



# Charmonia in Heavy Ion Collisions should we go **back to SPS ?**

- charmonia in A+A : the current (simplified) picture –
  - back to SPS : the CHIC picture –

## • Motivations

- Quarkonia suppression is a prediction of lattice QCD calculations, for instance :

state	J/ψ(1S)	χ <sub>c</sub> (1P)	ψ'(2S)	Υ(1S)	χ <sub>b</sub> (1P)	Υ(2S)	χ <sub>b</sub> (2P)	Υ(3S)
T <sub>d</sub> /T <sub>c</sub>	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

[H. Satz, J. Phys. G 32 \(2006\)](#)

## • Experimental setups

SPS/CERN – NA38, NA50, NA60 ( $\sqrt{s_{NN}} = 17 - 30$  GeV): fixed target experiments

- ✦ **Statistic** : 100 000's J/ψ
- ✦ **Data sets** : p+A w/ A=p, d, Be, Al, Cu, Ag, W, Pb; S+U, In+In, Pb+Pb
- ✦ **Small rapidity coverage** (typically  $y_{CMS} \in [0,1]$ )

RHIC/BNL – Phenix experiment ( $\sqrt{s_{NN}} = 200$  GeV): collider experiments

- ✦ **Statistic** : 1000's J/ψ (10000's since 2007)
- ✦ **Data sets** : p+p, d+Au, Cu+Cu, Au+Au
- ✦ **Large rapidity coverage** ( $y_{CMS} \in [-0.5,0.5]$ ,  $y_{CMS} \in [-2.2,-1.2]$  and  $y_{CMS} \in [1.2,2.2]$ )

LHC/CERN experiments ( $\sqrt{s_{NN}} = 5,5$  TeV): collider experiments

- ✦ Collider experiments
- ✦ **Statistic** : 100000's J/ψ
- ✦ **Data sets** : p+p, Pb+Pb, p+Pb
- ✦ **Large rapidity coverage** ( $|y_{CMS}| < 2.5$  ATLAS/CMS,  $|y_{CMS}| < 0.9$  and  $-4.0 < y_{CMS} < -2.5$  ALICE)

- Sequential suppression **in a QGP**

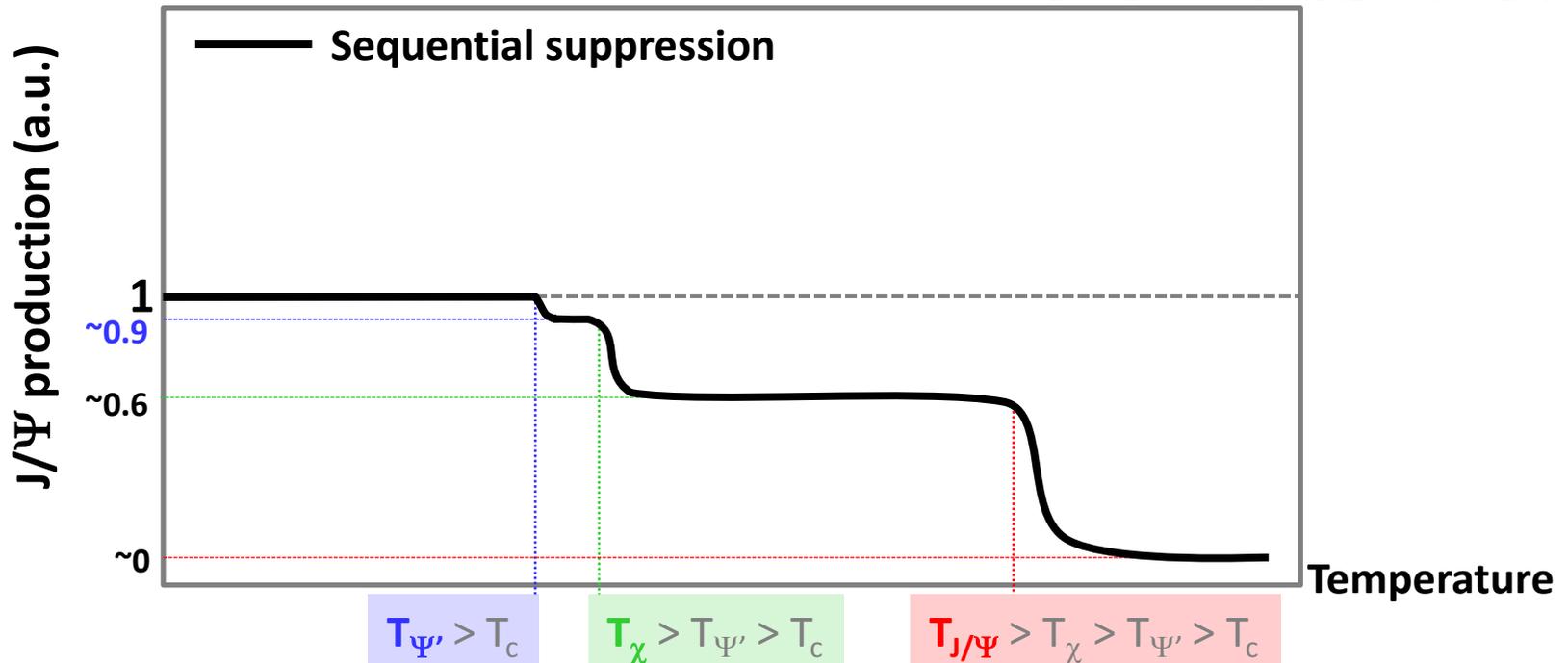
inclusive  $J/\Psi$  yield  $\sim$  **60%** direct  $J/\Psi$  + **30%**  $\chi_c \rightarrow J/\Psi + \gamma$  + **10%**  $\Psi' \rightarrow J/\Psi + X$

