



Heavy quarks and QGP with LHCb / SMOG

Physics case

LHCb / SMOG

• Heavy quarks and Quark Gluon Plasma (QGP)

Heavy quarks are "special" QGP probes : $m_Q \gg$ QGP critical temperature T_c (~ 170 MeV),

→ Heavy quarks should be produced in **initial** nucleon-nucleon collisions only, the **QGP phase shouldn't modify the overall heavy quark yields**,

→ **QGP phase should modify relative heavy quark (hidden/open) bound state yields**

Heavy quark hadronization ($c\bar{c}$ example):

- $\sim 90\%$ of $c\bar{c}$ pairs \rightarrow open charm
- $\sim 10\%$ of $c\bar{c}$ pairs \rightarrow hidden charm (charmonia)

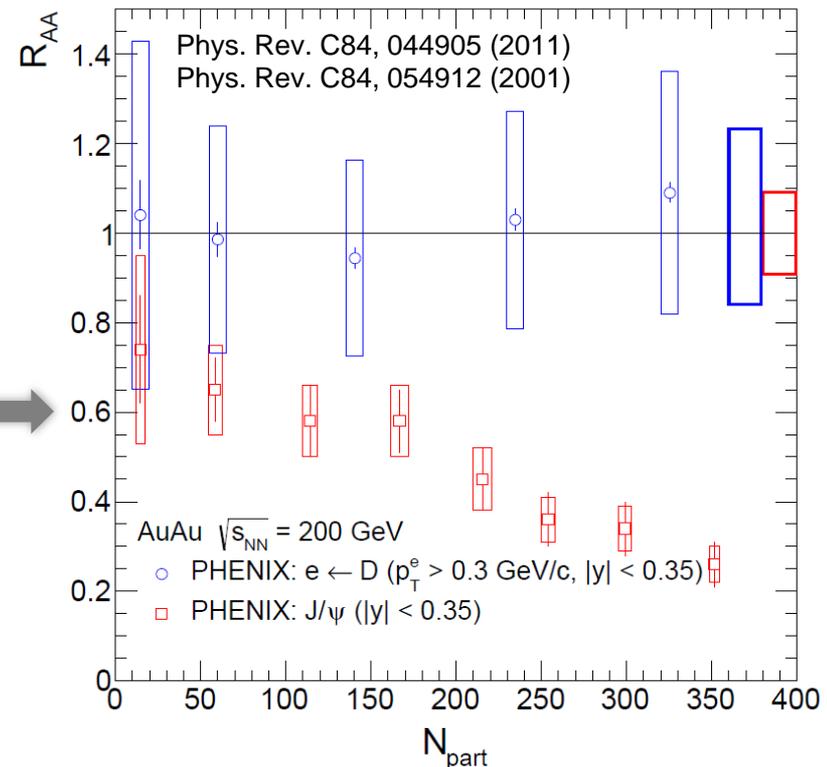
Since most of the produced $c\bar{c}$ pairs hadronize into open charm ($\sim 90\%$), **open charm production reflects the original charm quark yield**.

PHENIX Au+Au collisions @ $\sqrt{s_{NN}} = 200$ GeV

Blue = open charm

Red = hidden charm

- no (little) modification of open charm yield
- modification of J/Ψ ($c\bar{c}$ bound state) yield



- **Heavy quarks and Quark Gluon Plasma (QGP)**

Heavy quarks are "special" QGP probes : $m_Q \gg$ QGP critical temperature T_c (~ 170 MeV),

→ Heavy quarks should be produced in *initial* nucleon-nucleon collisions only, the **QGP phase shouldn't modify the overall heavy quark yields**,

→ *QGP phase should modify relative heavy quark (hidden/open) bound state yields*

- **Modifying hidden bound state yields**

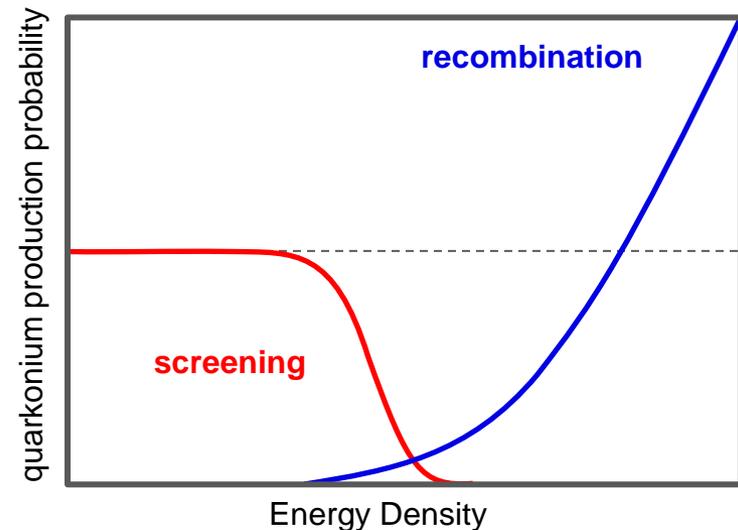
Possible QGP effects on quarkonium:

- **Color screening**: $Q\bar{Q}$ bound states suppression

- Color screening in a QGP decreases quarkonium binding
- Color screening should lead to a suppression of quarkonium production yields

- **Recombination**: $Q\bar{Q}$ bound states enhancement

- at sufficiently high $\sqrt{s_{NN}}$, heavy quarks are abundantly produced.
- After thermalisation, statistical combination can lead to an enhancement of quarkonium production yields



- **Experimentally, charmonium is a privileged probe**

- Charmonium production in A+A collisions studied at:

- | | | |
|------------|------------------------|------------------------------|
| • CERN-SPS | ($\sqrt{s}=17$ GeV) | NA38, NA50, NA60 experiments |
| • BNL-RHIC | ($\sqrt{s}=200$ GeV) | PHENIX, STAR experiments |
| • CERN-LHC | ($\sqrt{s}=2.76$ TeV) | ALICE, CMS experiments |

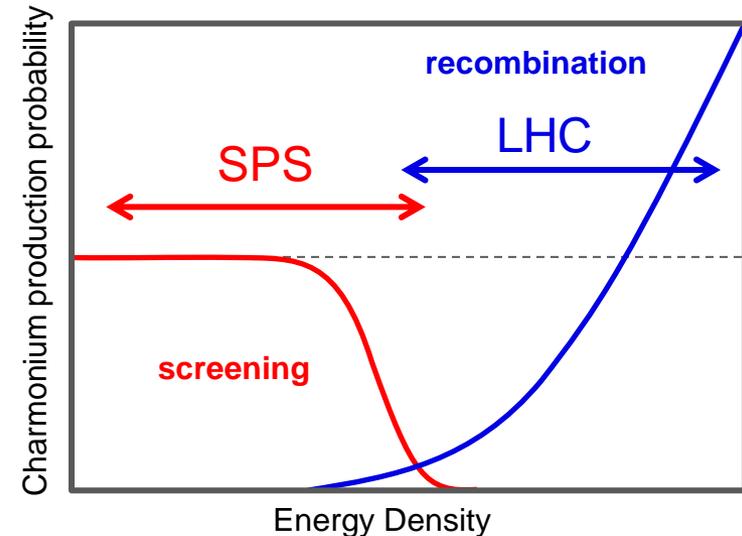
- Short summary for J/Ψ :

