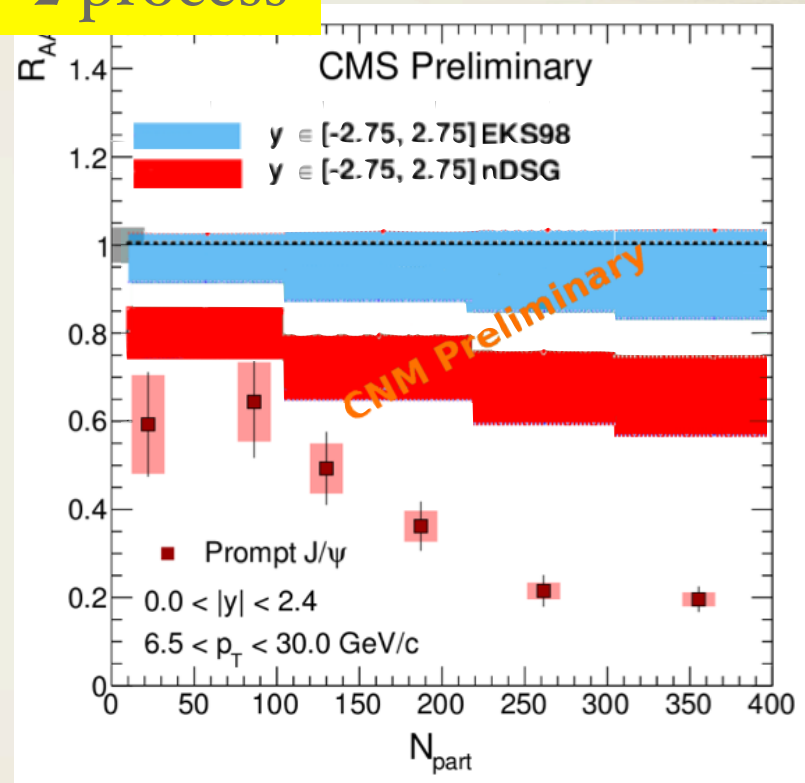
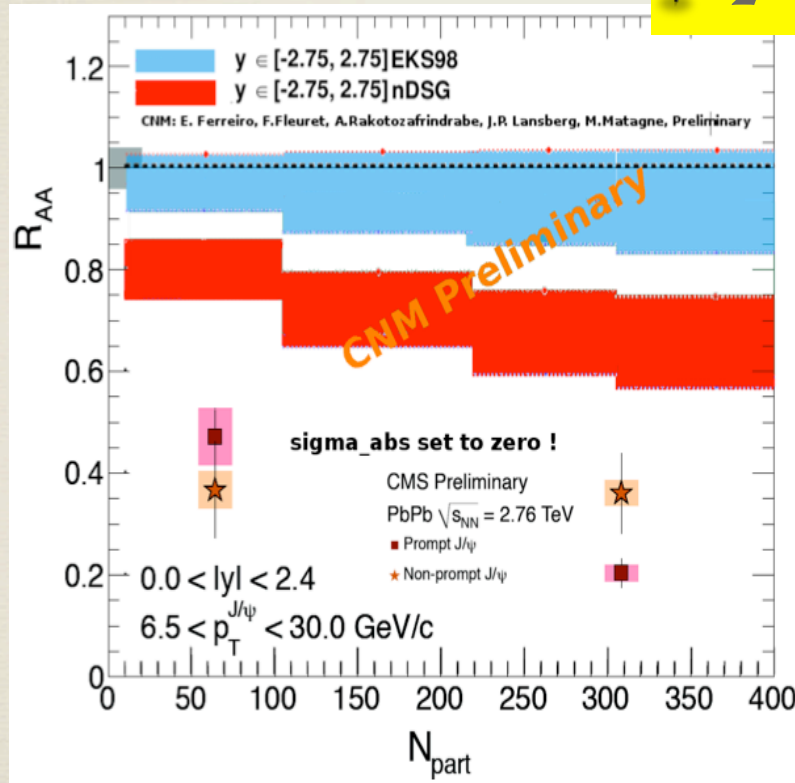
- ✓ Alice data dominated by low  $p_T$   $J/\psi$
- ✓ PbPb suppression : about the same or slightly more than CNM effects
- ➔ hint of a large recombination effect as expected at low  $p_T$  ?

# CNM effects at LHC : $J/\psi$ in PbPb

Going to mid-y, adding a  $p_T$  cut  $> 6.5$  GeV/c to the most conservative model and to the data :

$2 \rightarrow 2$  process

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg, N. Matagne  
and A. R., in preparation



- ☑ PbPb suppression goes much further than CNM effects
- ➔ confirms that recombination is small at high  $p_T$  ?

# Conclusion and outlook (1/2)

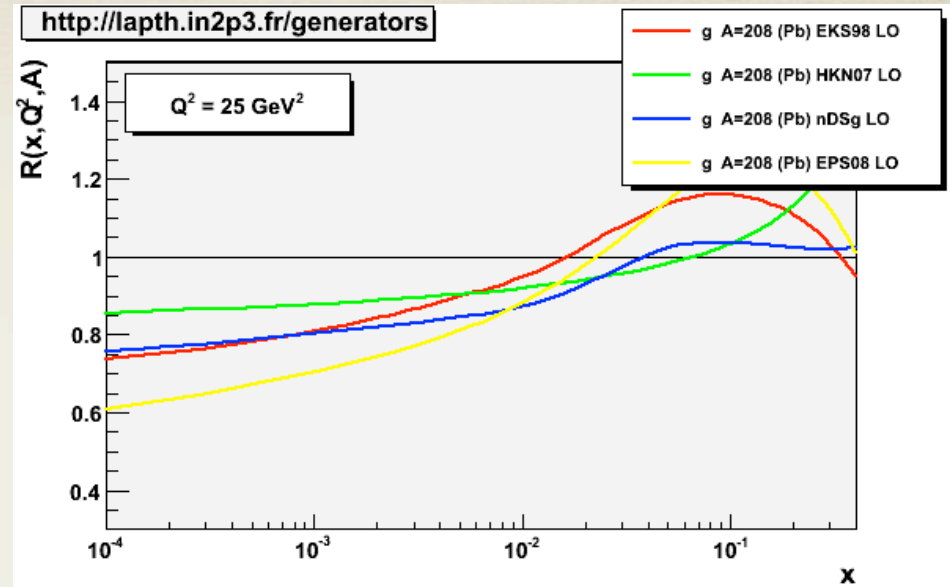
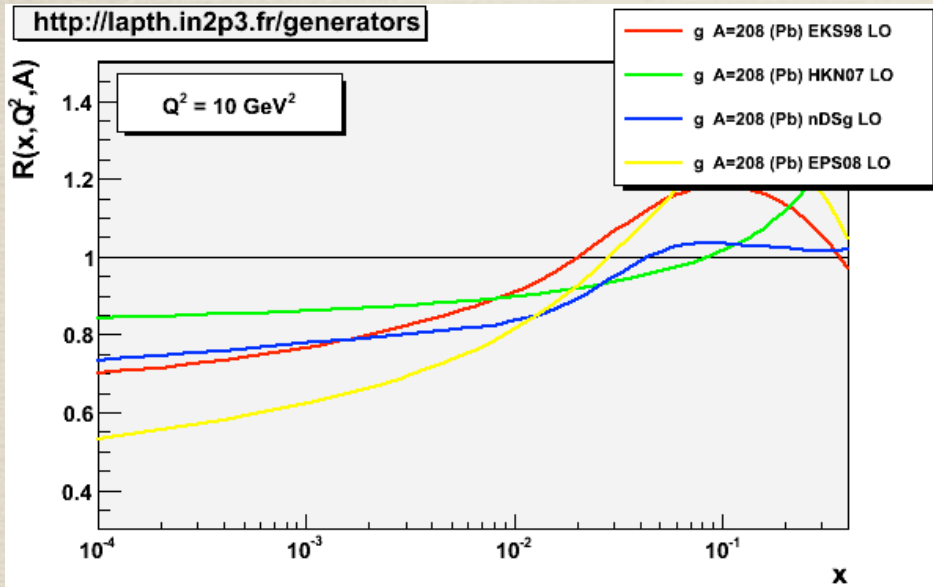
- Shadowing computations can benefit from an improved knowledge of quarkonia production process
- Within the extrinsic scheme (2 to 2 process) :
  - more complex kinematics
  - for a given  $y$ ,  $x$  is larger  $\Rightarrow$  larger absorption cross-section needed to match RHIC data, and antishadowing peak shifted to larger  $y$
  - Upsilon could be used as a tool to explore antishadowing and EMC effect

# Conclusion and outlook (2/2)

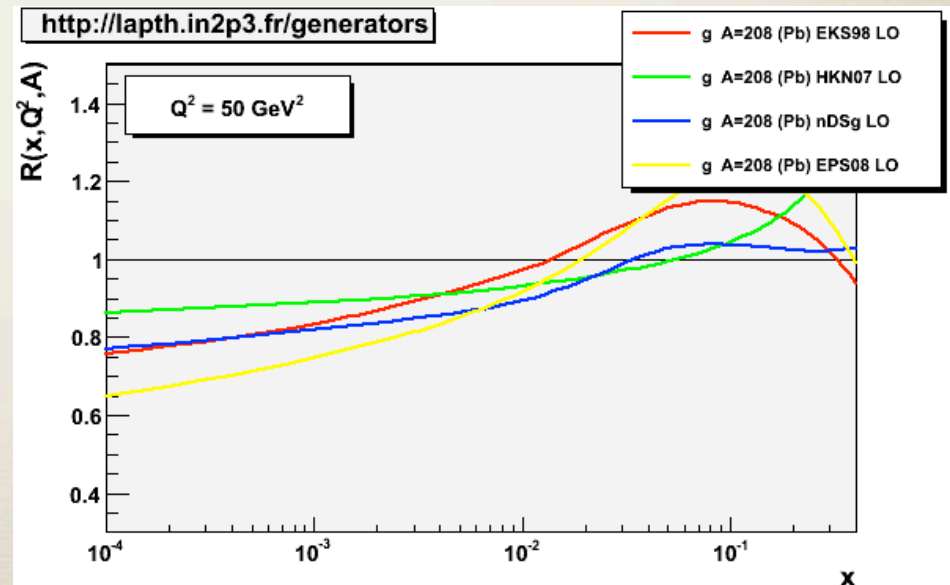
- Shadowing computations suffer from large uncertainties. Especially large is the uncertainty due to  $\mu_F$  for the  $J/\psi$ .
- Hints of recombination at low  $p_T$ , and of a negligible recombination at large  $p_T$ .
- It will be difficult to understand the present suppression pattern of the  $J/\psi$  if we do not reduce these uncertainties on shadowing computations.
- pA run with a good precision on the  $y$ ,  $p_T$  and centrality dependence
- move to heavier quarkonium states at LHC ?

# BACK-UP

# nPDF for Pb

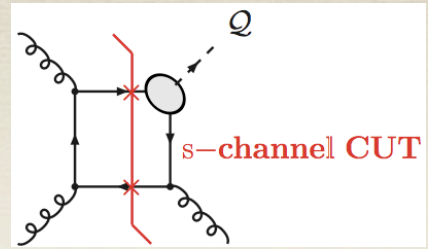
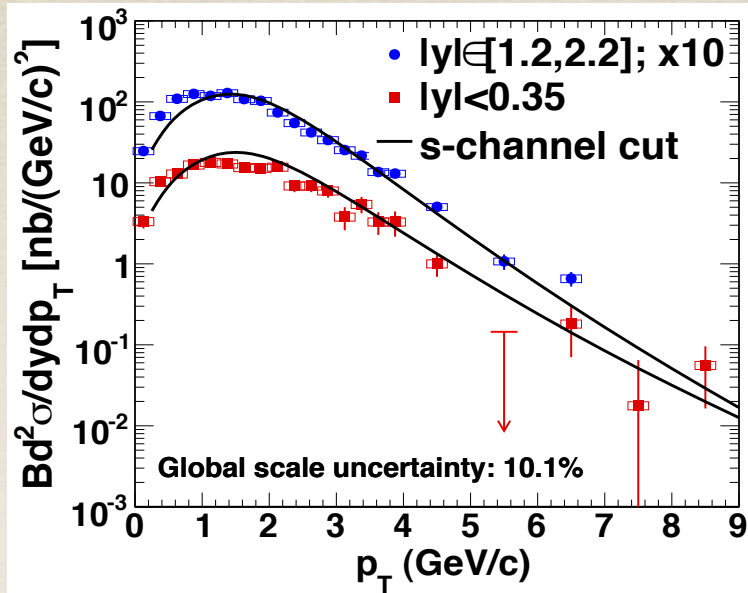


