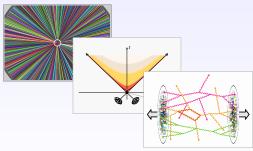
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The Quark-Gluon plasma in the LHC era

Journées de prospective IN2P3-IRFU, Giens, Avril 2012



François Gelis IPhT, Saclay

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Quarks and gluons

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Matter : quarks ; Interaction carriers : gluons

$$\sum_{i}^{a} g(t^{a})_{ij} \qquad \sum_{i}^{a} g(T^{a})_{bc}$$

- *i*, *j* : quark colors ; *a*, *b*, *c* : gluon colors
- $(t^{a})_{ij}$: 3 × 3 SU(3) matrix ; $(T^{a})_{bc}$: 8 × 8 SU(3) matrix

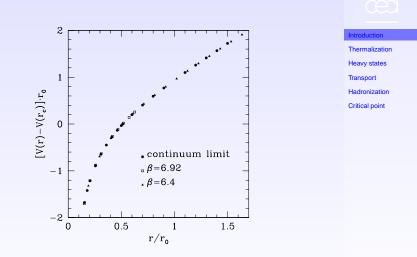
Lagrangian

$$\mathcal{L} = -\frac{1}{4} F^2 + \sum_{f} \overline{\psi}_{f} (i \not\!\!D - m_{f}) \psi_{f}$$

• Free parameters : quark masses m_f , scale Λ_{ocd} (or $\alpha_s(M_z)$)

Color confinement

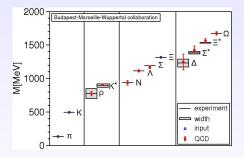
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• Lattice evidence : linear rise of the $Q\overline{Q}$ potential

Color confinement

- In nature, we do not see free quarks and gluons (the closest we have to actual quarks and gluons are jets)
- Instead, we see hadrons (quark-gluon bound states):



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- The hadron spectrum is uniquely given by $\Lambda_{_{
 m QCD}}, m_{f}$
- But this dependence is non-perturbative (it can now be obtained fairly accurately by lattice simulations)

Debye screening

What happens when many hadrons are packed in a small volume?

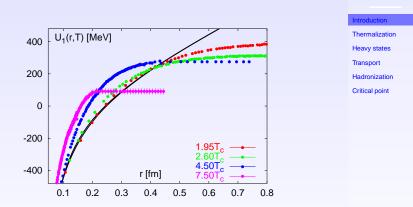
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- In a dense medium, color charges are screened by their neighbors
- The interaction potential decreases exponentially beyond the Debye radius r_{debve}
- Hadrons whose radius is larger than r_{debve} cannot bind

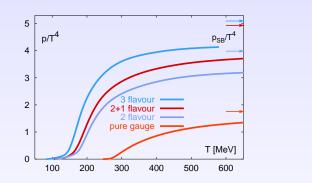
Debye screening

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 In lattice calculations, one sees the QQ potential flatten at long distance as T increases

Deconfinement transition



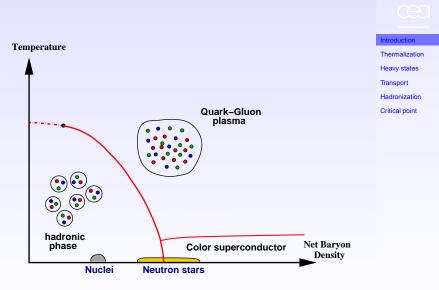
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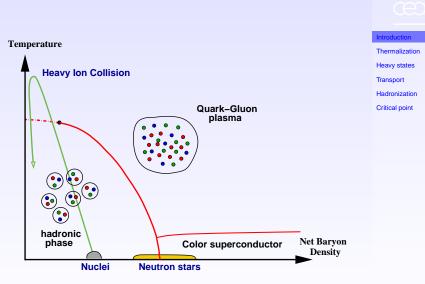
- Rapid increase of the pressure :
 - at T ~ 270 MeV, with gluons only
 - at $T \sim 150$ to 180 MeV, with light quarks

▷ interpreted as the increase in the number of degrees of freedom due to the liberation of quarks and gluons

QCD phase diagram



Heavy ion collisions



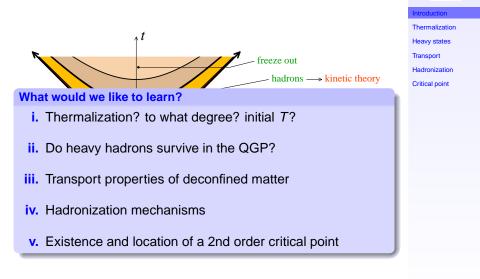
Stages of a nucleus-nucleus collision

ffreeze out hadrons \rightarrow kinetic theory gluons & quarks in eq. \rightarrow ideal hydro gluons & quarks out of eq. \rightarrow viscous hydro strong fields \rightarrow classical dynamics Z

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Stages of a nucleus-nucleus collision



Is the QGP in equilibrium?