Charmonium temperatures  
of dissociation

state	$J/\psi(1S)$	$\chi_c(1P)$	$\psi'(2S)$
$T_d/T_c$	2.10	1.16	1.12

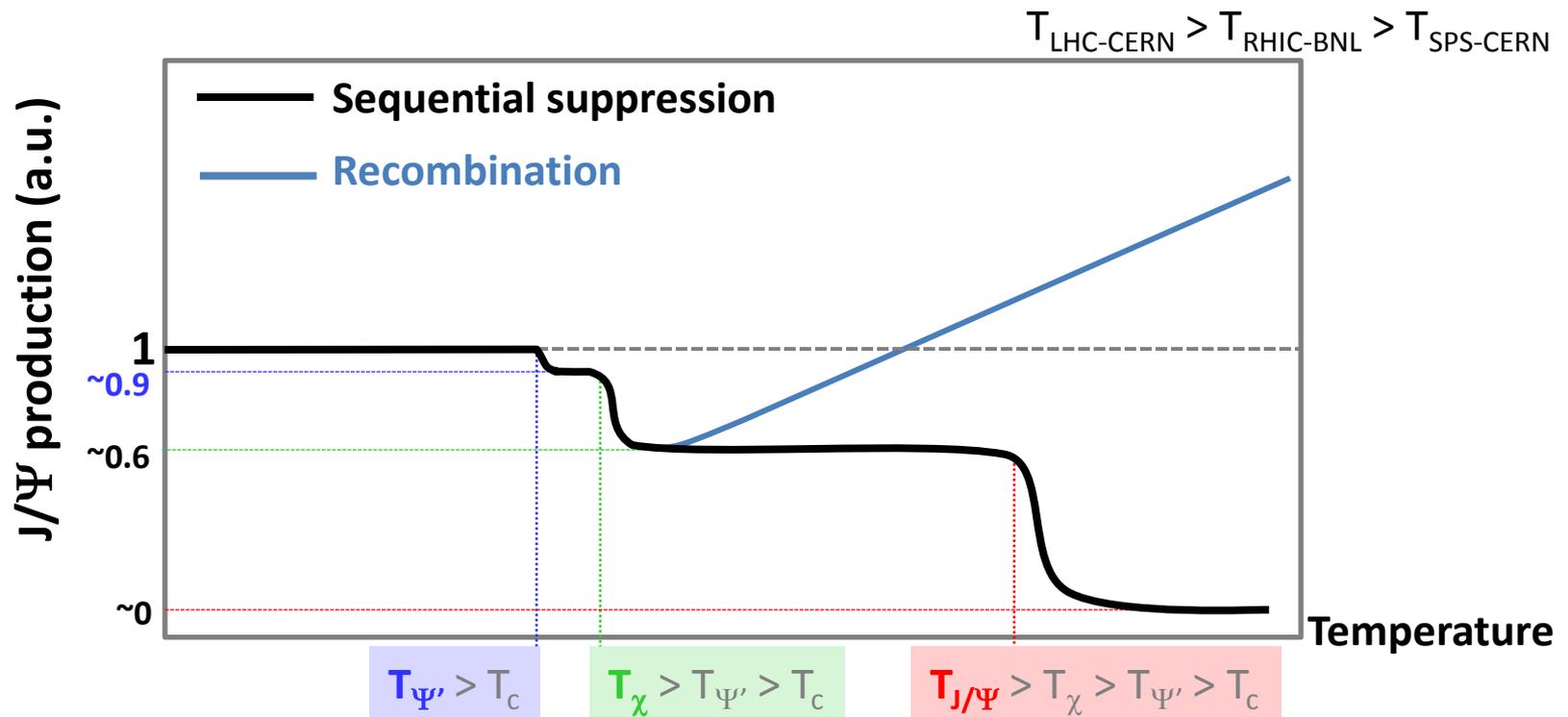
[H. Satz, J. Phys. G 32 \(2006\)](#)

$T_{\text{LHC-CERN}} > T_{\text{RHIC-BNL}} > T_{\text{SPS-CERN}}$



- Recombination **in a QGP**

If QGP at work → **c and  $\bar{c}$  quarks can combine to form a  $J/\Psi$**   
 (require a large number of  $c\bar{c}$  pairs → RHIC ? LHC ?)



- Suppression by comovers (Alternative scenario)**

- Suppression by comovers: [\(Eur.Phys.J.C58:437-444,2008\)](#)