EKS<sub>98</sub>  
HKN<sub>07</sub>  
nDSg  
EPS<sub>08</sub>



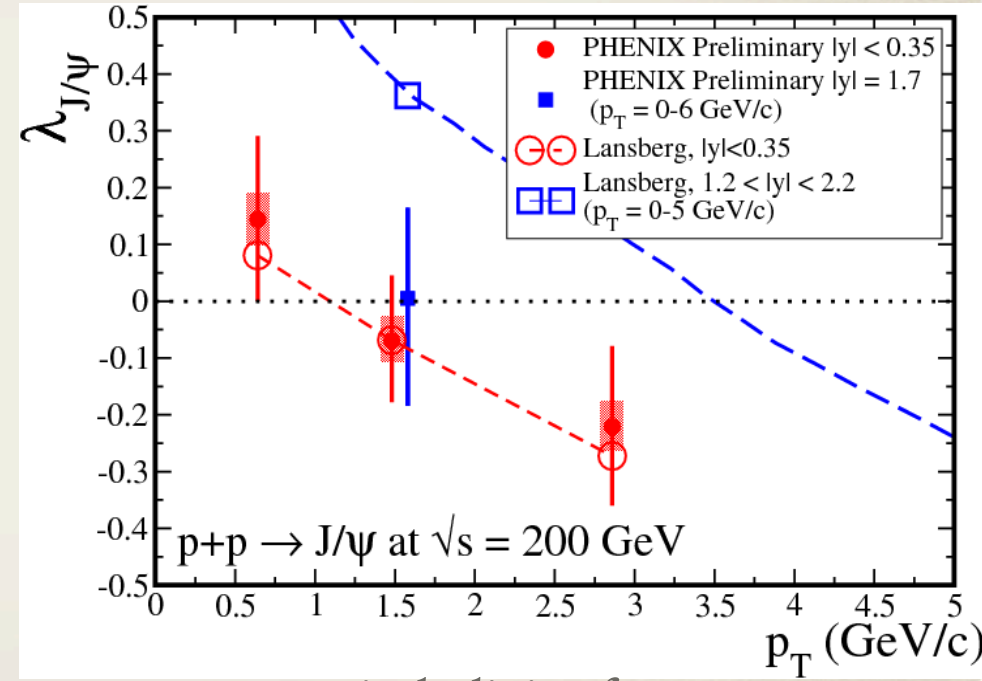
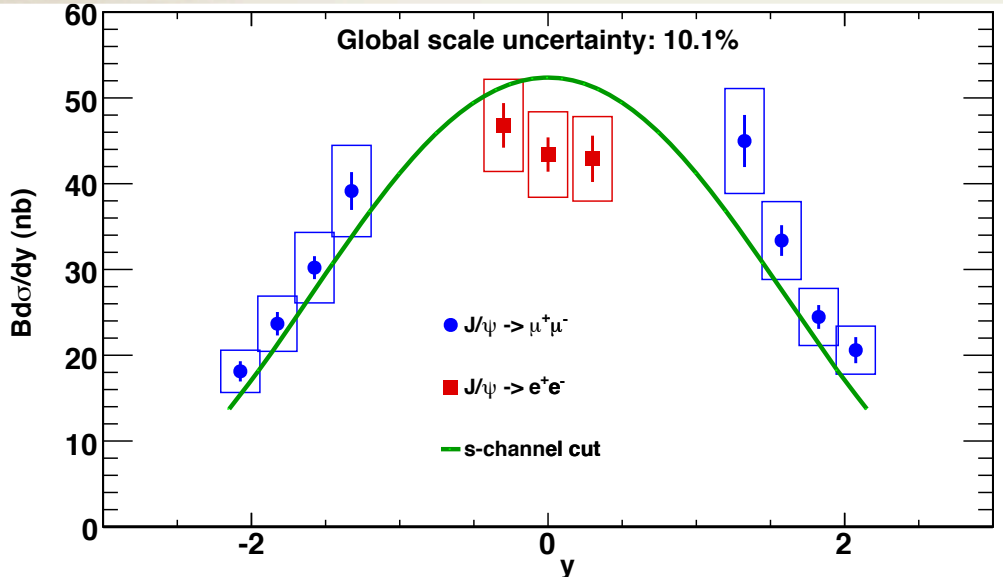
# CSM + s-channel cut at RHIC in p+p

yield vs  $p_T$



polarisation vs  $p_T$  at mid-y and fwd-y

data : phenix run5 p+p spectra



in helicity frame

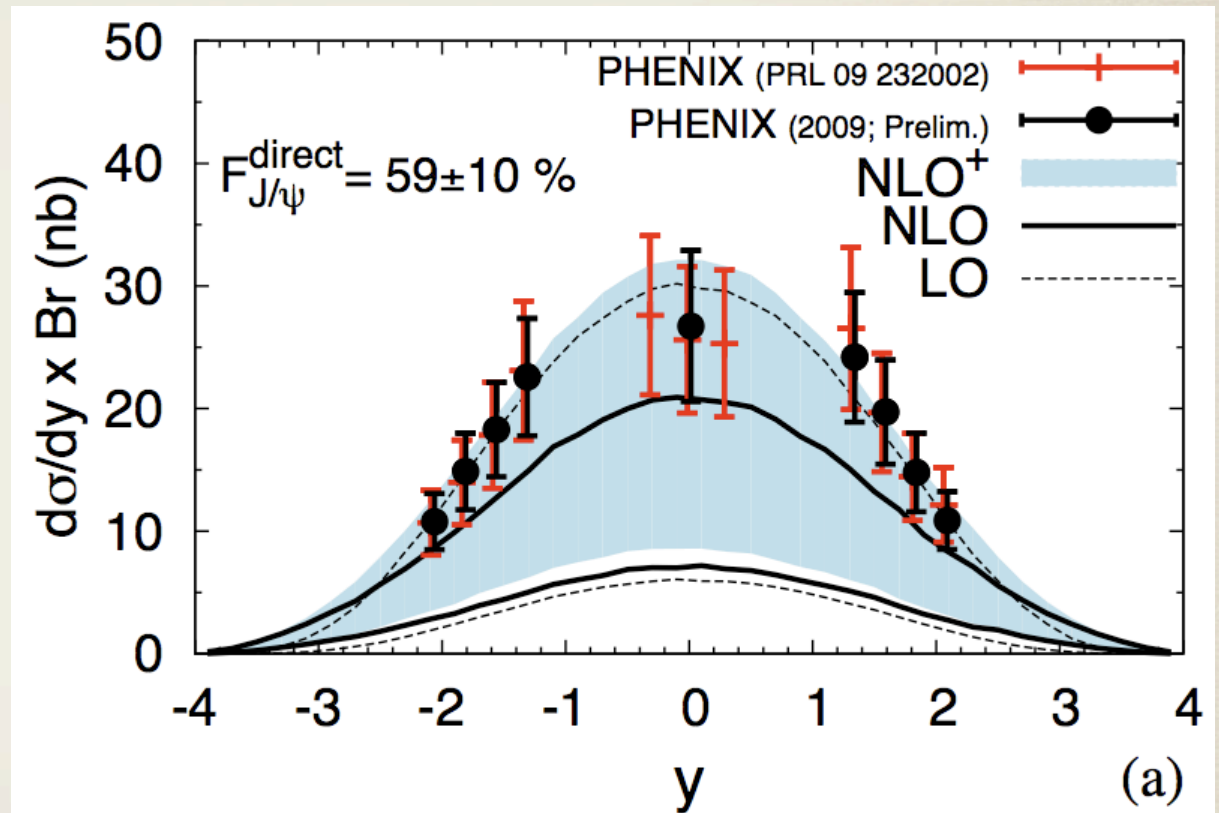
yield vs rapidity



# CSM LO, NLO, $NLO+cg$ vs the $y$ spectra in $p+p$ at RHIC

direct  $J/\psi$

points :  $y$  spectra for  
the *direct*  $J/\psi$ ,  
as extrapolated from  
PHENIX spectra



S. J. Brodsky and J. P. Lansberg, arXiv:0908.0754

# The $y$ spectra in $p+p$ at RHIC

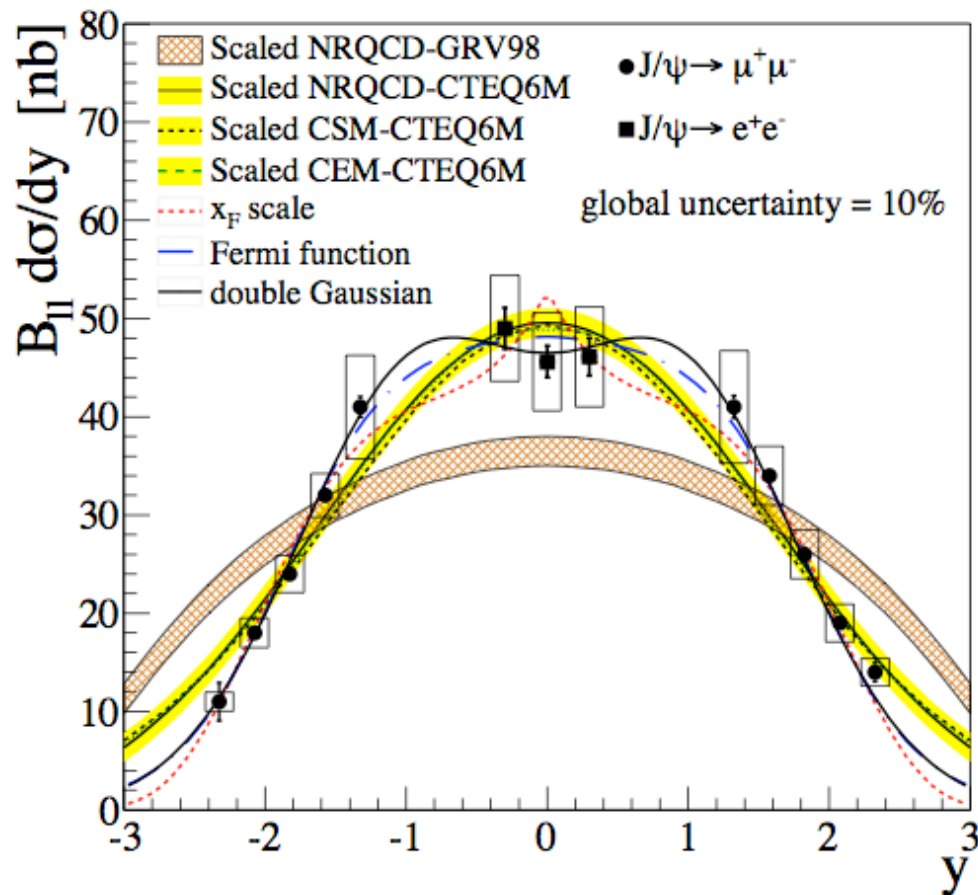
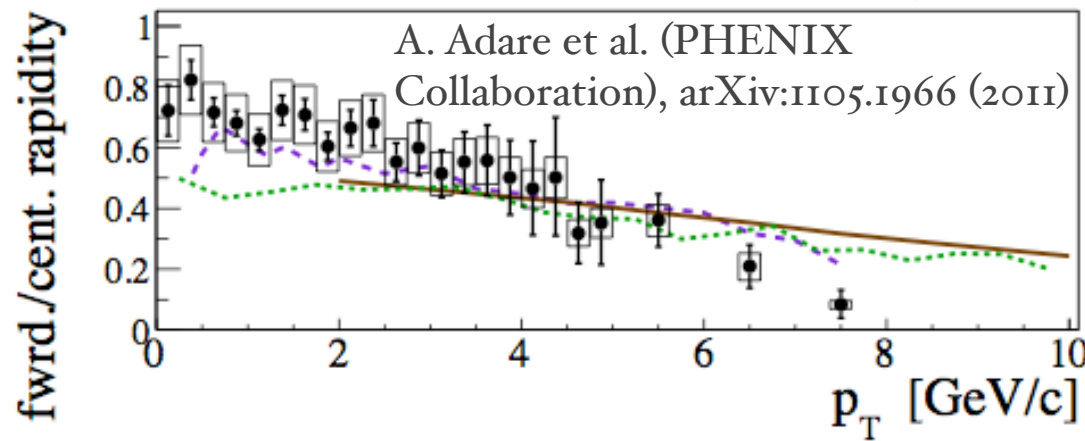
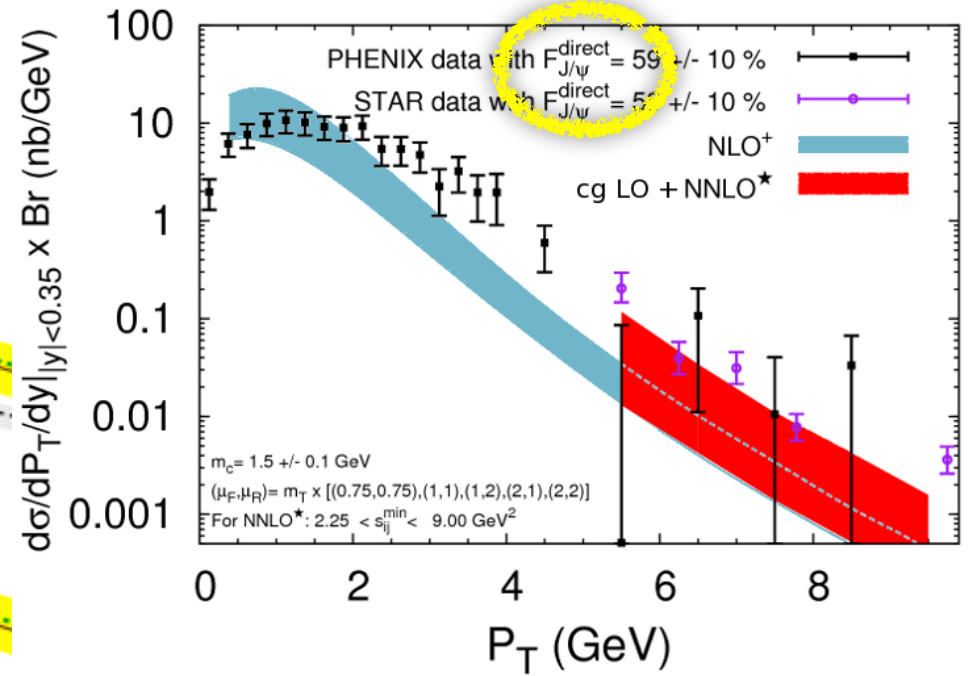
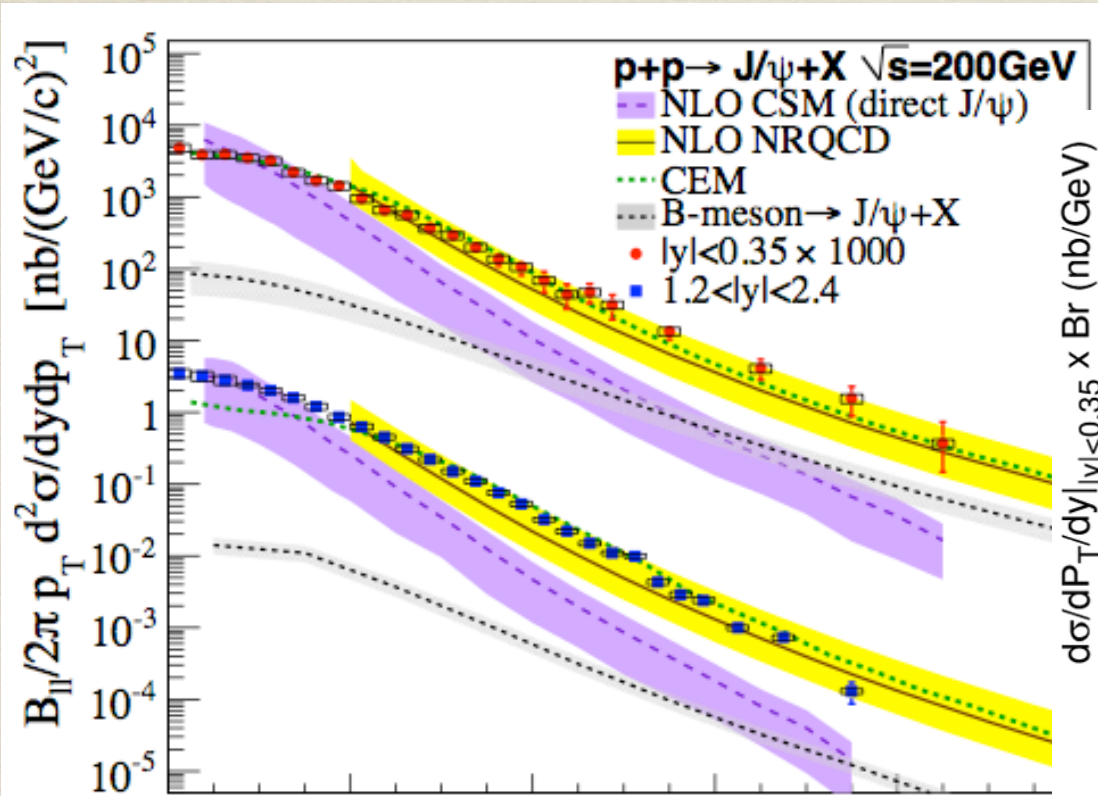