- | | |
|----------------------|-----------------------------------------------------|
| • NA50 (PbPb@SPS) | observed an <i>anomalous</i> J/Ψ suppression |
| • PHENIX (AuAu@RHIC) | observed a <i>similar</i> suppression (than NA50) |
| • ALICE (PbPb@LHC) | observed a <i>smaller</i> suppression (than PHENIX) |

➔ Possible Color screening starting at SPS

➔ Possible recombination occurring at LHC

- **Within the SPS+RHIC+LHC energy range, charm seems to be the adequate probe to investigate both screening and recombination.**



• What next to be done with charmonium

To confirm (and study) charmonium color screening and recombination, one must compare charmonium and open charm production in A+A collisions

- Since most of the produced $c\bar{c}$ pairs hadronize into open charm ($\sim 90\%$), open charm production reflects the original $c\bar{c}$ pair production
- Open charm is therefore an (the?) appropriate reference to calibrate charmonium screening/recombination studies.

– TeV scale: Charmonium recombination

- Both J/Ψ and open charm will be measured in PbPb at large energy densities at LHC

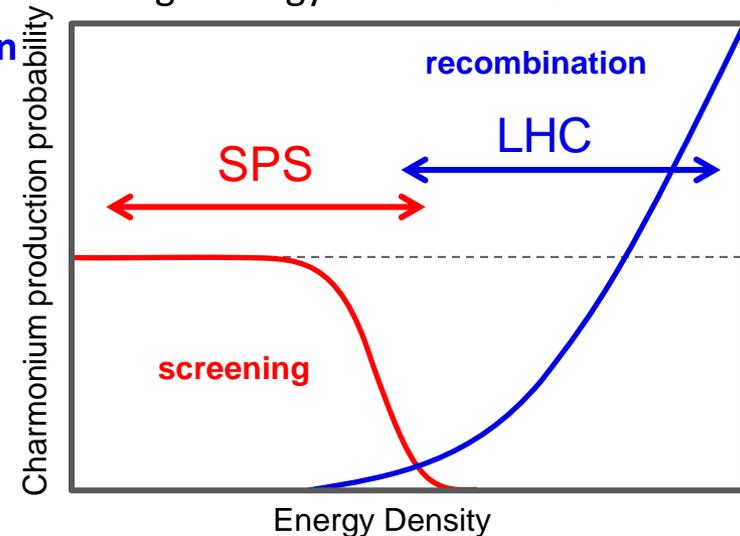
→ Best place to study charmonium recombination

– 20-GeV Scale: Charmonium color screening

- At SPS energies, in Pb+Pb collisions, J/Ψ suppression occurs in the middle of the accessible energy density range

→ Best place to study color screening

- Needs measurement of open charm yields
- Needs precise measurements of several $c\bar{c}$ states to test if color screening leads indeed to a sequential suppression



- **Quarkonium sequential suppression**

- **Quarkonium sequential suppression** in a Quark Gluon Plasma is a prediction of **lattice QCD** :

[H. Satz, J. Phys. G 32 \(2006\)](#)

quarkonium dissociation temperature
critical QGP temperature →

state	$J/\psi(1S)$	$\chi_c(1P)$	$\psi'(2S)$	$\Upsilon(1S)$	$\chi_b(1P)$	$\Upsilon(2S)$	$\chi_b(2P)$	$\Upsilon(3S)$
T_d/T_c	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

- Because of feed-downs and different T_d , **sequential suppression** should show up.

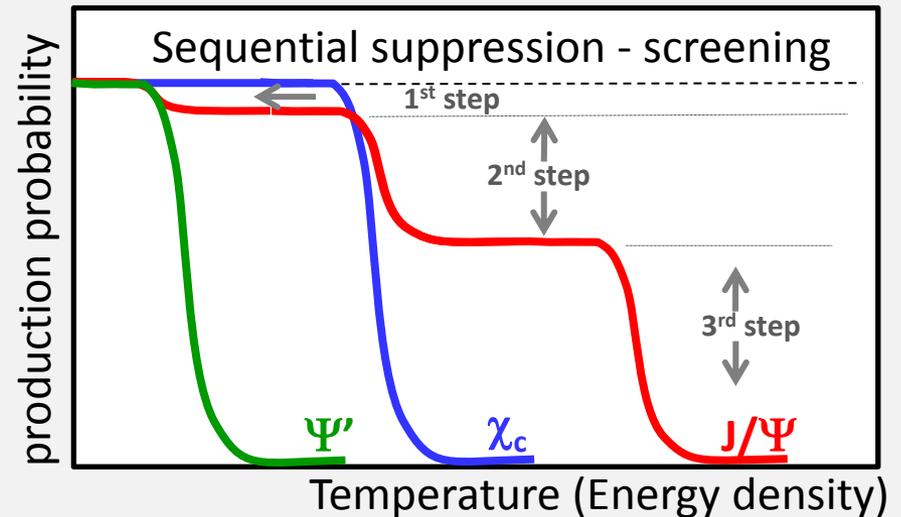
Feed-downs contributing to J/Ψ inclusive yield

60% direct J/Ψ
+ 30% $\chi_c \rightarrow J/\Psi + \gamma$
+ 10% $\Psi' \rightarrow J/\Psi + X$

Inclusive J/Ψ yield

According to lattice calculations,
 $T_d(\Psi') < T_d(\chi_c) < T_d(J/\Psi)$

→ **One should observe a step-like suppression pattern**



- Anomalous suppression at SPS**

[Eur.Phys.J.C49:559-567,2007](http://www.eur-phys.org/ViewArticle.aspx?ajph=1007)