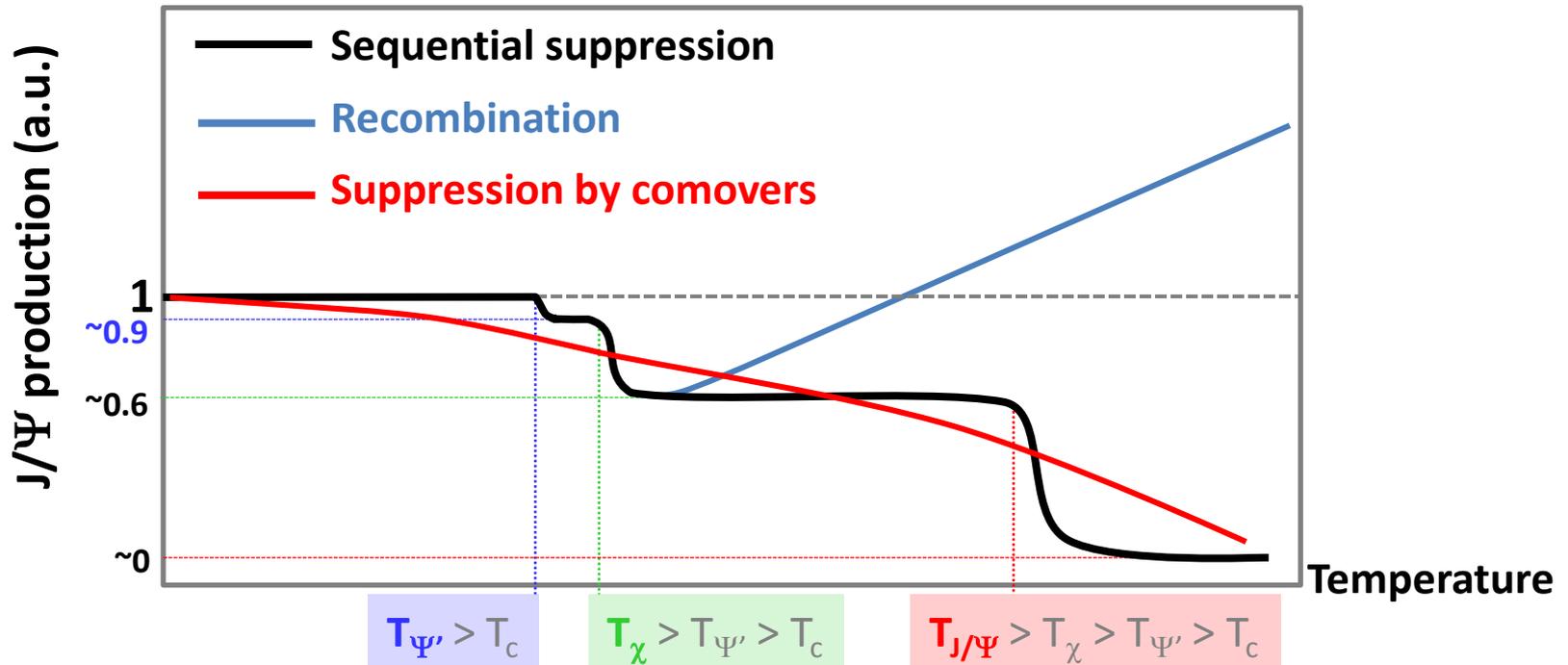
- quarkonia can be broken by interaction with comoving hadrons



Two parameters

Interaction cross section  $\sigma_{co}$       Hadron density  $N^{co}$

$$\tau \frac{dN_{J/\psi}}{d\tau}(b, s, y) = -\sigma_{co} N^{co}(b, s, y) N_{J/\psi}(b, s, y)$$



- SPS (17 GeV): NA38, NA51, NA50, NA60

## Two major results :

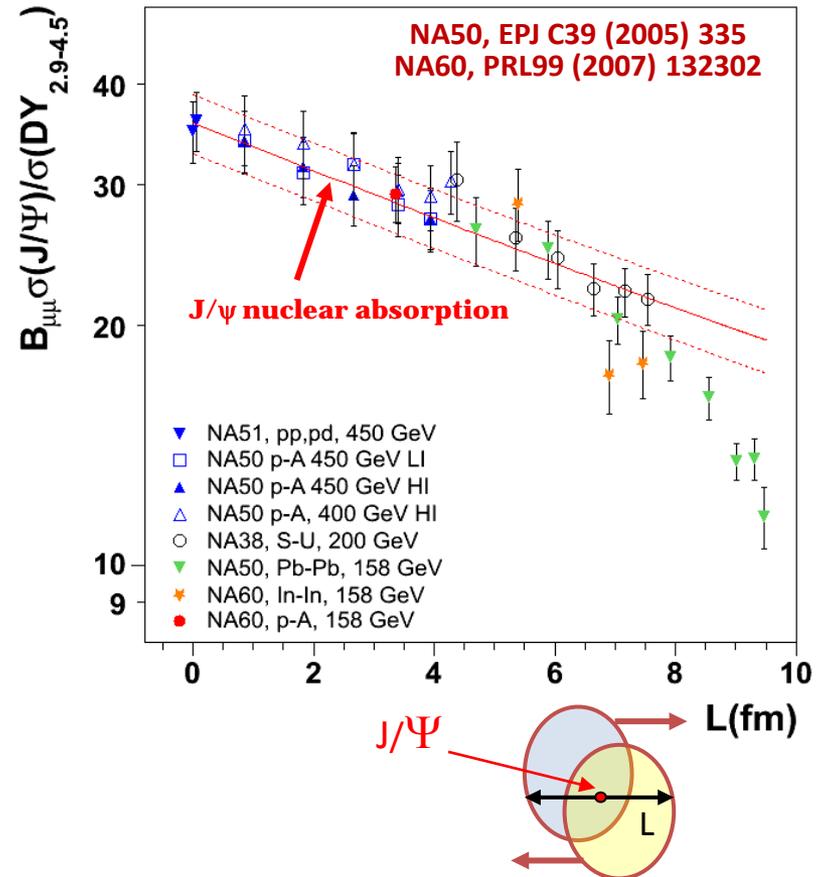
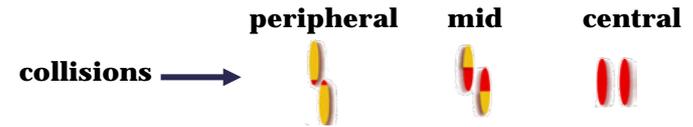
- Observation of **Cold Nuclear Matter effects** : Absorption by nuclear matter

- Suppression observed from p+p to peripheral Pb+Pb
- J/ψ survival probability :

$$S(J/\Psi) \propto e^{-\rho\sigma_{abs}L}$$

- Fit to data:  $\sigma_{abs} = 4.18 \pm 0.35 \text{ mb}$

- Observation of **Anomalous suppression** in Pb+Pb (NA50) central collisions when compared with Cold Nuclear Matter effects.