TABLE V: Comparison of the measured  $J/\psi$  cross section with the three models considered in this text. Direct  $J/\psi$  cross sections are obtained assuming that the  $\chi_c$  and  $\psi'$  feed-down fractions measured at midrapidity are the same at forward rapidity. Type A, type B and type C errors are quadratically summed in the measured result.

	direct $J/\psi$	inclusive
CEM	-	$169 \pm 30$ nb
NLO CSM	$53 \pm 26$ nb	-
LO NRQCD	-	$140 \pm 5$ nb
Measured	$105 \pm 26$ nb	$181 \pm 22$ nb

A. Adare et al. (PHENIX Collaboration), arXiv:1105.1966 (2011)

# The $p_T$ spectra in $p+p$ at RHIC

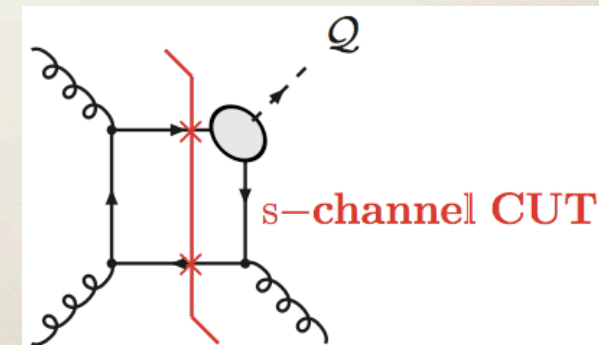
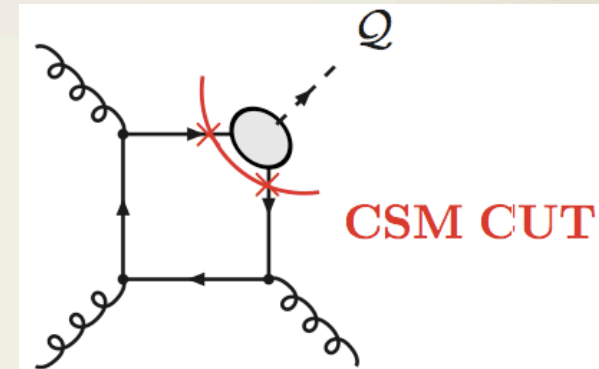
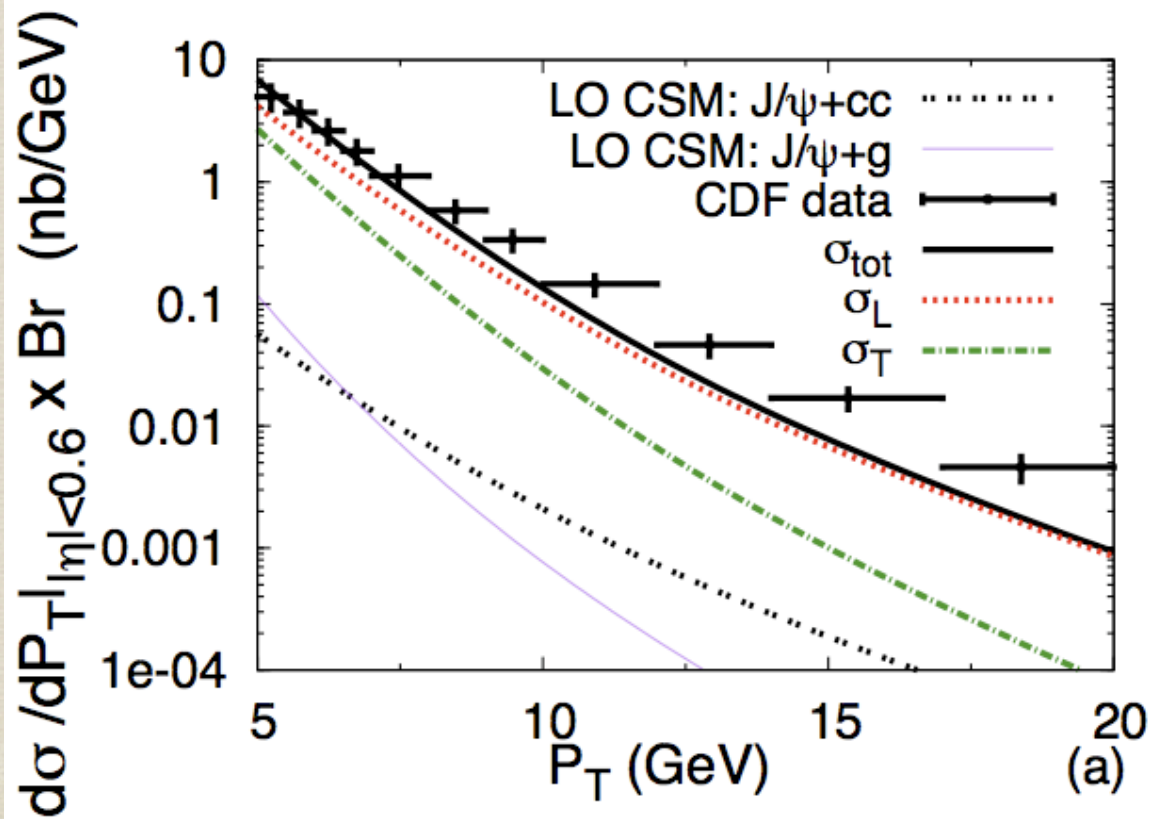


Brodsky, Lansberg, PRD 81:051502 (R) (2010)



# S-channel cut model vs Tevatron $p_T$ spectra

H. Habertzettl et J. P. Lansberg,  
PRL 100, 032006 (2008)



# «Clean» exp. results ...



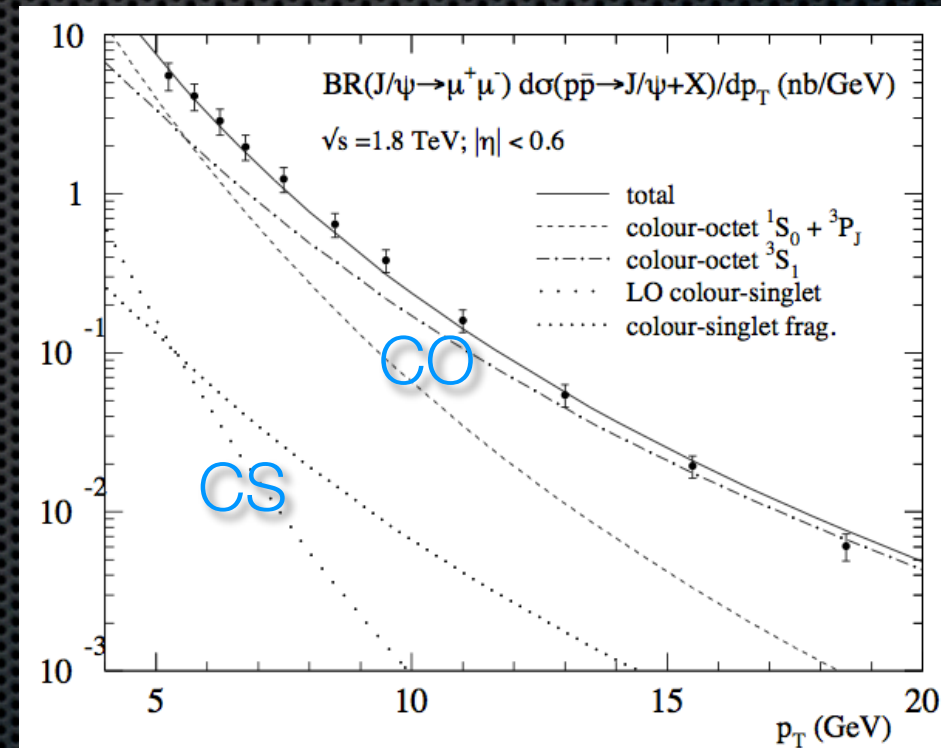
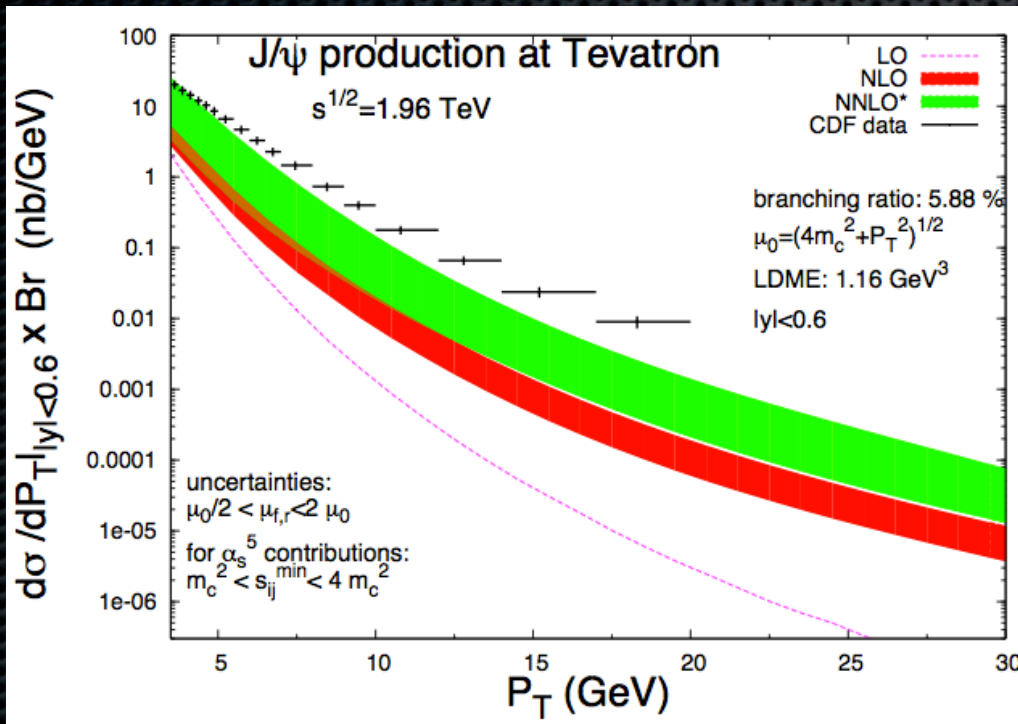
... not so clear underlying theory

Artoisenet, Lansberg, Maltoni, PLB 653:60 (2007); Campbell, Maltoni, Tramontano, PRL 98:252002 (2007); Gong, Wang, PRL 100:232001 (2008)

Artoisenet, Campbell, Lansberg, Maltoni, Tramontano (in progress) Kramer, Prog. Part. Nucl. Phys. 47:141 (2001)

CSM NNLO\*

NRQCD - CO dominance



Direct J/ψ vs p<sub>T</sub> :

- ✦ not satisfactory for CSM even with higher order corrections included
- ✦ good agreement for NRQCD, with matrix elements tuned with CDF data

# CSM @ LO and higher orders

$\gamma(1S)$

P. Artoisenet, J. Campbell, J.P. Lansberg, F. Maltoni, PRL 101 (2008) 152001.

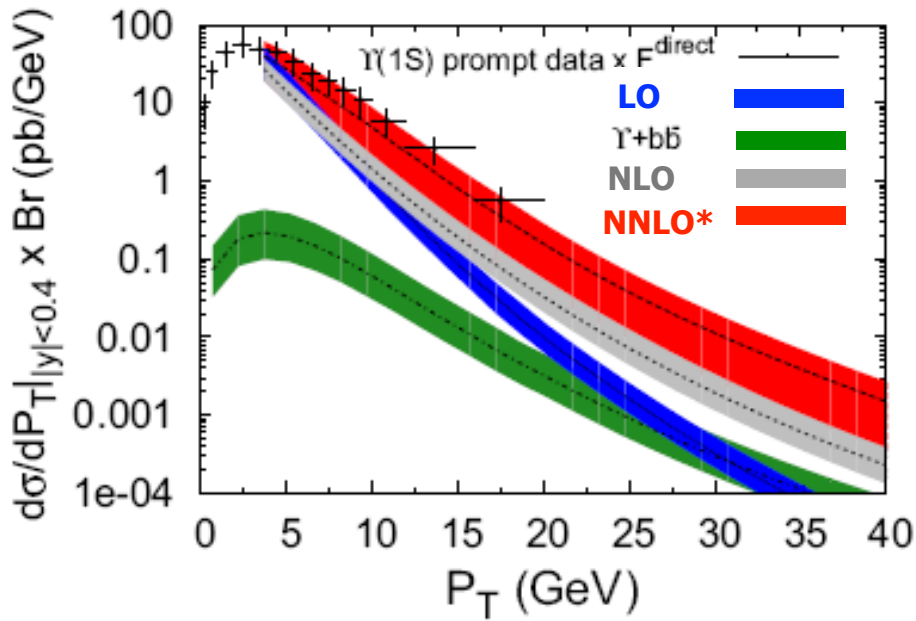
D. Acosta et al. (CDF Collaboration), PRL 88 (2002) 161802.