L = length of nuclear matter seen by quarkonium state

Expected = measured yields in p+A extrapolated to large L

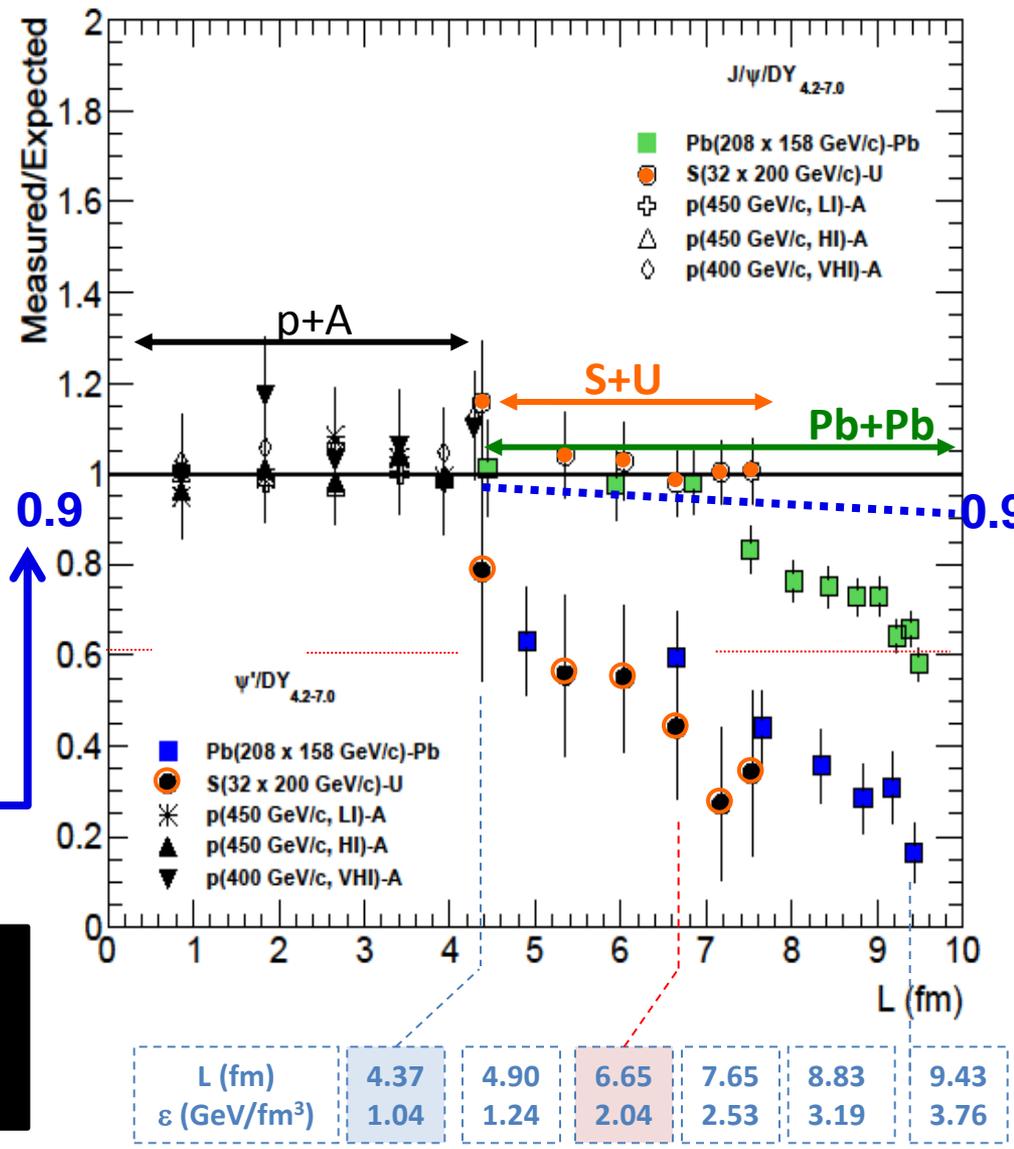
Color screening ?

NA50 measured J/Ψ and Ψ' , but,

- too small $\Psi' \rightarrow J/\Psi$ feed-down
- too fragile Ψ'

to answer the question

→ need of a larger feed-down fraction
 → Need of a stronger bound state
 → Need to measure χ_c yield !



- Anomalous suppression at SPS [Eur.Phys.J.C49:559-567,2007](http://eur.phys.j.c49:559-567,2007)

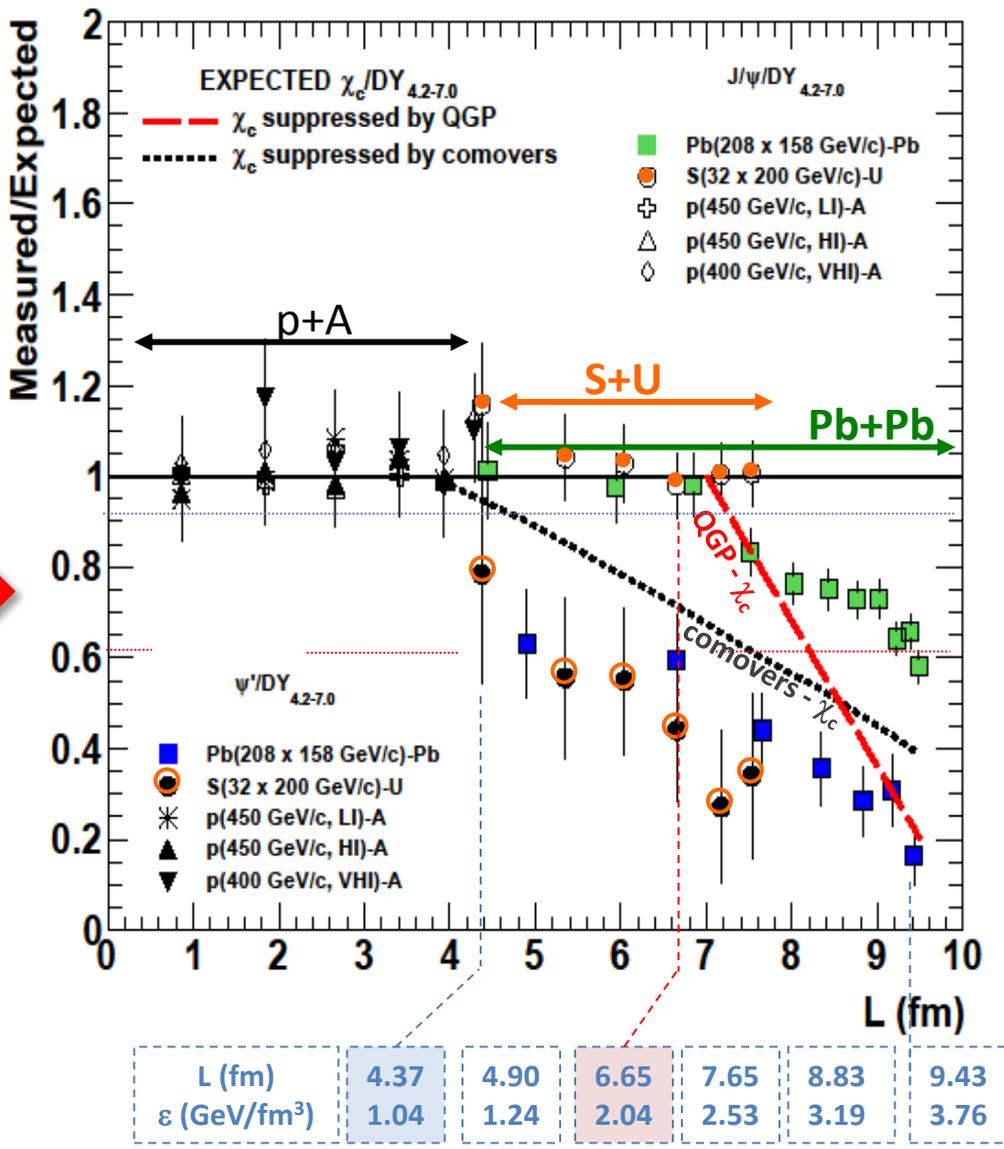
Color screening ?