Why is it important?

- If the QGP is not in equilibrium, then it remembers many details about the initial condition ⇒ what we measure is not a simple property of the QGP
- Theory predictions are much easier in equilibrium (lattice QCD works ONLY in equilibrium)
- Hydrodynamics is applicable only to systems close enough to equilibrium

Questions

- At what time is the system thermalized?
- What is its temperature at that time?

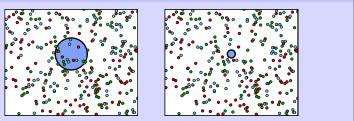
Probes

- Indirect evidence from hydrodynamical models
- Thermal direct photons

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Do heavy states survive?

Large hadrons melt, while the small ones should survive



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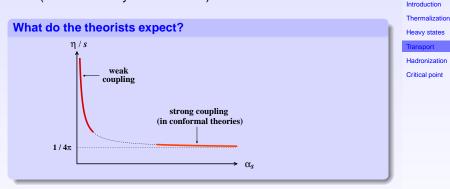
- Bound state sizes decrease as their mass increase
- Excited states are larger than ground states
 - \implies Melting order: light mesons, χ_c , ψ' , J/ψ , Υ'' , Υ' , Υ

Probes

- Quarkonia yields, up to the upsilon family
- Open heavy flavors would help disentangle various models
- · Reference yields in proton-proton and proton-nucleus

Transport properties (I)

 The viscosity η tells us how good a fluid the QGP is (small viscosity = better fluid)

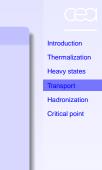


What we would like to learn from flow measurements

- How small is it really?
- So far, only a gross average value: can we get the viscosity at several stages of the evolution?
- Do heavier partons (c and b quarks) flow?

Transport properties (II)

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- Are jets completely absorbed?
- What fraction of hadrons manage to get out?

How are parton cascades modified in a dense medium?

- Modifications of angular correlations
- Modifications of jet fragmentation functions
- Mass dependence: do heavy quarks also lose energy?

Hadronization by fragmentation (dominant in vacuum)

• $V(Q, \overline{Q})$ increases

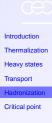


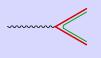
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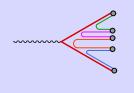
Hadronization by fragmentation (dominant in vacuum)

- $V(Q, \overline{Q})$ increases
- At some point, there is enough energy to extract a new qq pair from the vacuum





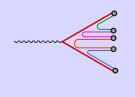
Hadronization by fragmentation (dominant in vacuum)



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- This process repeats, and hadrons are formed at the end of the cascade

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Hadronization by fragmentation (dominant in vacuum)

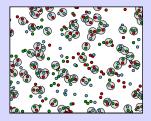


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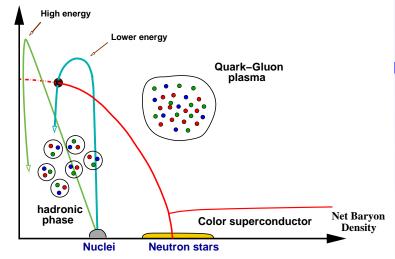
Hadronization by coalescence



- Nearby partons in the QGP bind together
- · Larger yield of baryons
- More strange hadrons if $T \gtrsim 2m_s$

Critical point

Temperature



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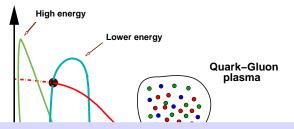
Heavy states

Critical point

Transport Hadronization

Critical point

Temperature



- At zero density: crossover (known from lattice QCD)
- At high density: first order transition line?
- If true, there should be a 2nd order critical endpoint
- How do we get there? collisions at lower energy
- · What do we look for? unusually large fluctuations

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Outlook

CMS Experiment at the LHC, CERN

Data recorded: 2010-Nov-08 10:22:07 828203 GMT(11:22:07 CB Run / Event: 150431 / 541464

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Projects (see J. Castillo and M. Estienne presentations)

- ALICE upgrades: Inner tracking, Calorimetry, Muon forward tracking, Vzero
- CBM experiment at FAIR (GSI, Germany)
- CHIC, AFTER