Tevatron (1.8 TeV):

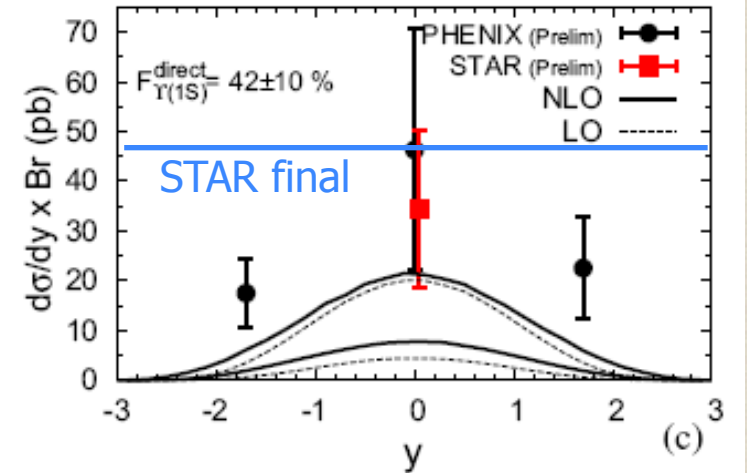
CSM LO sufficient to describe low  $p_T$  data

RHIC (200 GeV):

CSM LO below the data ?



S.J. Brodsky, J.P. Lansberg, PRD 81 (2010) 014004.

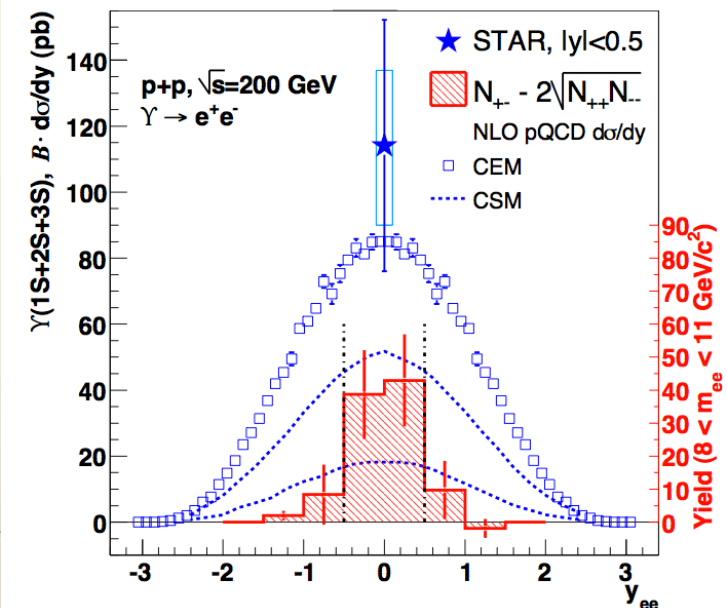


STAR preliminary [STAR Collaboration, J. Phys. G 34 (2007) S947.]

$$BR \times (d\sigma/dy)_{y=0} = 91 \pm 28 \text{ (stat.)} \pm 22 \text{ (syst.)}$$

STAR final [STAR Collaboration, arXiv:1001.2745.]

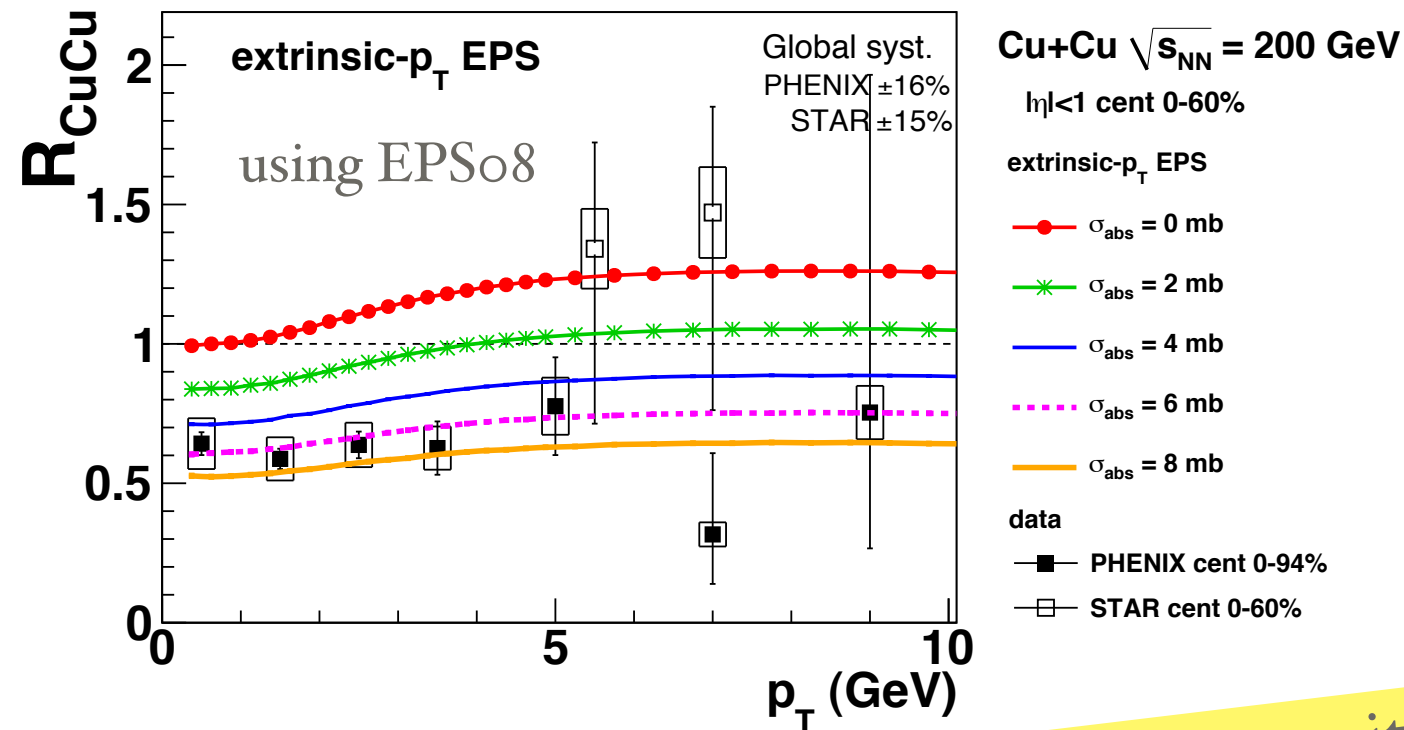
$$BR \times (d\sigma/dy)_{y=0} = 114 \pm 38^{+23}_{-24}$$





# CNM effects at RHIC : $J/\psi$ $R_{\text{CuCu}}$ vs $p_T$

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg and A. R.  
PRC 81 (2010) 064911

  $2 \rightarrow 2$  process



 CNM effects with  $2 \rightarrow 2$  process as input also in the game

 Less suppression due to shadowing with increasing  $p_T$

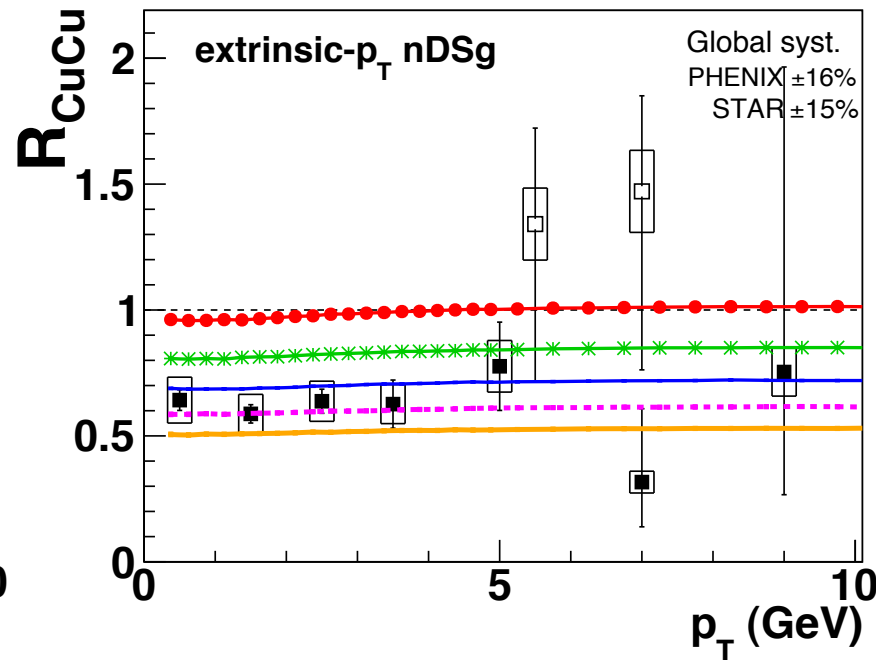
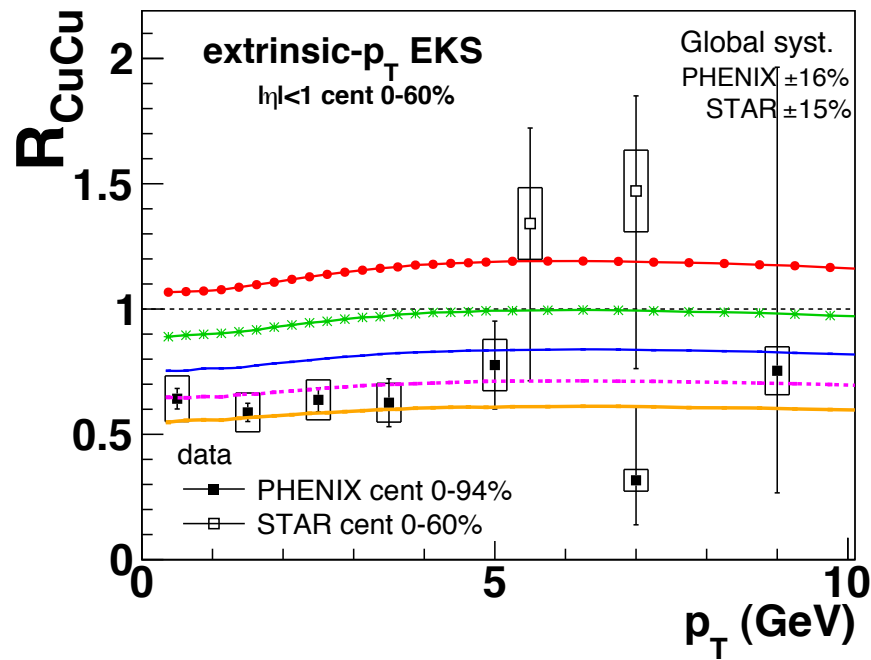
# CNM effects : $J/\psi$ $R_{\text{CuCu}}$ vs $p_T$

**2 → 2 process**

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg and A. R.  
PRC 81 (2010) 064911

EKS98

nDSg



**Cu+Cu  $\sqrt{s_{\text{NN}}} = 200$  GeV**  
 $|\eta| < 1$  cent 0-60%

**extrinsic- $p_T$  nDSg**

- $\sigma_{\text{abs}} = 0$  mb
- \*  $\sigma_{\text{abs}} = 2$  mb
- $\sigma_{\text{abs}} = 4$  mb
- ⋯  $\sigma_{\text{abs}} = 6$  mb
- $\sigma_{\text{abs}} = 8$  mb