Take advantage of large $\chi_c \rightarrow J/\Psi$ feed-down fraction

60% direct J/Ψ
 + 30% $\chi_c \rightarrow J/\Psi + \gamma$
 + 10% $\Psi' \rightarrow J/\Psi + X$
Inclusive J/Ψ yield

Measuring J/Ψ , Ψ' and χ_c suppression patterns will give the answer

- **Alternative (no QGP) scenario: suppression by comoving hadrons**
 - Smooth suppression
 - Same suppression-starting point
 - Slopes related to binding energy : $S_{\Psi'} > S_{\chi} > S_{J/\Psi}$



- **Hint of color screening at SPS**
 - Anomalous suppression of J/ψ
 - Anomalous suppression of ψ'
- **But, need to be confirmed**
 - by measuring together J/ψ , ψ' , χ_c and open charm
- **Measuring χ_c may also permit to**
 - Constrain J/ψ T_c
 - test the phase transition

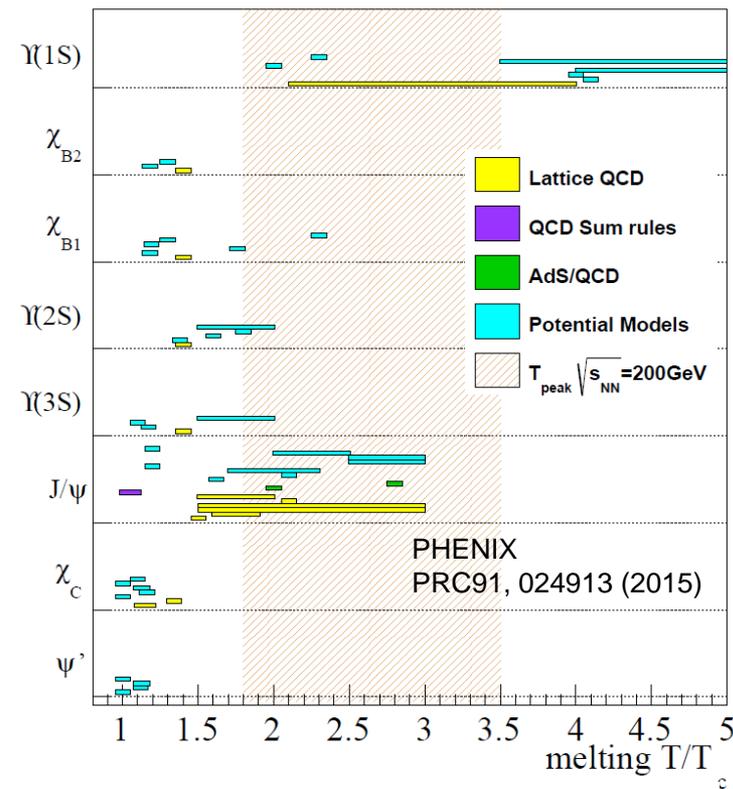


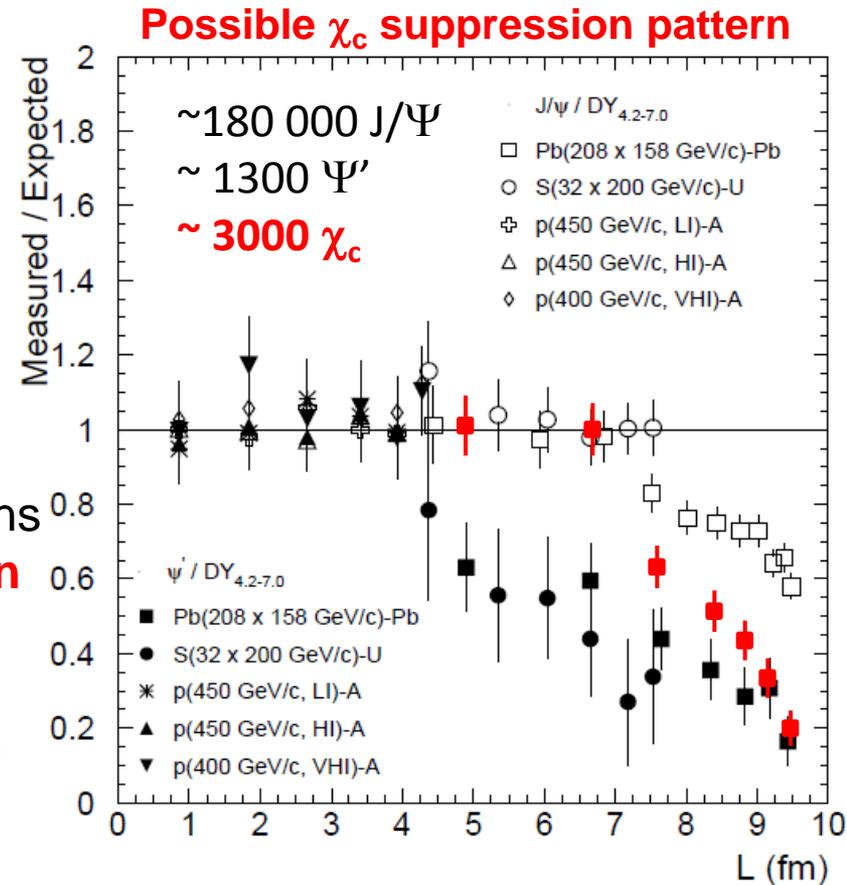
FIG. 1. (Color online) Compilation of medium temperatures relative to the critical temperature (T_e) where quarkonium states are dissociated in the quark-gluon plasma. Note that these estimations were performed assuming different T_e values. Each horizontal bar corresponds to one estimation and its temperature extension (when applied) represents the range where the quarkonia state undergoes a mass/size modification until it completely melts. Techniques used in calculations: Lattice QCD [5–15], QCD sum rules [4, 16–20], AdS/QCD [21–24], effective field theories [27, 28] and potential models [15, 29–35]. The shaded band from 1.8 to 3.5 T/T_e represents the hydrodynamic estimation for the peak temperature reached in Au+Au collisions at 200 GeV [36].

- **Typical 40-day Pb+Pb run @SPS** ($10^7 \cdot s^{-1}$ Pb beam \rightarrow 10% λ_1 Pb target)