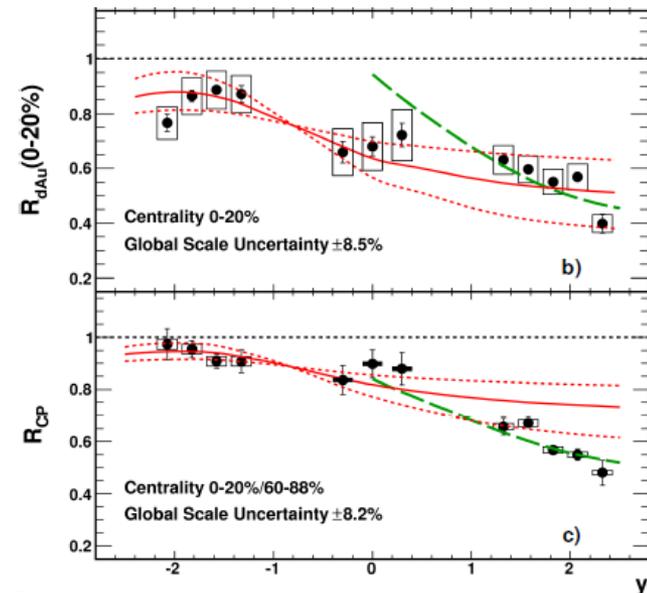
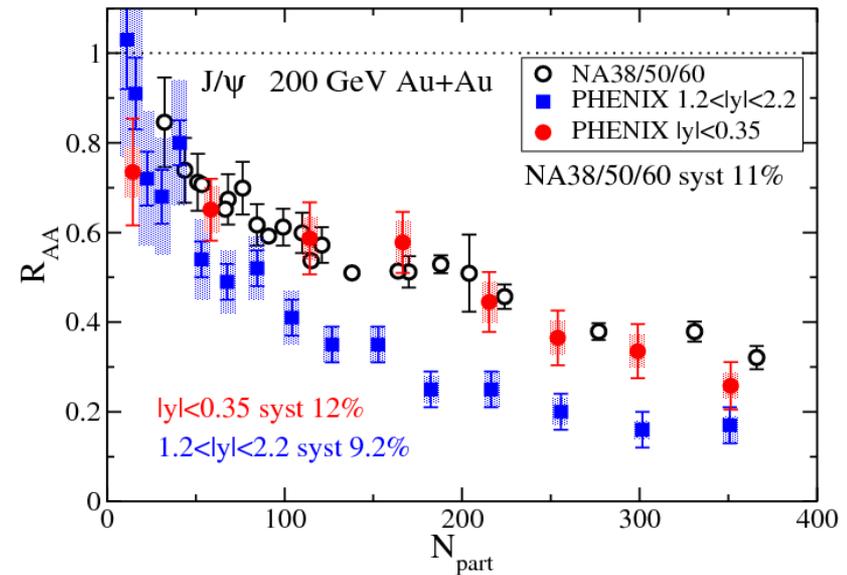
- **RHIC (200 GeV) .vs. SPS (17 GeV)**

1. **Hot and dense matter effects**

- Measure  $J/\Psi$  in Au+Au (RHIC) Pb+Pb (SPS)
- Compare at same rapidity (same  $y \sim$  same  $x_F$ )
  - $0 < y < 1$  at SPS (**NA50/NA60**)
  - $|y| < 0.35$  at RHIC (**PHENIX**)
- Expected larger suppression at RHIC due to larger energy density
- observe **SIMILAR SUPPRESSION at mid rapidity**
- Observe **LARGER SUPPRESSION at forward rapidity**

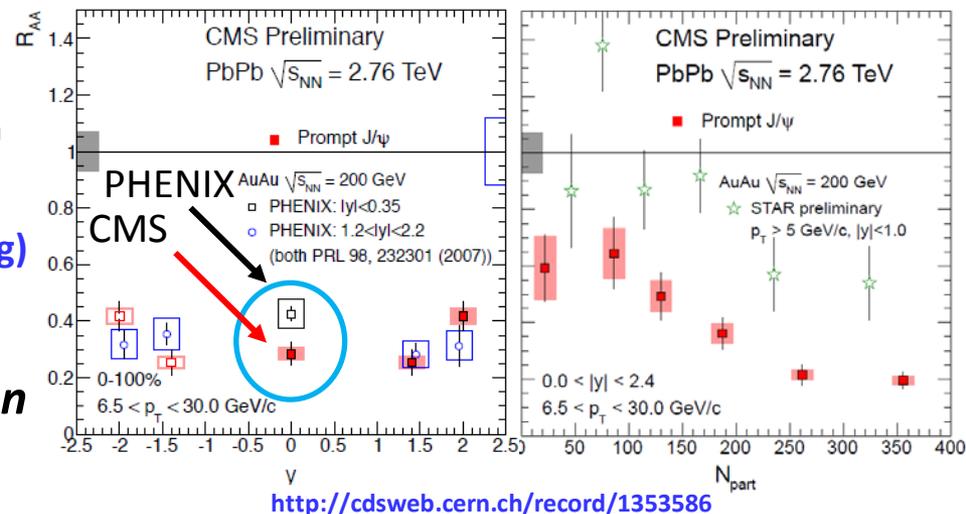
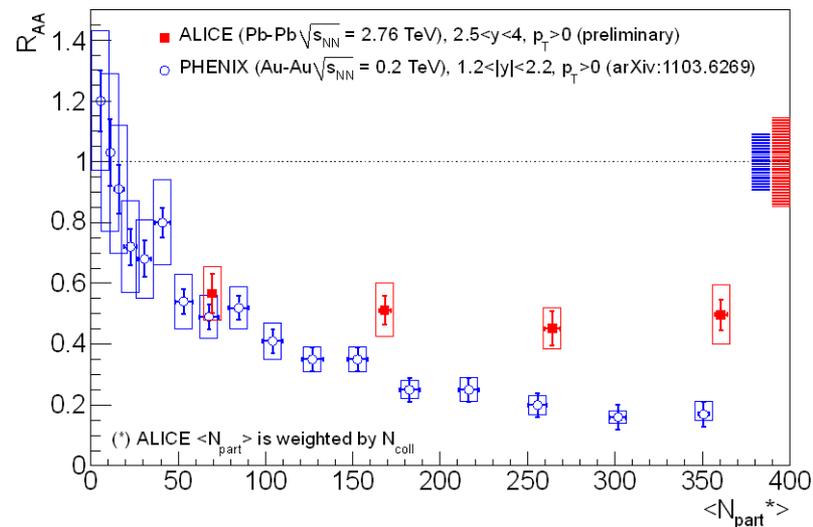
2. **Cold Nuclear Matter effects at RHIC**