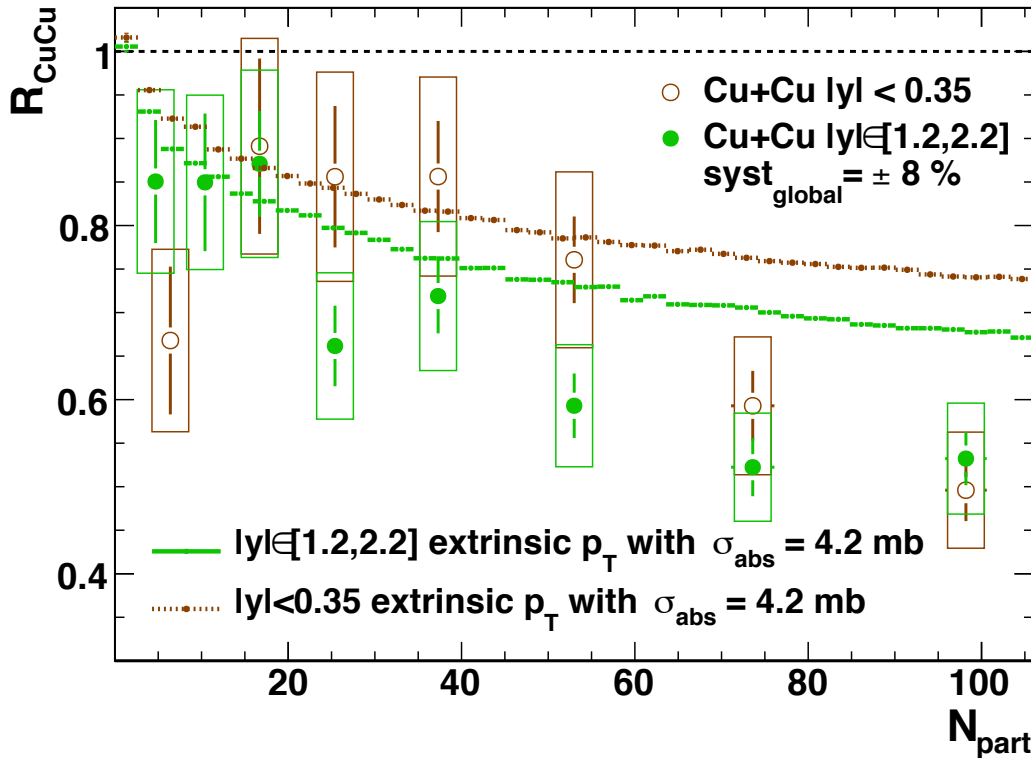
**data**

- PHENIX cent 0-94%
- STAR cent 0-60%



# CNM effects at RHIC : $J/\psi$ in CuCu

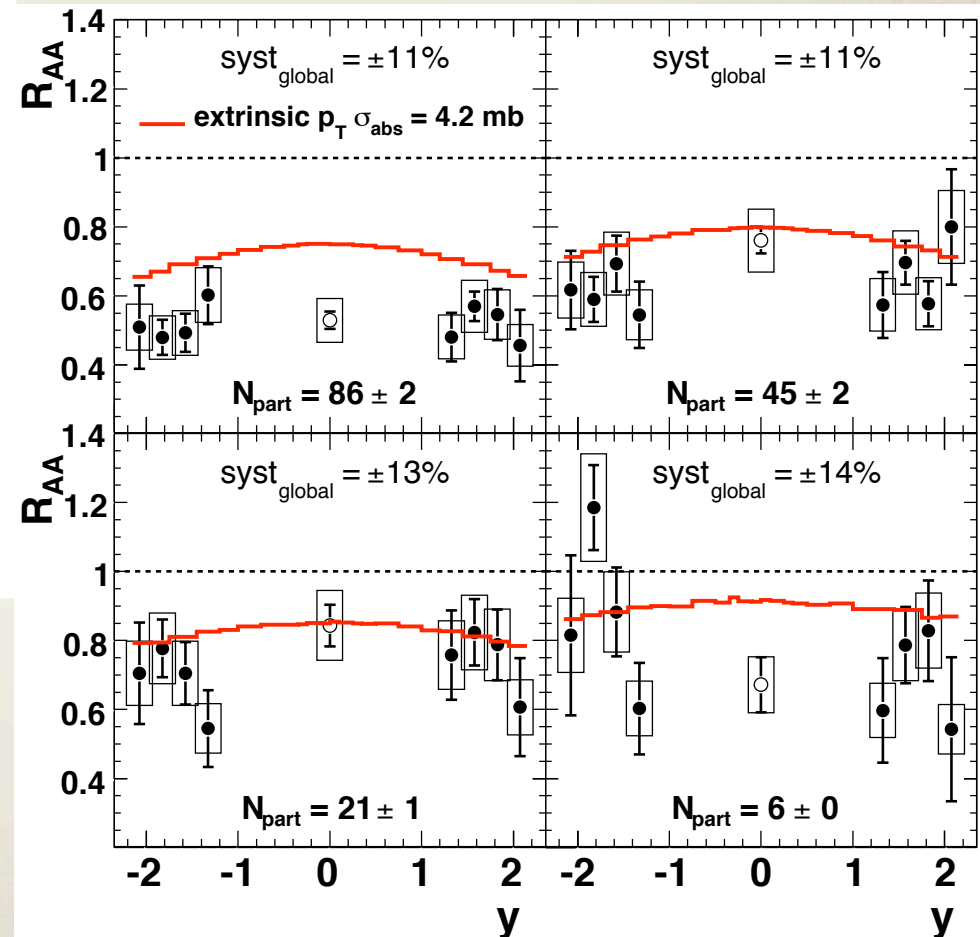
2 $\rightarrow$ 2 process : mid-y & fwd-y



E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg and A. R.  
PLB 680, 50-55 (2009)

2  $\rightarrow$  2 process

using EKS98



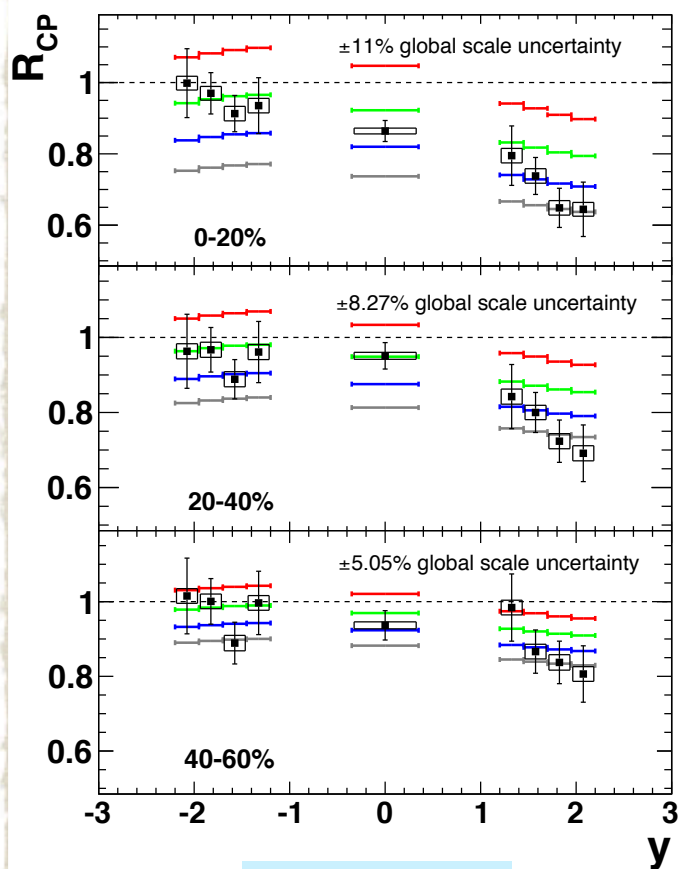
# CNM effects at RHIC : $J/\psi$ in dAu

E. G. Ferreiro, F. Fleuret, J. P. Lansberg  
and A. R., PRC 81 (2010) 064911

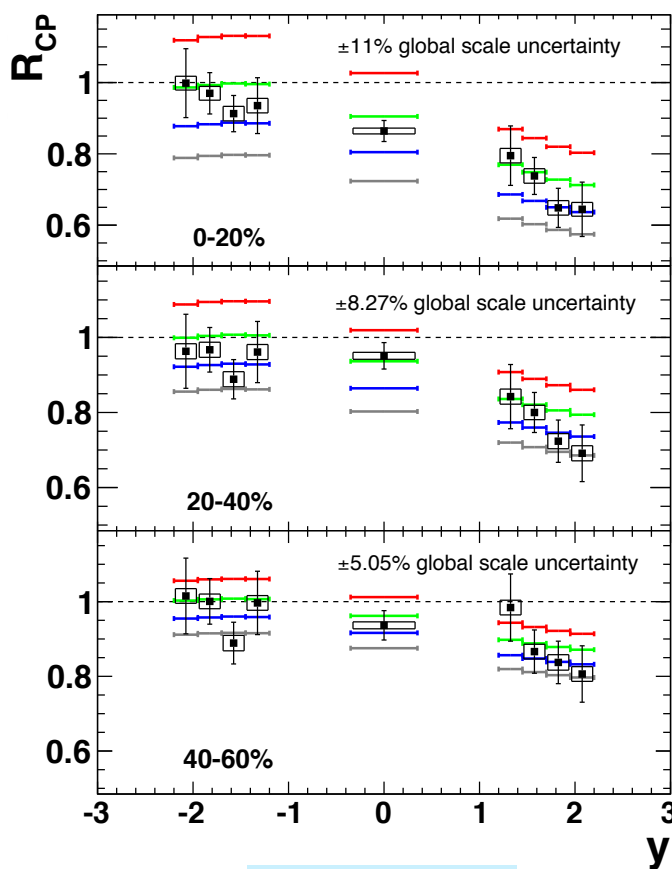
$2 \rightarrow 2$  process

absorption : 0, 2, 4, 6 mb

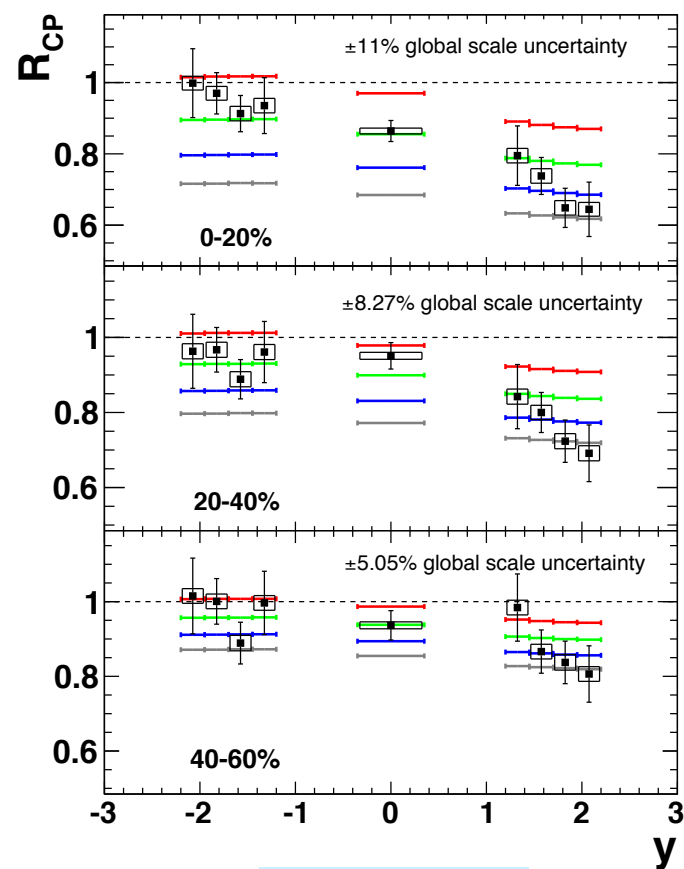
$R_{CP}$  in d+Au vs rapidity



EKS98



EPS08



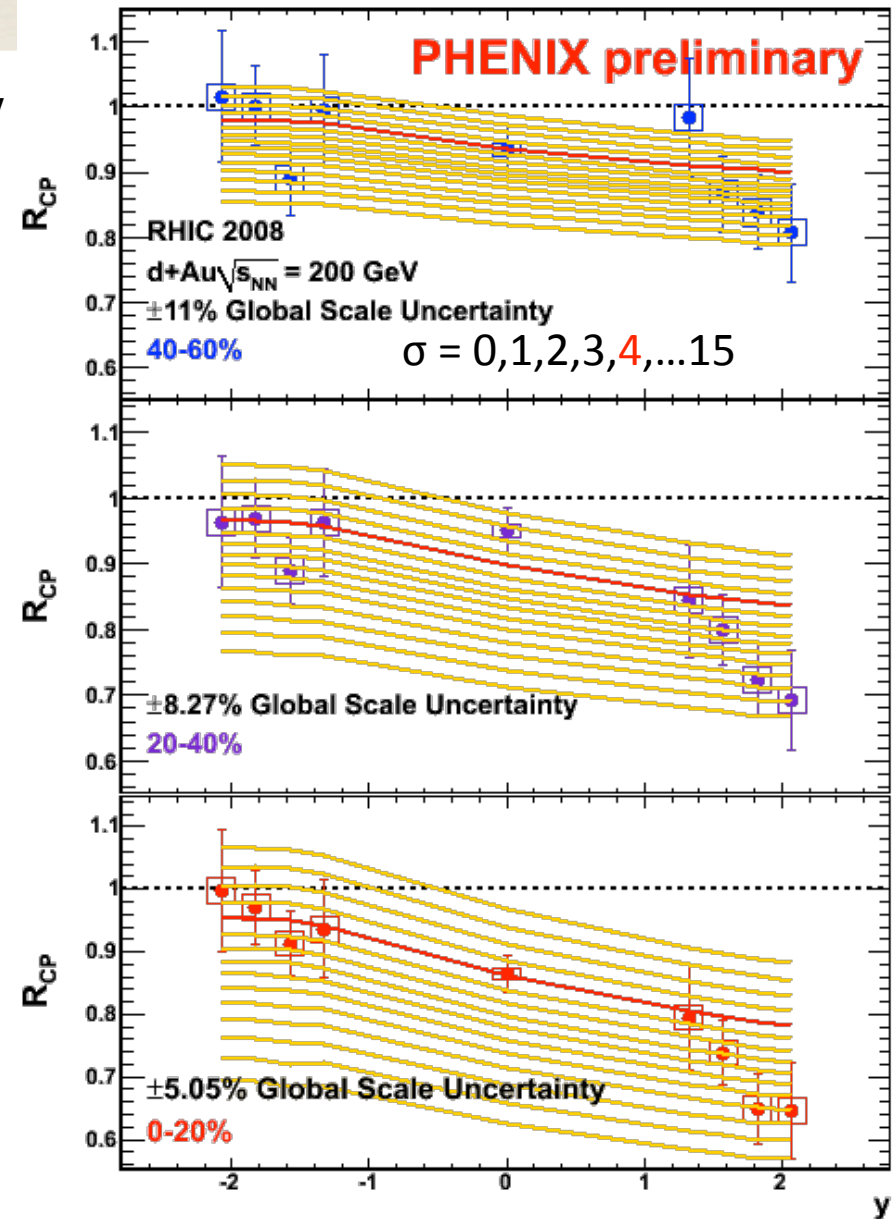
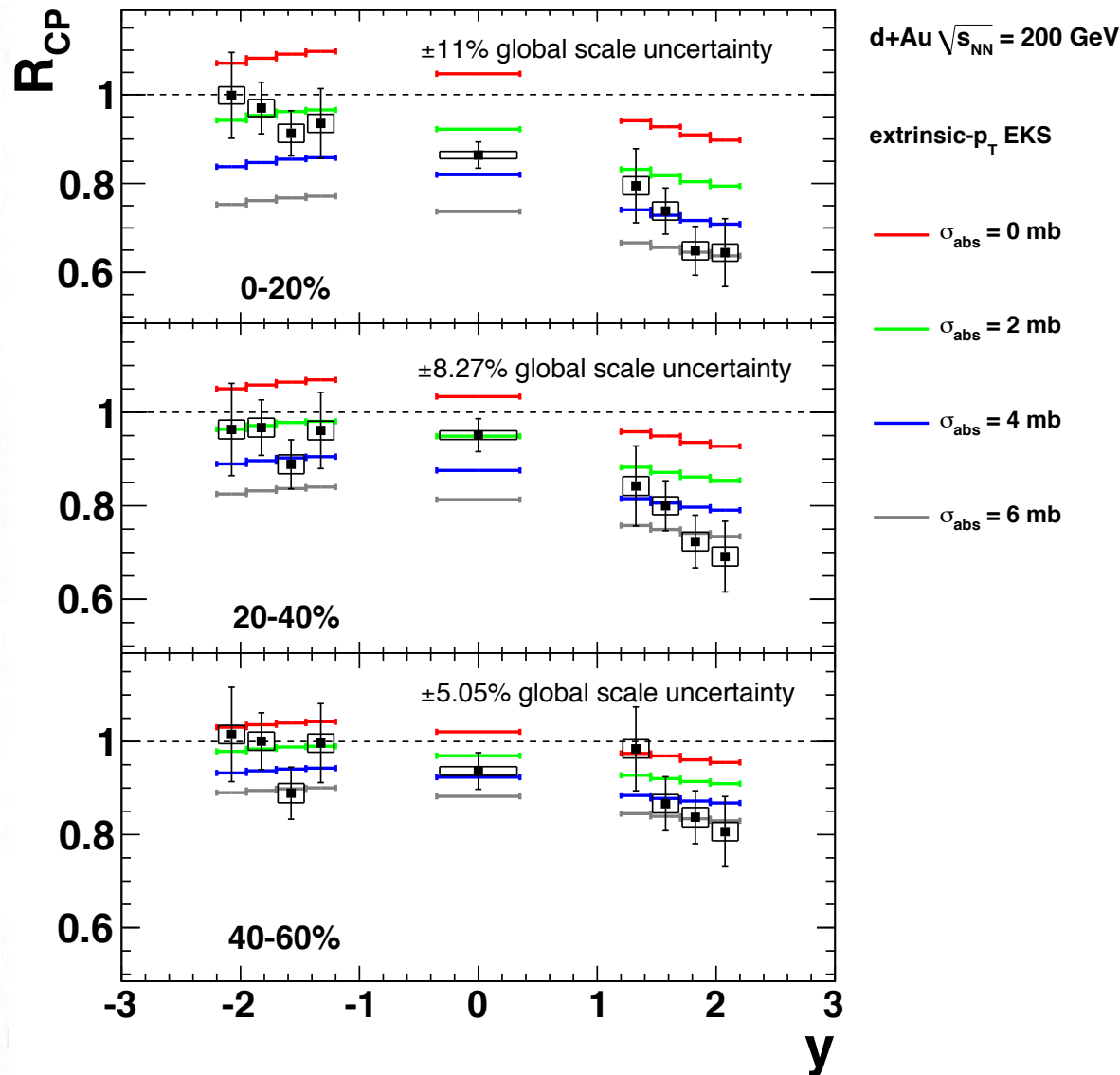
nDSg