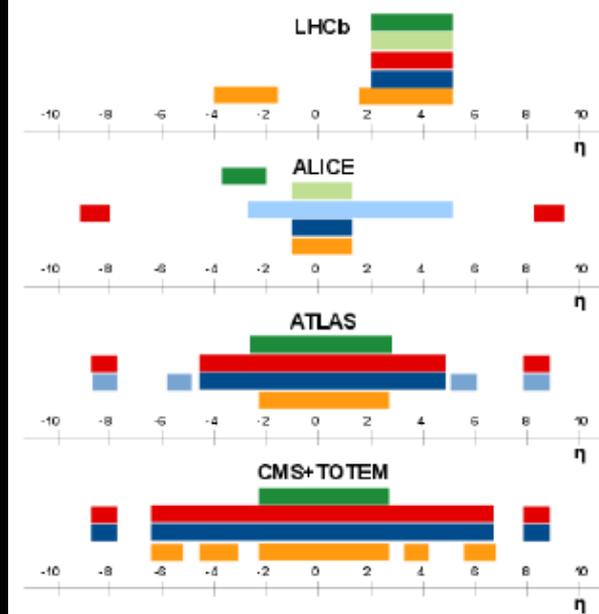
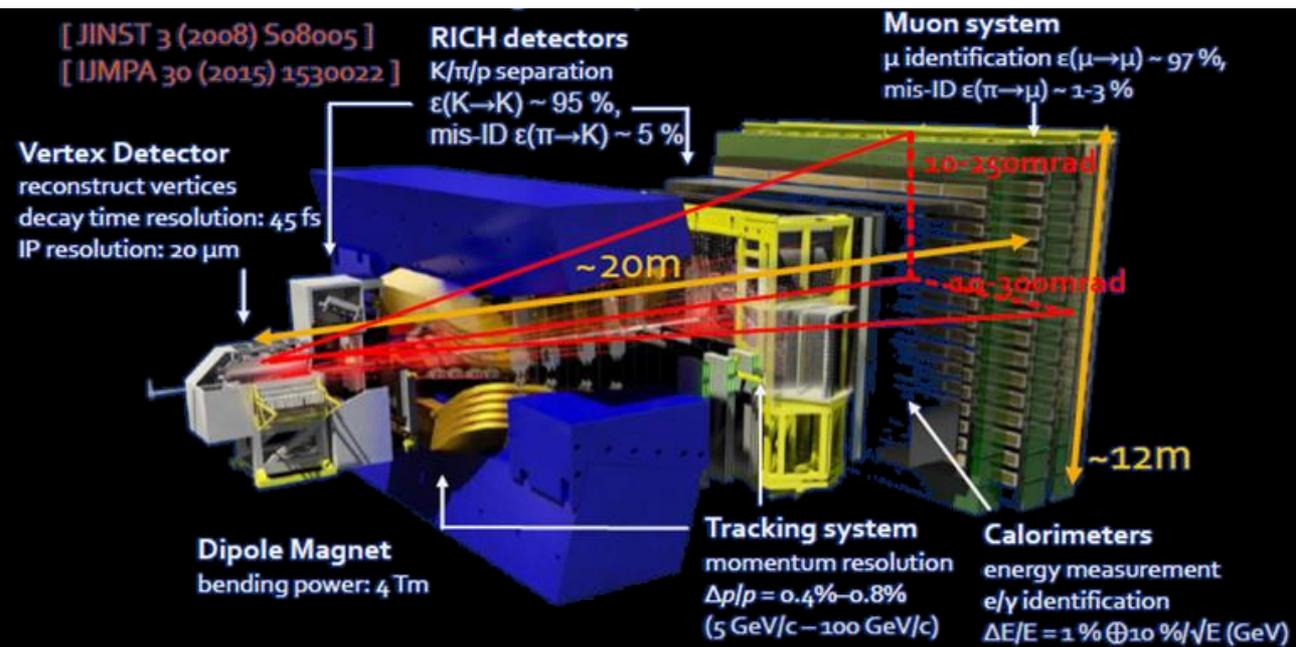
- $\sim 180\,000$ $J/\Psi \rightarrow \mu^+\mu^-$ recorded by NA50
- ~ 1300 $\Psi' \rightarrow \mu^+\mu^-$ recorded by NA50
- If capable of measuring $\chi_c \rightarrow J/\Psi\gamma$,
in NA50 acceptance : $3000 < \chi_c < 7000$

But

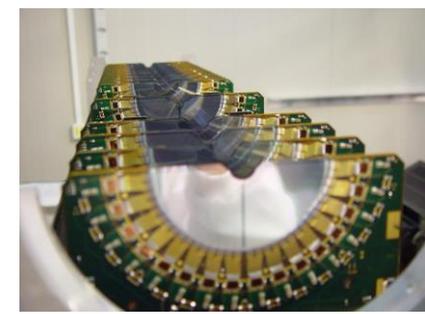
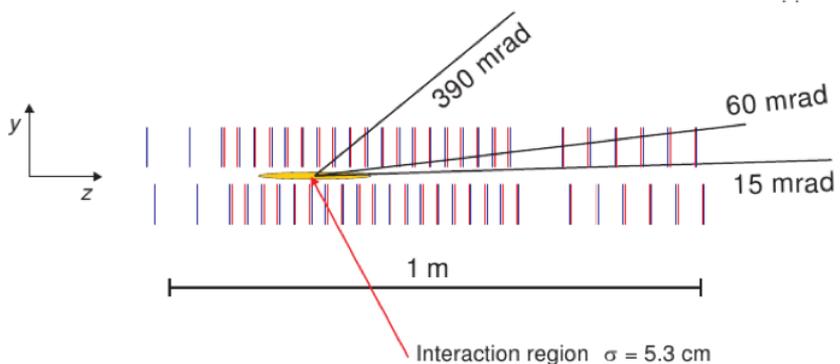
- NA50 was not instrumented to measure photons
 \rightarrow **not able to measure χ_c production**
- NA50 was not instrumented to perform PID
 \rightarrow **not able to measure D^0 production**

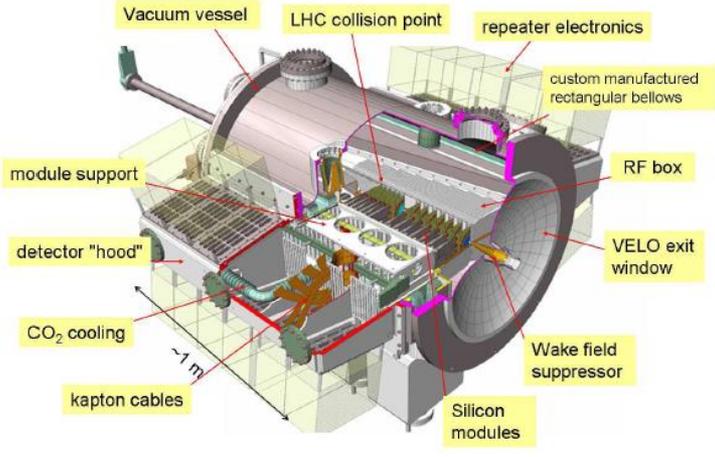
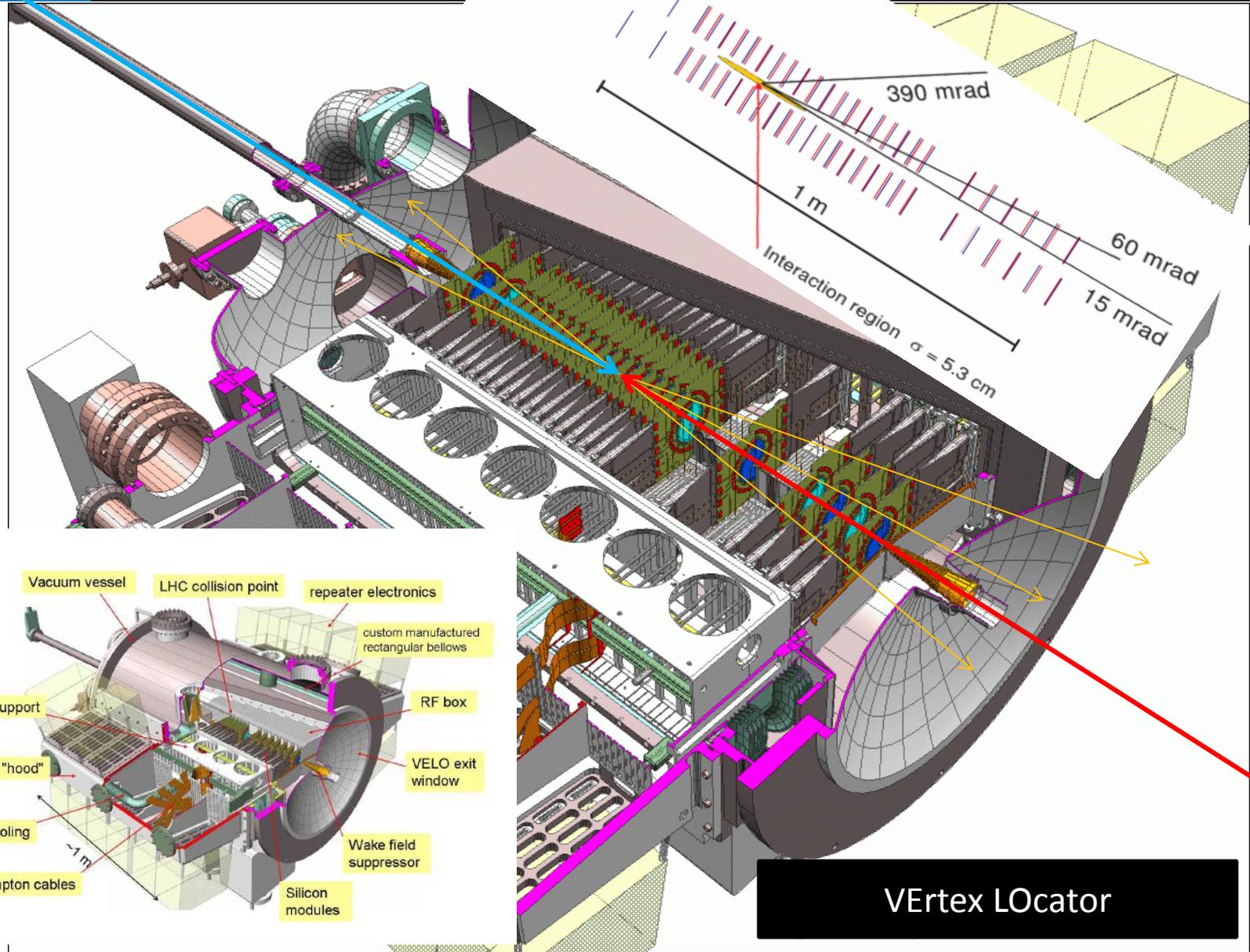