- Measure  $J/\Psi$  production in d+Au collisions
- Observe **LARGER SUPPRESSION at forward rapidity (small  $x_2$ )**
- Pattern still not fully understood
- Difference forward.vs.mid rapidity may explain larger suppression observed in forward Au+Au



- **RHIC (200 GeV) .vs. LHC (2.76 TeV) at forward rapidity**
  - Compare PHENIX vs ALICE
    - $1.2 < |y| < 2.2$  at RHIC/PHENIX
    - $2.5 < y < 4$  at LHC/ALICE
  - **LESS SUPPRESSION** at LHC .vs. RHIC
  - Could be due to **recombination** effects
  
- **RHIC (200 GeV) .vs. LHC (2.76 TeV) at mid-rapidity**
  - Compare PHENIX, STAR vs CMS
    - $|y| < 0.35$  at RHIC/PHENIX
    - $|y| < 1$  at RHIC/STAR
    - $|y| < 1$  at LHC/CMS
  - **MORE SUPPRESSION** at LHC .vs. RHIC
    - $p_T > 6.5$  GeV/c  $\rightarrow$  in principle no recombination applies
    - larger suppression due to **QGP effects ?**
  - Hint of **sequential suppression ? (J/ $\Psi$  melting)**

**Caution : Need CNM effects comparison**



- Overall possible  $J/\Psi$  (simplified) picture

- Similar suppression** at SPS vs. RHIC

$\Psi'$  and  $\chi_c$  suppression only ?

- CMS: Larger suppression** at LHC

$p_T > 6.5 \text{ GeV}/c \rightarrow$  « outside » recombination regime ?

Hint of sequential suppression ?

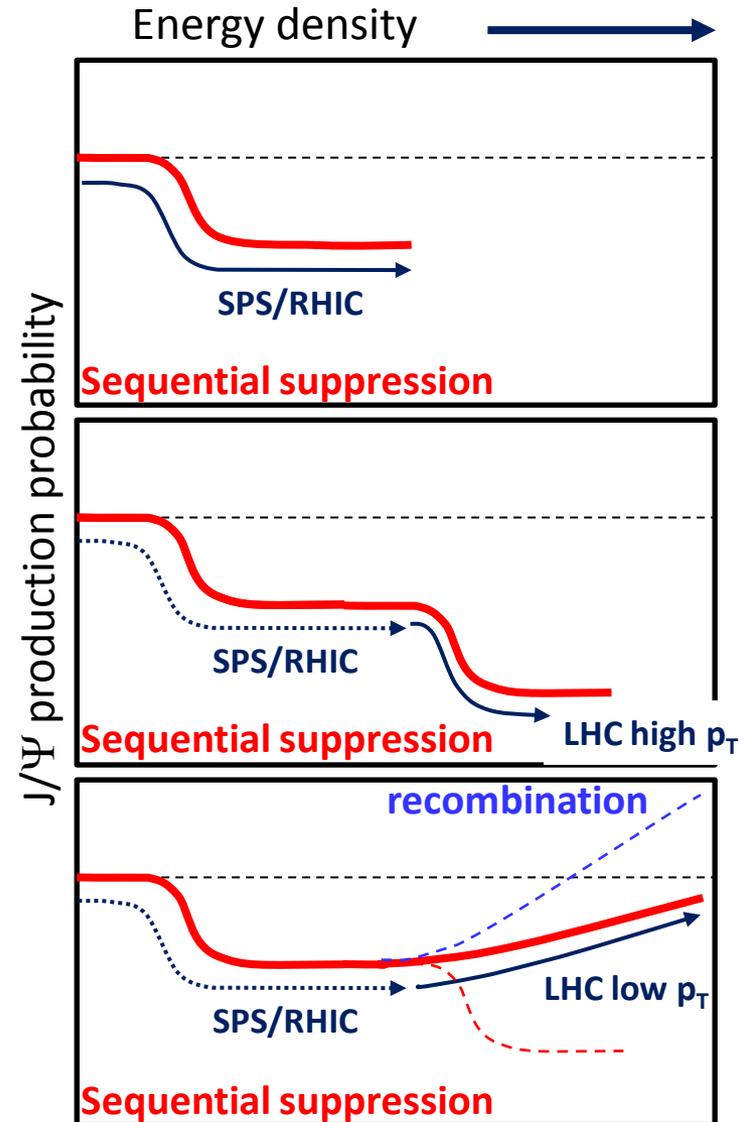
(assuming CNM effects are the same or smaller)

- ALICE  $|y| > 2.5$ : Smaller suppression** at LHC

« inside » recombination regime ?

Hint of recombination ?

(assuming CNM effects are the same or larger)



- **Answers to these questions are mandatory :**
  - **What are CNM effects at LHC ?** → p+Pb run
    - *Shadowing should be large at forward rapidity*
    - *Shadowing should be small at high  $p_T$*
    - *Resonance break-up cross section should be small*
  
  - **Is recombination mechanism at work ?** → ALICE.vs.CMS at  $|y|=0$ 
    - *If smaller suppression observed at mid-rapidity and low  $p_T$*
  
  - **Is sequential suppression at work ?**
    - *Need several (at least two) resonances*
    - *$\Psi'$  is not a good probe because of comovers*
    - *Should measure  $\chi_c$*  → unreachable

- **Measuring  $\chi_c$  in A+A:**
  - test charmonia sequential suppression
  - How  $\chi_c$  is suppressed relative to J/ $\Psi$  ? Dependence with  $y$ ,  $p_T$ , centrality?
  - ➔ **Mandatory to draw the whole picture (SPS .vs. RHIC .vs. LHC)**
  