# A closer look at $R_{CP}$ vs $y$ in dAu with

$$g + g \rightarrow J/\psi + g$$

EKS98

$$g + g \rightarrow c\bar{c}$$

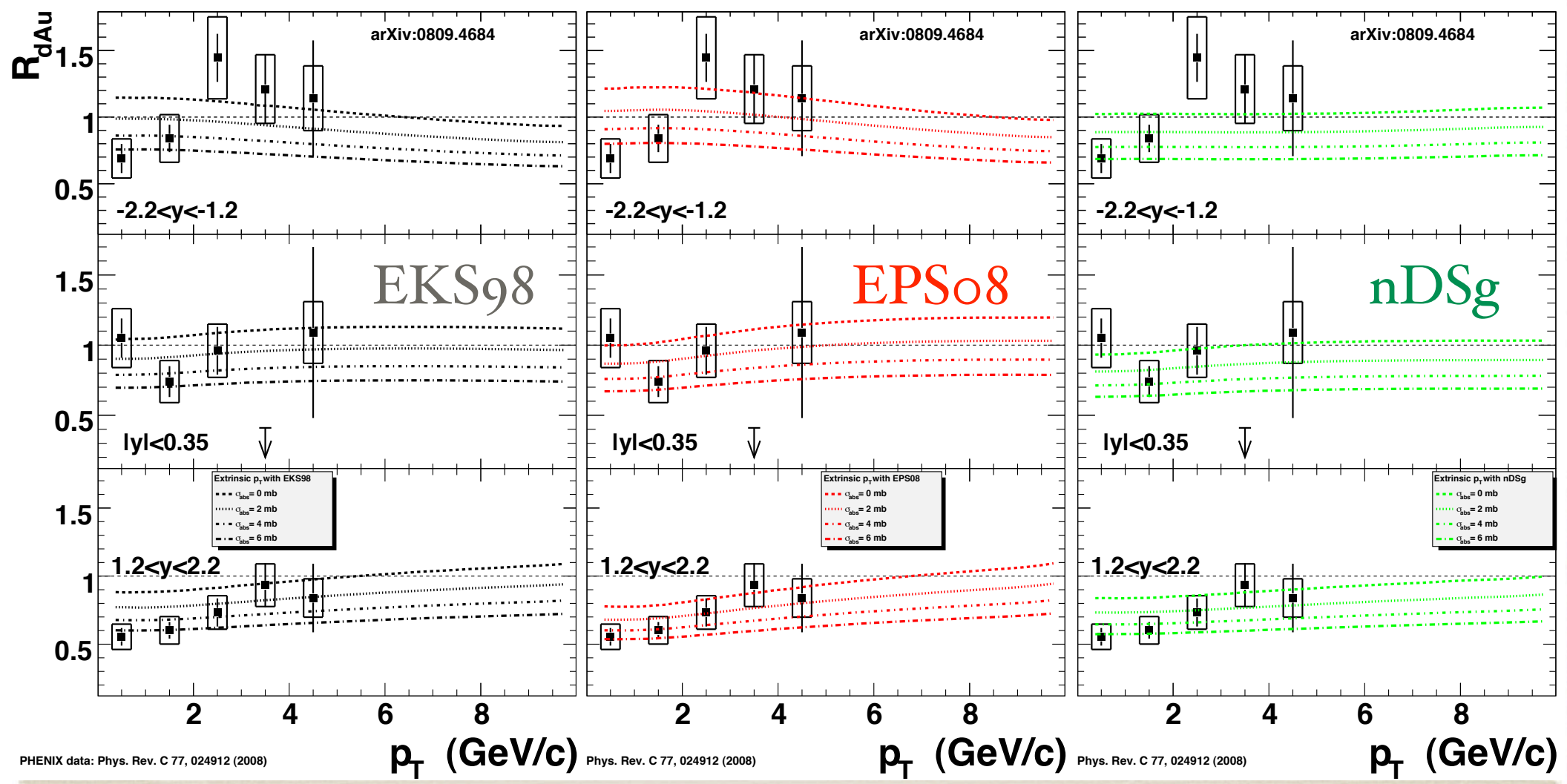


# CNM effects : $J/\psi$ $R_{dAu}$ VS $p_T$

**2 → 2 process**

E. G. Ferreiro, F. Fleuret,  
J. P. Lansberg and A. R.  
PLB 680, 50-55 (2009)

absorption : 0, 2, 4, 6 mb

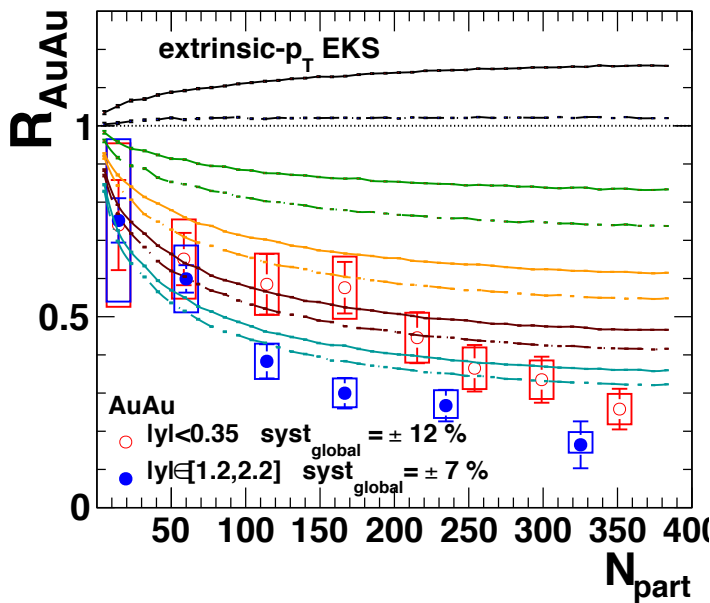


# CNM effects : $J/\psi$ in AuAu

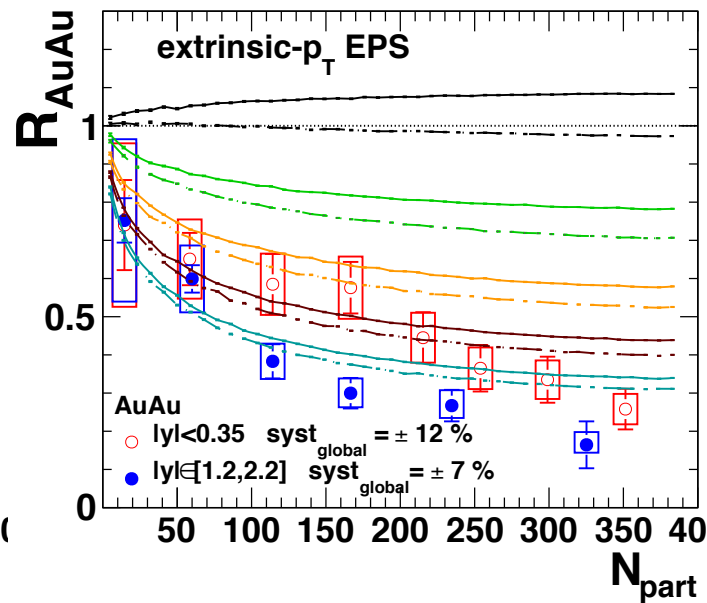
$2 \rightarrow 2$  process

absorption : 0, 2, 4, 6, 8 mb

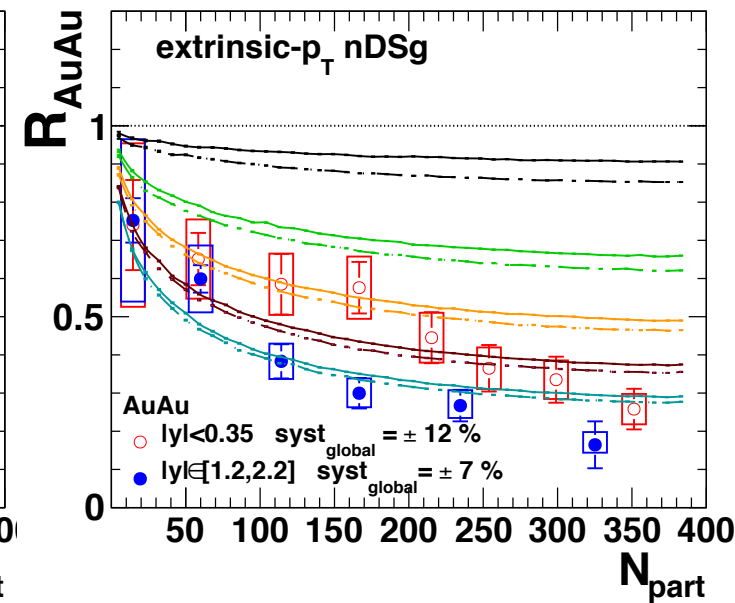
E. G. Ferreiro, F. Fleuret, J. P. Lansberg  
and A. R., PRC 81 (2010) 064911