Today, these measurements are possible with LHCb / SMOG



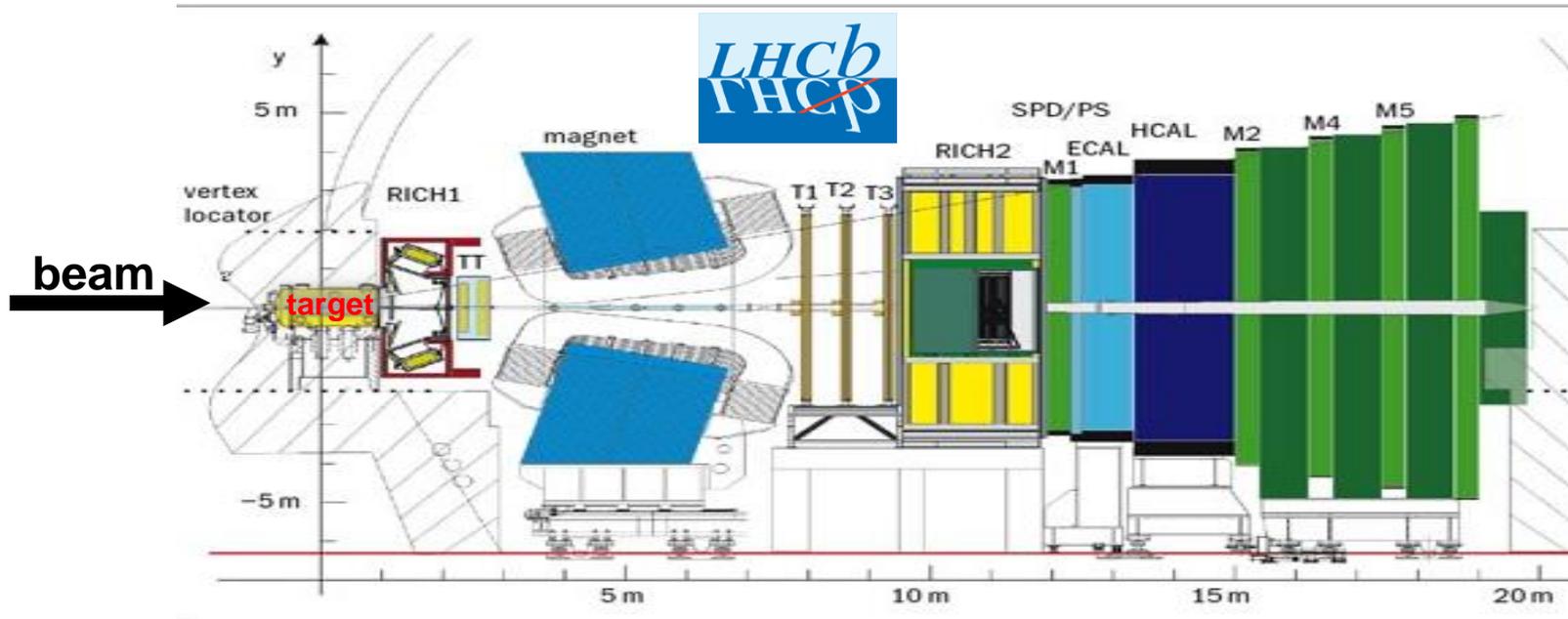
Le VELO (Vertex LOcator)





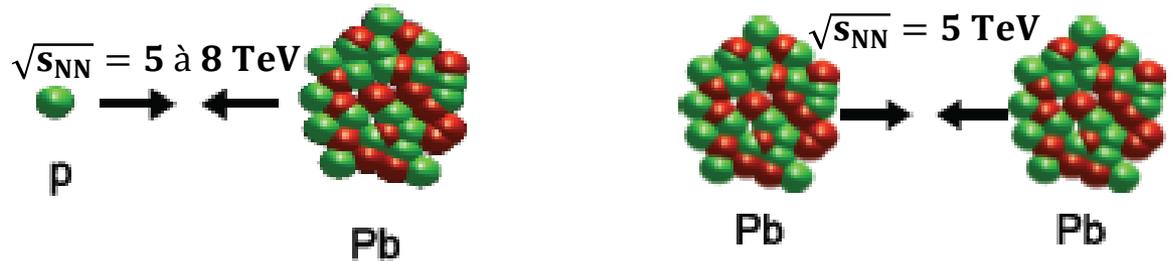
VERtEX LOcator

- Injecting gas in LHCb Vertex Locator (VELO) region
 - Primary role : luminosity measurement
 - Can be used as an internal gas target
 - Noble gas only : (very low chemical reactivity)
 - He (4), Ne (20), Ar (40), Kr (84), Xe (131)
 - Gaz pressure : 10^{-7} to 10^{-6} mbar