- **Should measure  $\chi_c$  at SPS. Why at SPS ?**
  - *If we understand SPS, we understand RHIC (same suppression)*
  - Anomalous suppression has been seen at SPS
  - Appropriate range of energy density: can investigate  $\Psi'$ ,  $\chi_c$  and J/ $\Psi$  suppression
  - **On average, 0.1  $c\bar{c}$  pair/event**
    - ➔ **No recombination at SPS**
  
- **Fixed target experiment ?**
  - Can operate many target species
  - ➔ **Better control of CNM effects**

## Charmonia suppression

### At SPS

60% direct  $J/\Psi$   
 + 30%  $\chi_c \rightarrow J/\Psi + \gamma$   
 + 10%  $\Psi' \rightarrow J/\Psi + X$   
**Inclusive  $J/\Psi$  yield**

Two possible scenarios:

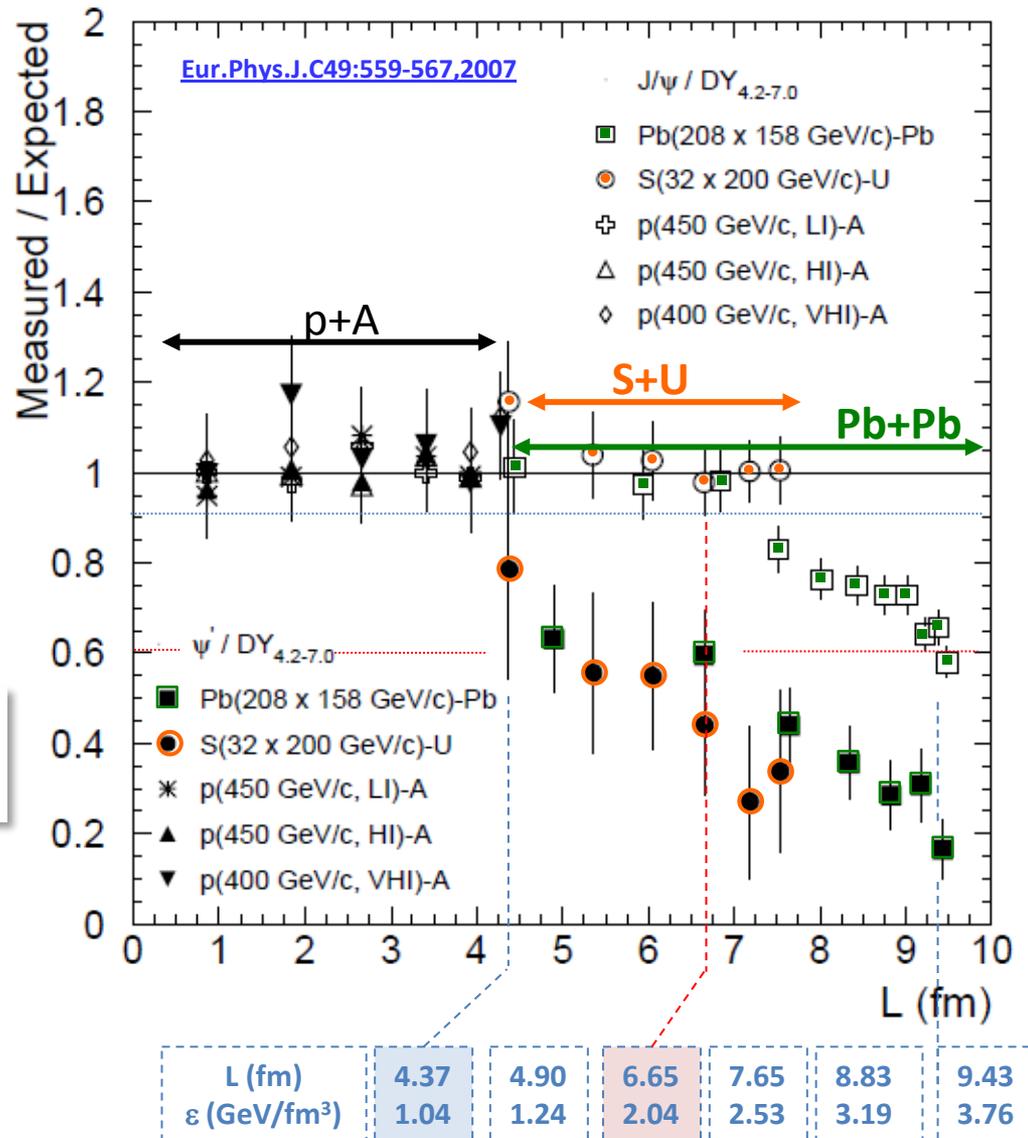
- sequential suppression (QGP)
- comovers (no QGP)

### Temperature of dissociation

state	$J/\psi(1S)$	$\chi_c(1P)$	$\psi'(2S)$	$\Upsilon(1S)$	$\chi_b(1P)$	$\Upsilon(2S)$	$\chi_b(2P)$	$\Upsilon(3S)$
$T_d/T_c$	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

### Binding energy

state	$\eta_c$	$J/\psi$	$\chi_{c0}$	$\chi_{c1}$	$\chi_{c2}$	$\psi'$
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
$\Delta E$ [GeV]	0.75	0.64	0.32	0.22	0.18	0.05



## Two possible scenarios

### 1. QGP (sequential suppression)

state	$\eta_c$	$J/\psi$	$\chi_{c0}$	$\chi_{c1}$	$\chi_{c2}$	$\psi'$
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
$\Delta E$ [GeV]	0.75	0.64	0.32	0.22	0.18	0.05