EKS98



EPS08



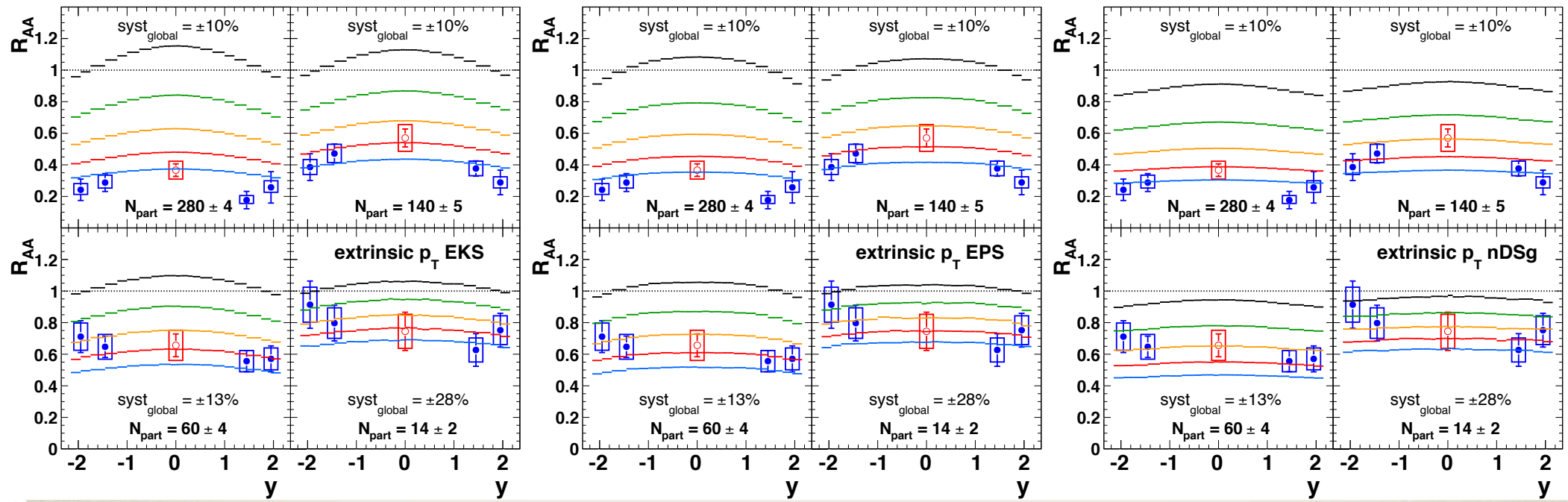
nDSg

# CNM effects : $J/\psi$ in AuAu

$2 \rightarrow 2$  process

absorption : 0, 2, 4, 6, 8 mb

E. G. Ferreiro, F. Fleuret, J. P. Lansberg  
and A. R., PRC 81 (2010) 064911



EKS98

EPS08

nDSg

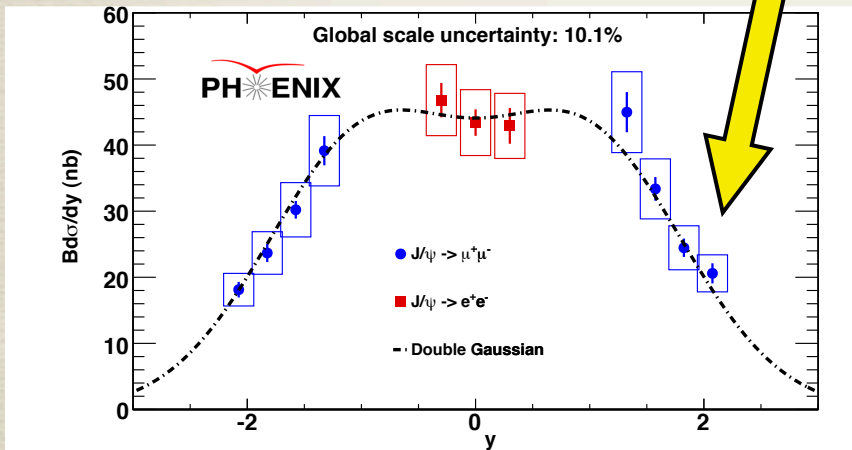
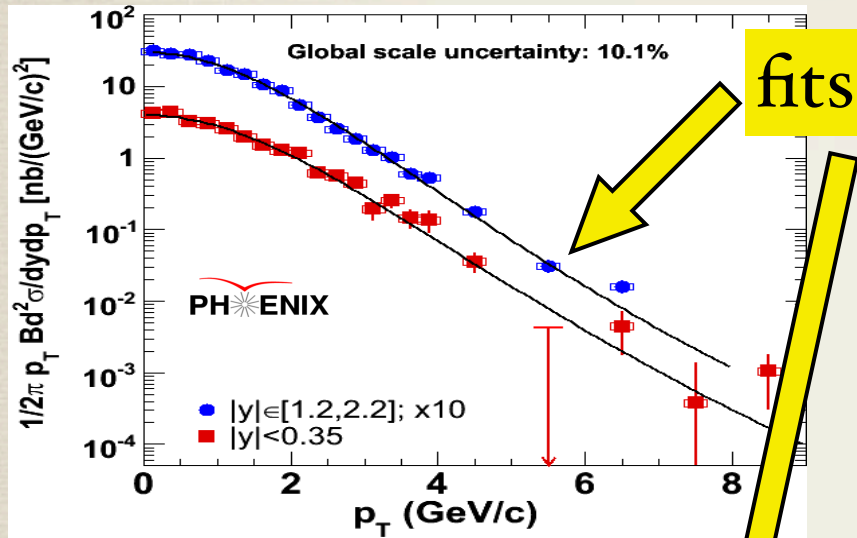
# Adding the $p_T$ dependence

$$(x_1, x_2) \xleftrightarrow[\text{physical constraints}]{c\bar{c} \text{ hard production process}} (y, p_T)$$

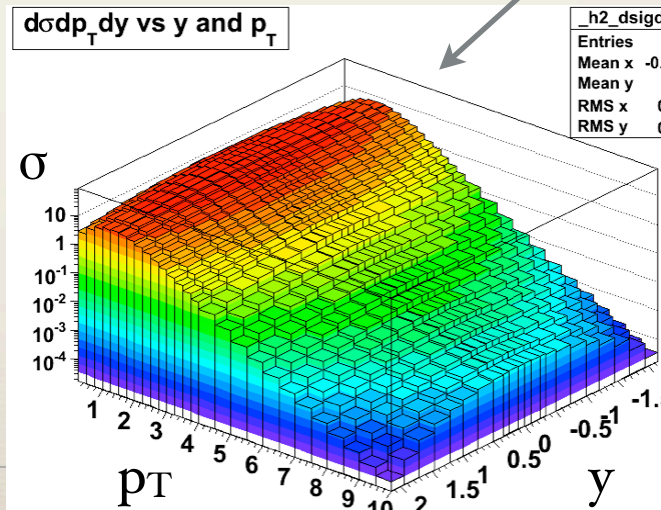
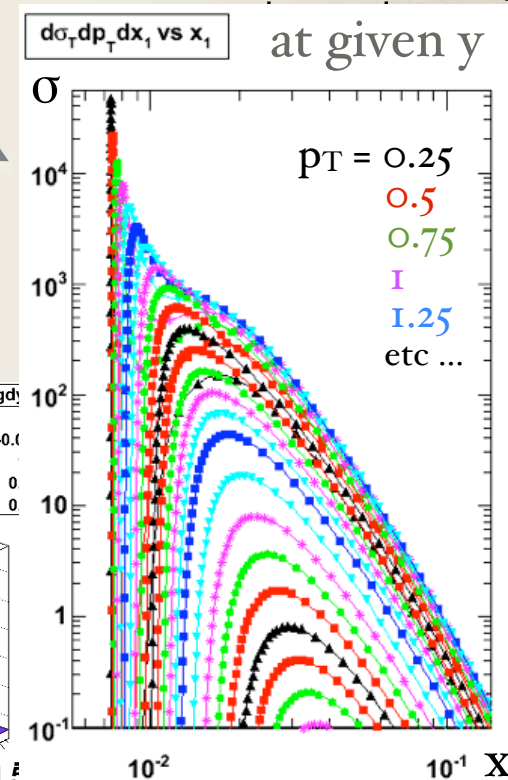
● Intrinsic scheme

● Extrinsic scheme

model



$$x_1 \quad x_2 \quad g(x_1)g(x_2) \frac{d^4\sigma}{dy dp_T dx_1 dx_2} \quad p_T \quad y$$

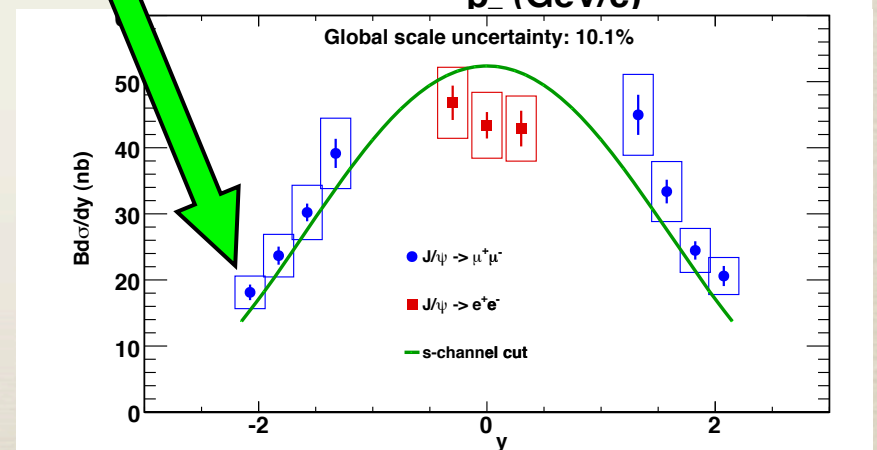
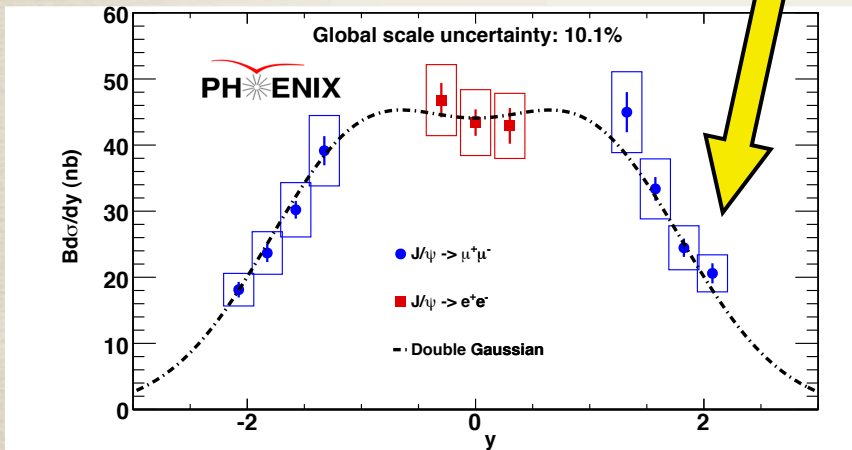
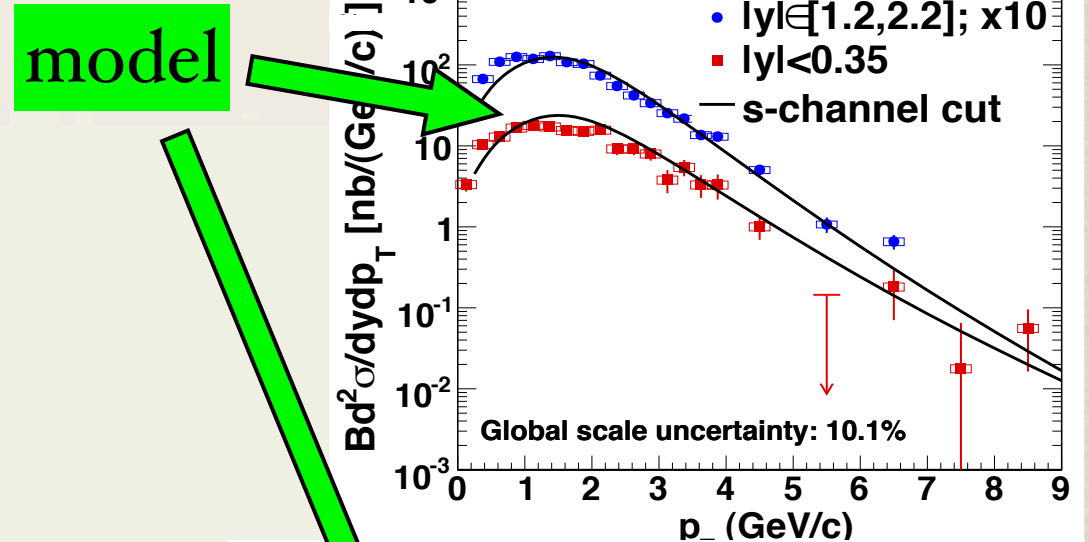
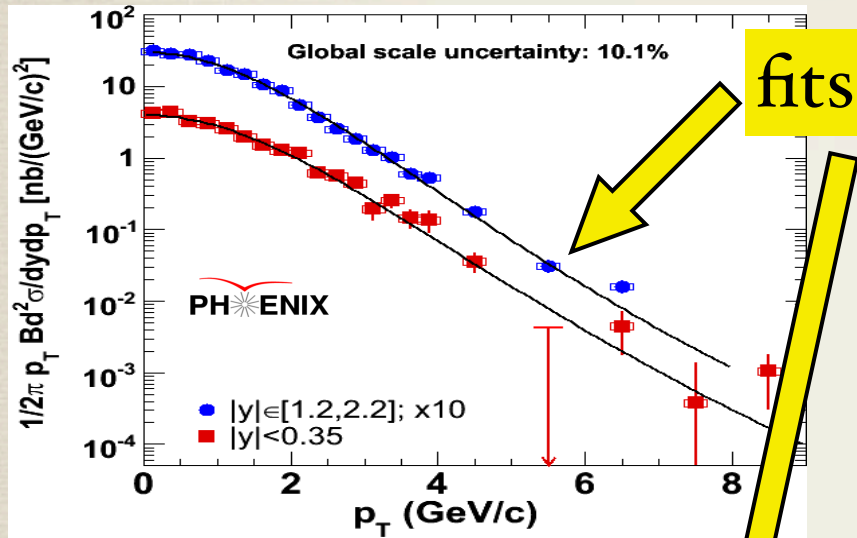


# Adding the $p_T$ dependence

$(x_1, x_2) \xleftrightarrow[\text{physical constraints}]{c\bar{c} \text{ hard production process}} (y, p_T)$

● Intrinsic scheme

● Extrinsic scheme





# Our Monte-Carlo approach for J/ψ production

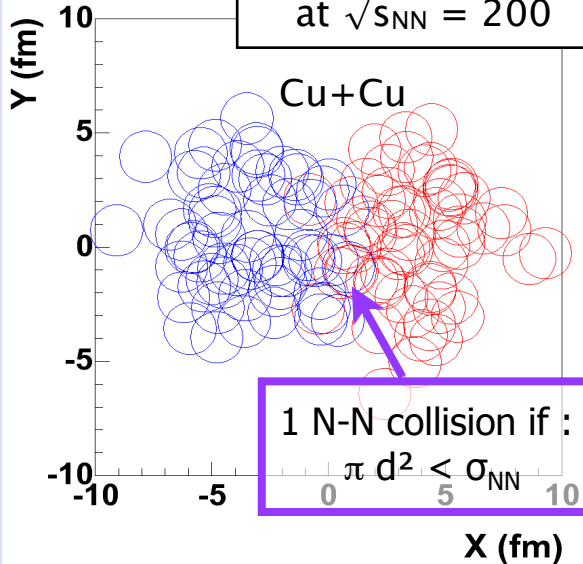
1

2

3

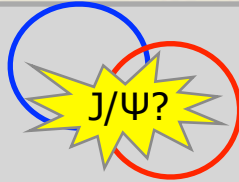
## Glauber MC

$\sigma_{NN} = 42\text{mb}$   
at  $\sqrt{s_{NN}} = 200$



Random :

- $b$  according to  $2\pi b db$
- position of nucleons  $\in A, B$  according to Woods-Saxon



For each N-N collision

## J/ψ candidate produced

- according to  $\sigma_{J/\psi} \leq \sigma_{NN}$

with random :