- LHCb fonctionnera en deux modes

- Mode collisionneur



- Mode « cible fixe »

$\sqrt{s_{NN}}^{SPS} \sim 20 \text{ GeV}$

$\sqrt{s_{NN}}^{RHIC} = 200 \text{ GeV}$

$\sqrt{s_{NN}}^{LHC} = 5 \text{ TeV}$

$\sqrt{s_{NN}} = 90 \text{ à } 110 \text{ GeV}$



$$\text{LHCb rapidity } 2.5 < y_{\text{LHCb}} < 4.5 \Rightarrow \begin{cases} 7 \text{ TeV beam:} & -2.3 < y_{\text{LHCb}}^* < -0.3 \\ 2.75 \text{ TeV beam:} & -1.8 < y_{\text{LHCb}}^* < 0.2 \end{cases}$$

• PbAr@70 GeV .vs. PbPb@17 GeV

- Multiplicity is related to event centrality and center-of-mass energy
- Multiplicity can be used to compare different A+B collisions at different $\sqrt{s_{NN}}$

System \ centrality	60 – 100%	50 – 60%	40 – 50%	30 – 40%	20 – 30%	10 – 20 %	0 – 10%
PbNe – 71 GeV	108.6	254.4	392.5	588.0	814.5	1086.0	1494.9
PbAr – 71 GeV	123,6	308,8	496,5	806,6	1228,3	1711,9	2372,7
PbKr – 71 GeV	196,9	533,6	919,1	1451,2	2205,5	2986,6	4084,3
PbPb – 17 GeV	124,2	331,6	605,9	919,6	1338,7	2035,8	2980,5

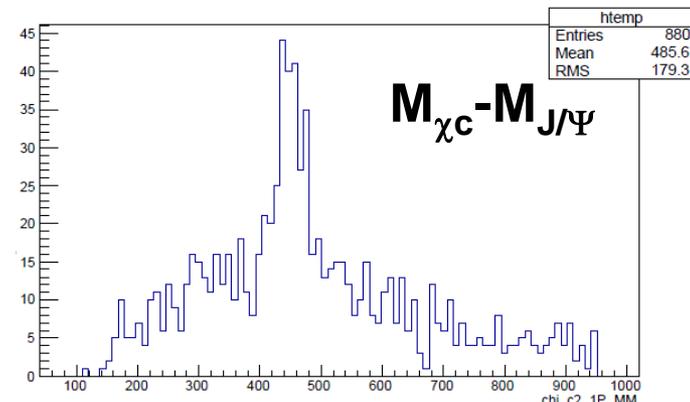
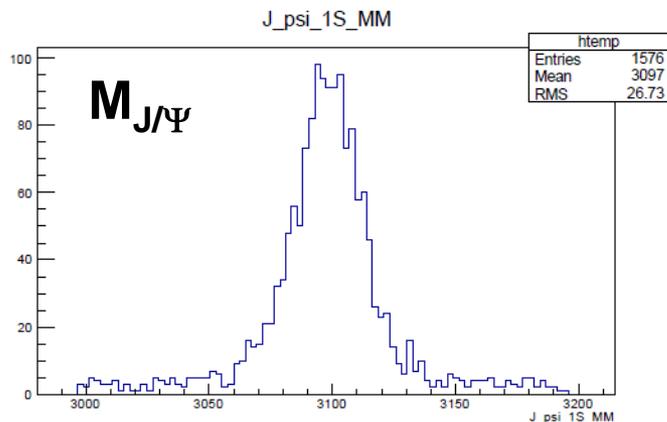
(based on EPOS-LHC-v3400)

- PbAr @ 71 GeV multiplicity \equiv PbPb@17 GeV multiplicity

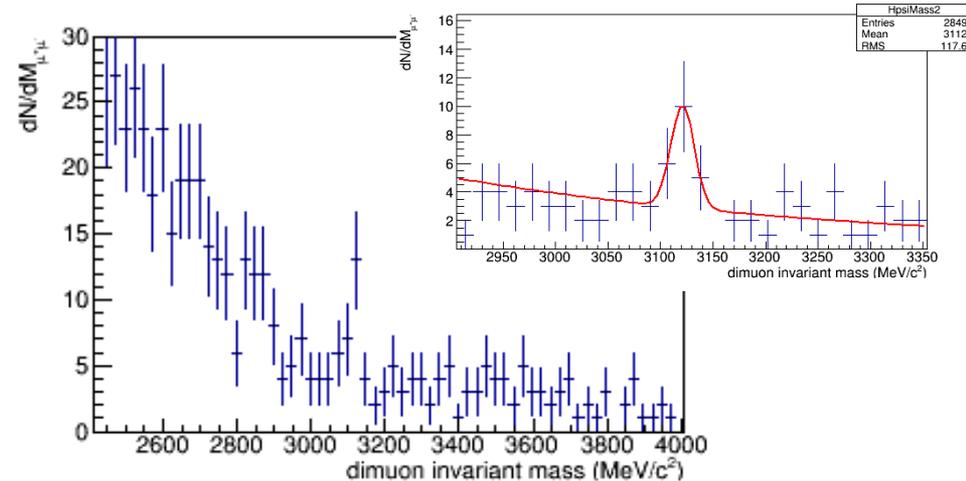
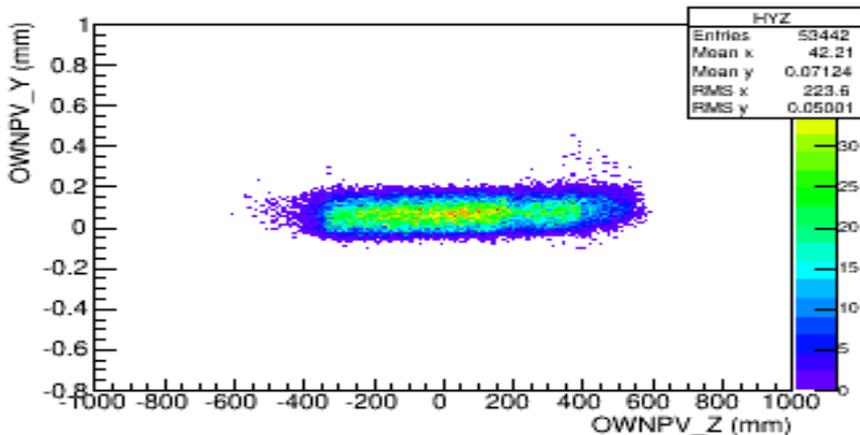
→ PbAr @ 71 GeV is a good starting point to compare with NA50

Acceptance \times efficiency
(PbAr full LHCb simulations)