Because  $\Delta E(\Psi') \sim 50$  MeV

- $\Psi'$  easily suppressed by comovers

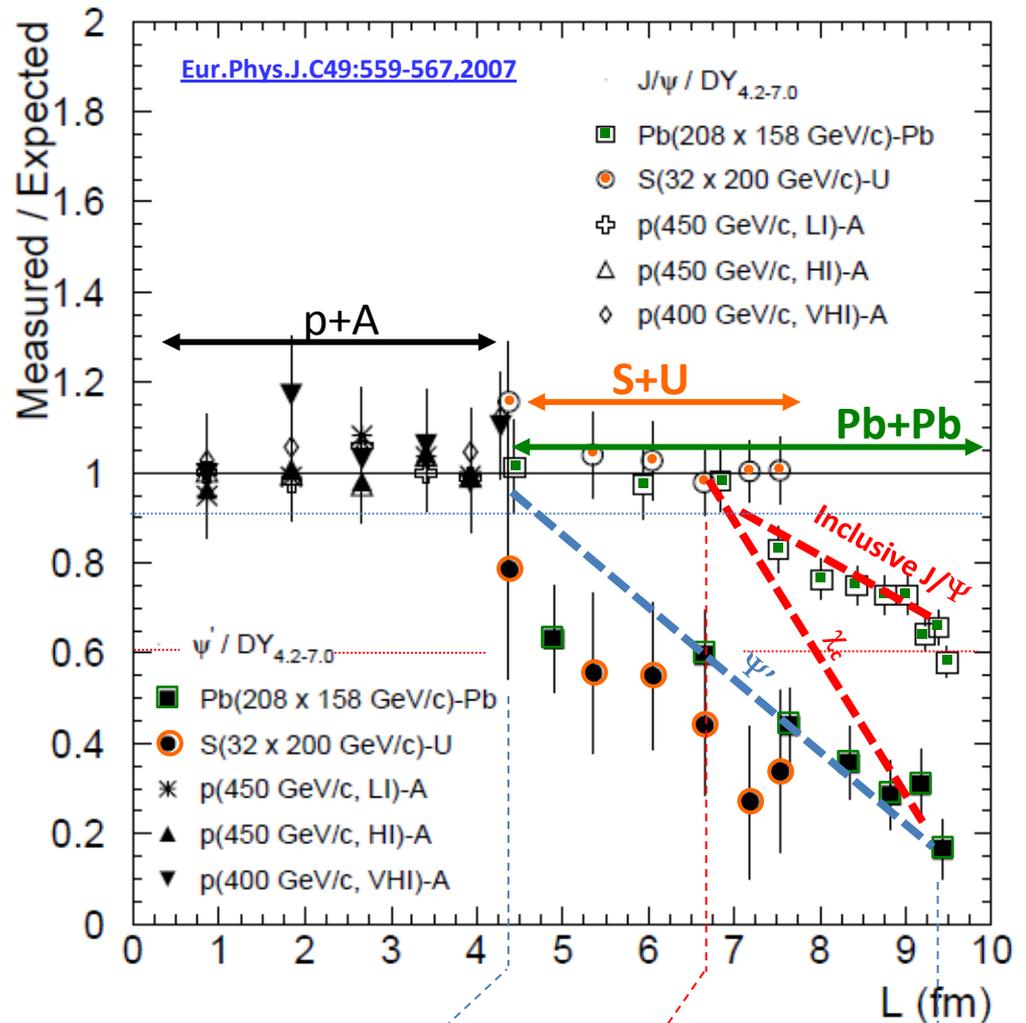
Because  $\Delta E(\chi_c) \sim 200$  MeV and  $\Delta E(J/\Psi) \sim 600$  MeV

- $\chi_c$  and  $J/\Psi$  hardly suppressed by comovers

If  $\chi_c$  suppressed by QGP,

- $\chi_c$  slope strongly steeper than  $J/\Psi$  and  $\Psi'$

Measuring  $\chi_c$  suppression pattern will (in)validate this



L (fm)	4.37	4.90	6.65	7.65	8.83	9.43
$\epsilon$ (GeV/fm <sup>3</sup> )	1.04	1.24	2.04	2.53	3.19	3.76

Note that direct  $J/\Psi$  can be experimentally estimated  
 $\text{Yield}_{\text{incl. } J/\Psi} - \text{Yield}_{\chi_c \rightarrow J/\Psi + \gamma} - \text{Yield}_{\Psi'} \sim \text{Yield}_{\text{direct } J/\Psi}$

## Two possible scenarios

### 2. No QGP (full comovers)

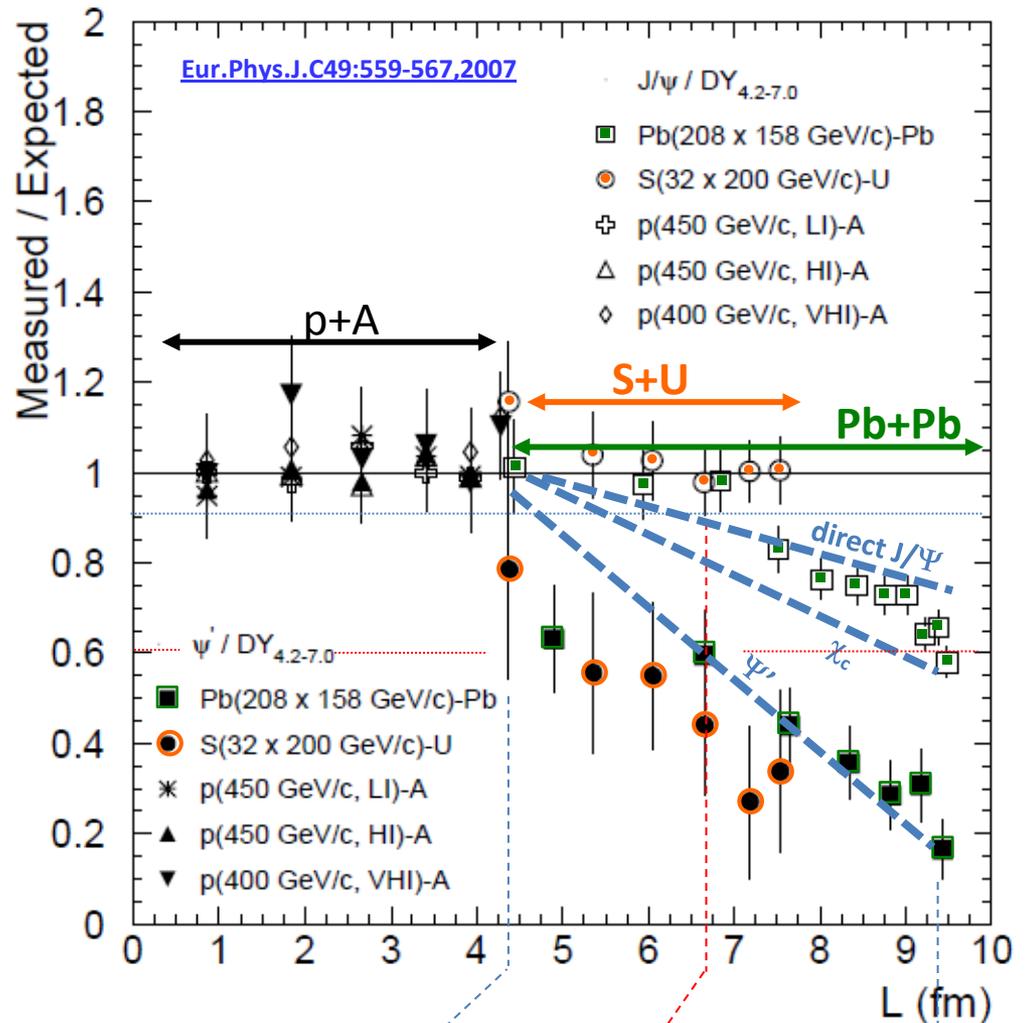
state	$\eta_c$	$J/\psi$	$\chi_{c0}$	$\chi_{c1}$	$\chi_{c2}$	$\psi'$
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
$\Delta E$ [GeV]	0.75	0.64	0.32	0.22	0.18	0.05