- $y$  and  $p_T$
- random  $p_T$  orientation  $\phi$  uniformly distributed in  $[0, 2\pi]$
- $x_1, x_2$  determined from intrinsic or extrinsic scheme

## Kinematics for J/ψ candidate:

$y, p_T, \phi, M \Rightarrow p_x, p_y, p_z, E$

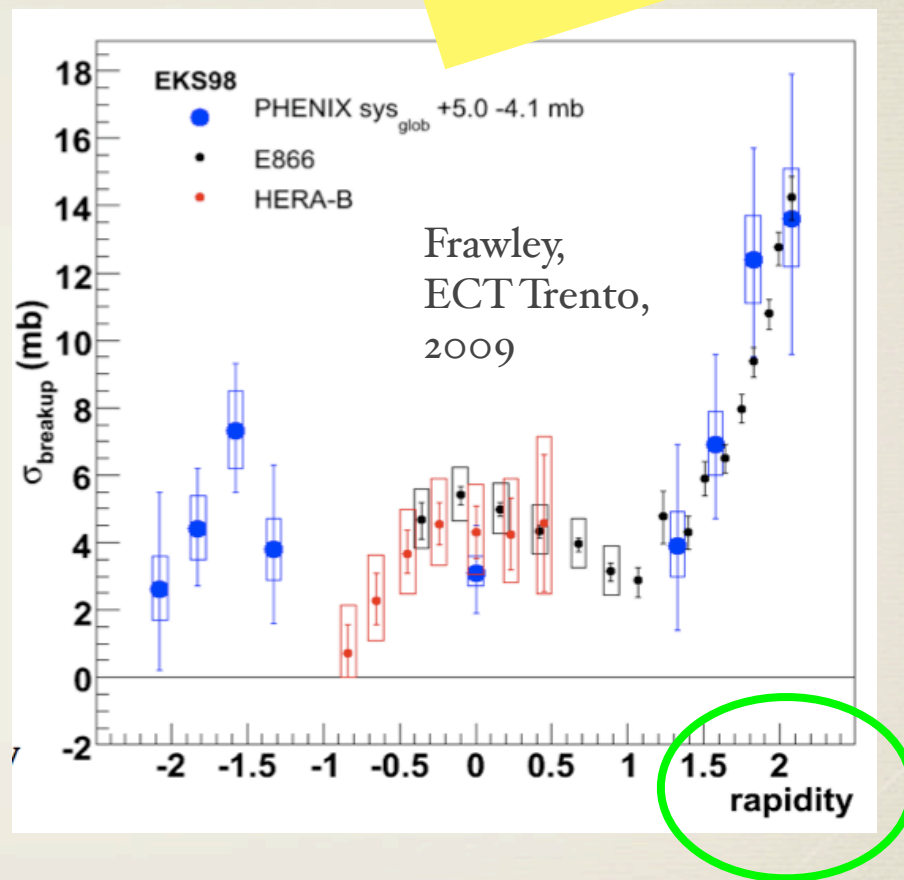
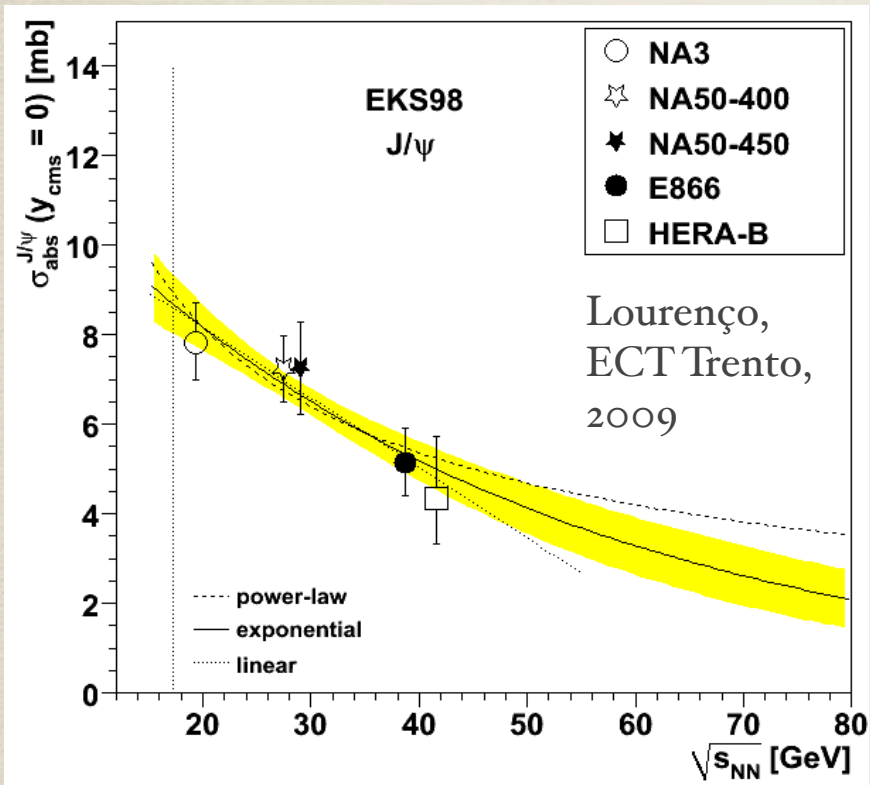
J/ψ candidate  $\Rightarrow$  real J/ψ if :  
 $\text{random}[0,1] < R_{\text{shadow}} \times \sigma_{J/\psi} / \sigma_{NN}$

computed using **EKS**

Nuclear modif. factor =  
 $dN_{\text{real J/}\psi} / dN_{\text{J/}\psi \text{ candidate}}$

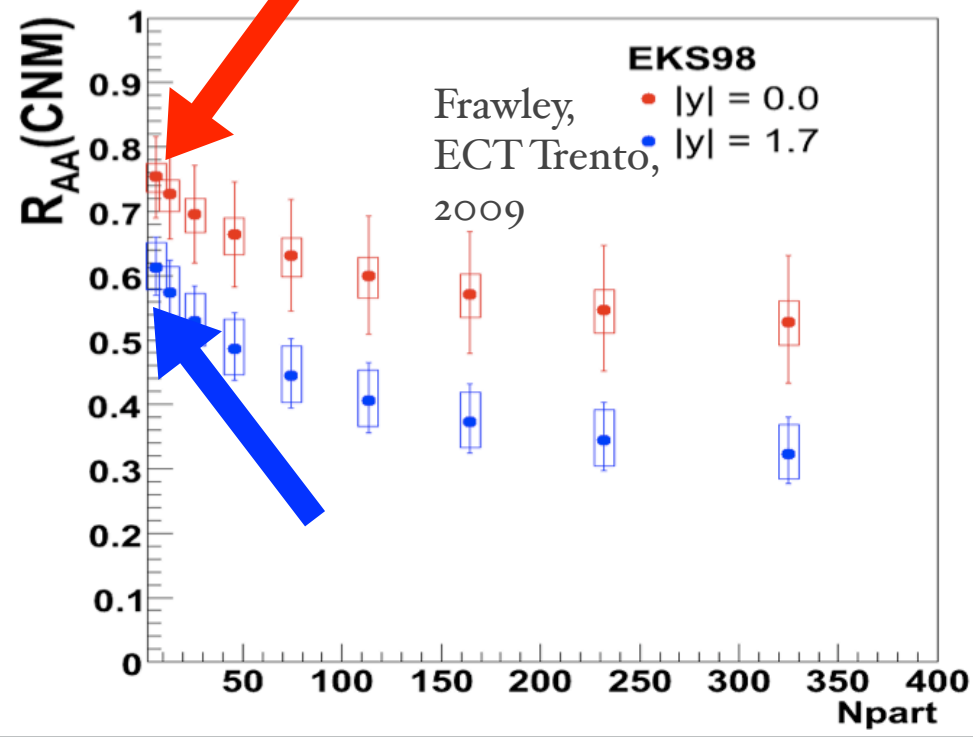
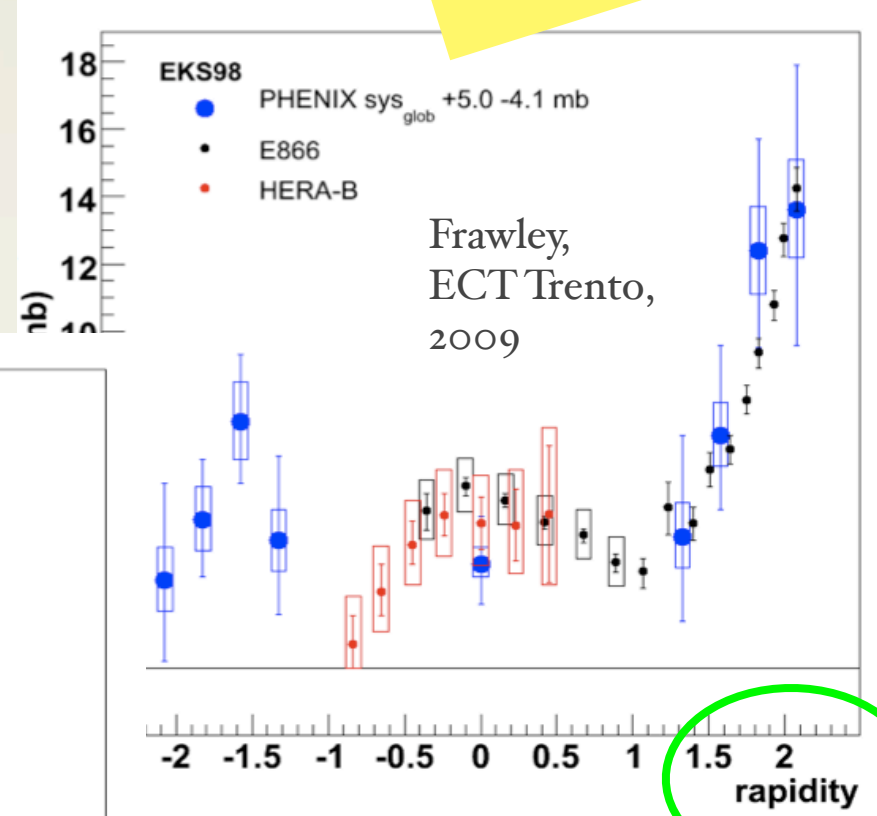
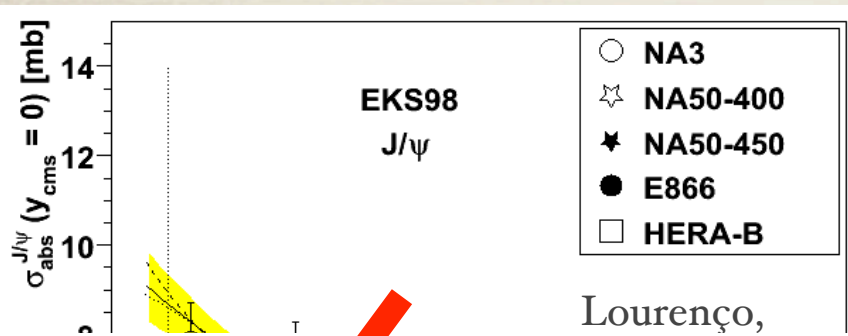
# What about any dependence of $\sigma_{\text{abs}}$ on kinematics ?

investigated assuming  $2 \rightarrow 1$  process only!



# What about any dependence of $\sigma_{\text{abs}}$ on kinematics ?

investigated assuming  $2 \rightarrow 1$  process only!



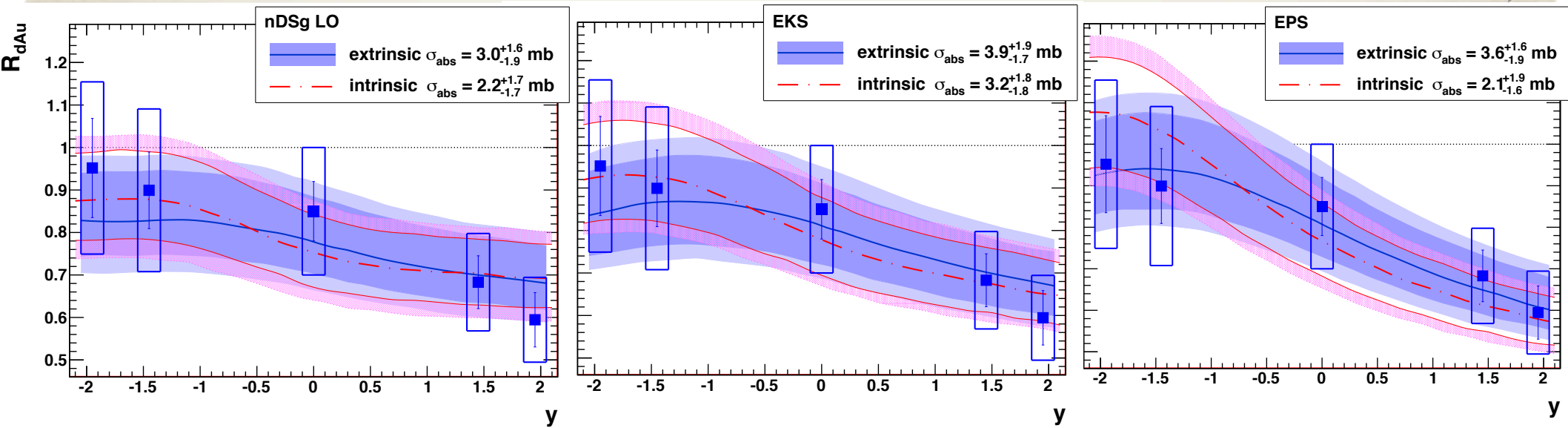
# Data constrains on the value of the effective $\sigma_{\text{abs}}$

A.R., E. G. Ferreiro, F. Fleuret and J. P. Lansberg, arXiv:1002.2351

hard to distinguish  $2 \rightarrow 1$  to  $2 \rightarrow 2$  processes with current exp. err.

$R_{\text{dAu}}$  vs rapidity

nPDF with stronger anti-shadowing 



nDSg

EKS<sub>98</sub>

EPS<sub>08</sub>