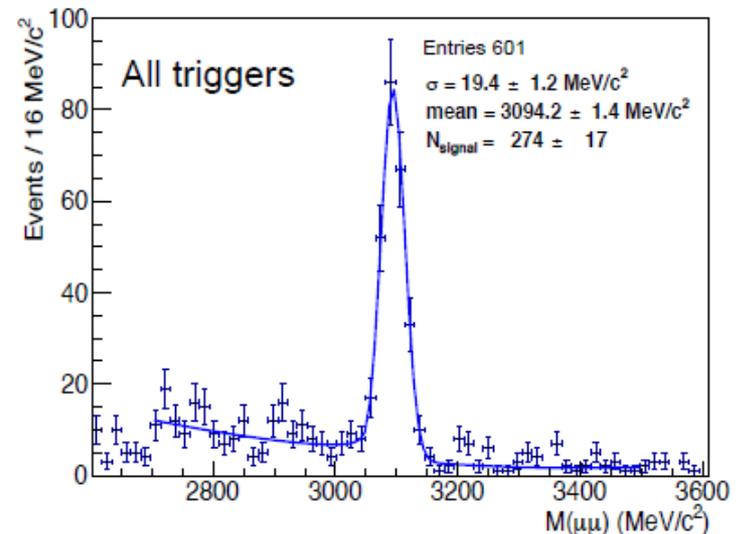
$J/\Psi \sim 20\%$, $\chi_c(J/\Psi \gamma) \sim 5\%$



- Feb. 2013 : 27 minutes Pb+Ne @ 54 GeV



- August 2015 : 12h p-Ne @ ~110 GeV
 - 36 protons bunches



- **Testing color screening**
 - Measuring together χ_c , J/Ψ and Ψ' is crucial
 - Open charm is the appropriate reference
- **LHCb**
 - **Fantastic opportunity** to perform this program
 - SMOG acts as a fixed target (He, Ne, Ar, Xe, Kr)
 - First full data taking in 2015
 - **October 2015 : 3 days of pAr (~20000 J/ψ)**
 - **November-december 2015 : 21 days of PbAr (~15000 J/ψ)**
 - **Leading role of LAL/LLR** team for heavy flavor: Francesco Bossu, Frédéric Fleuret, Giulia Manca, Laure Massacrié, Patrick Robbe, Yanxi Zhanj
- **Not discussed in this talk : a thorough pA program**
 - Capable of reaching **large x_2** (up to 0.3 for J/ψ , 0.9 for Y)
 - First look at the program proposed by the **AFTER** team (after.in2p3.fr)

backup

TABLE I: Centrality bin, number of NN collisions, nuclear overlap function, charm cross section per NN collision, and total charm multiplicity per NN collision, in $\sqrt{s_{NN}} = 200$ GeV Au+Au reactions.

Centrality	N_{coll}	T_{AA} (mb^{-1})	$\frac{1}{T_{AA}} \frac{dN_{c\bar{c}}}{dy} \Big _{y=0}$ (μb)	$N_{c\bar{c}}/T_{AA}$ (μb)
min. bias	258 ± 25	6.14 ± 0.45	$143 \pm 13 \pm 36$	$622 \pm 57 \pm 160$
0–10 %	955 ± 94	22.8 ± 1.6	$137 \pm 21 \pm 35$	$597 \pm 93 \pm 156$
10–20 %	603 ± 59	14.4 ± 1.0	$137 \pm 26 \pm 35$	$596 \pm 115 \pm 158$
20–40 %	297 ± 31	7.07 ± 0.58	$168 \pm 27 \pm 45$	$731 \pm 117 \pm 199$
40–60 %	91 ± 12	2.16 ± 0.26	$193 \pm 47 \pm 52$	$841 \pm 205 \pm 232$
60–92 %	14.5 ± 4.0	0.35 ± 0.10	$116 \pm 87 \pm 43$	$504 \pm 378 \pm 190$

Phys. Rev. Lett. 94, 082301 (2005)

In central Au+Au collisions @ 200 GeV

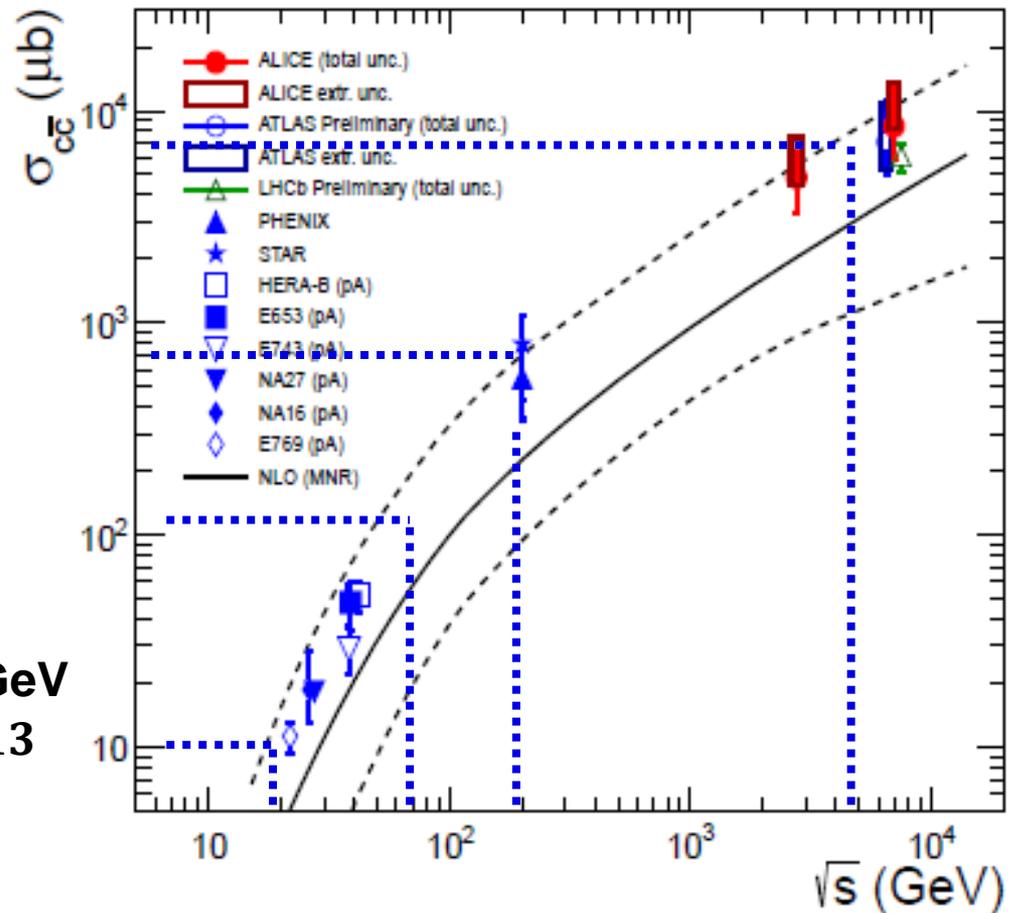
$$N_{c\bar{c}} \sim 597 \cdot 10^{-3} \text{mb} \times 22.8 \text{mb}^{-1} \sim 13$$

~0.1 $c\bar{c}$ @ 20 GeV

~1 $c\bar{c}$ @ 70 GeV

~10 $c\bar{c}$ @ 200 GeV

~100 $c\bar{c}$ @ 5500 GeV



$$\sigma_{c\bar{c}}^{5500 \text{ GeV}} \sim 10 \times \sigma_{c\bar{c}}^{200 \text{ GeV}} \sim 100 \times \sigma_{c\bar{c}}^{70 \text{ GeV}} \sim 1000 \times \sigma_{c\bar{c}}^{20 \text{ GeV}}$$