Because  $\sigma_{J/\Psi-co} \leq \sigma_{\chi_c-co} \leq \sigma_{\Psi'-co}$

- $\Psi'$  slope slightly steeper than  $\chi_c$
- $\chi_c$  slope slightly steeper than  $J/\Psi$

### Measuring

$\chi_c$  suppression pattern  
will (in)validate this



L (fm)	4.37	4.90	6.65	7.65	8.83	9.43
$\epsilon$ (GeV/fm <sup>3</sup> )	1.04	1.24	2.04	2.53	3.19	3.76

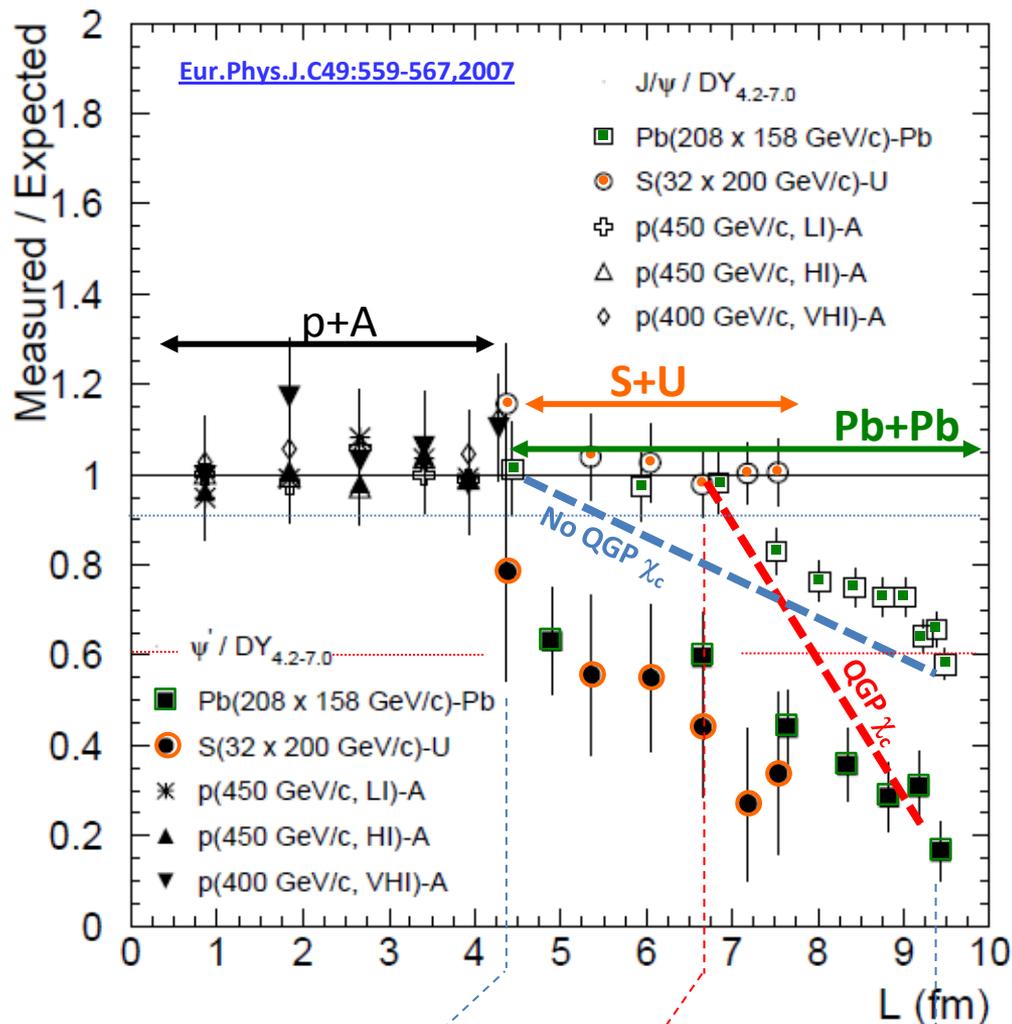
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- Conclusion :**

measuring  $\Psi'$ ,  $J/\Psi$  and  $\chi_c$  suppression pattern

will answer the question

- QGP
- no QGP



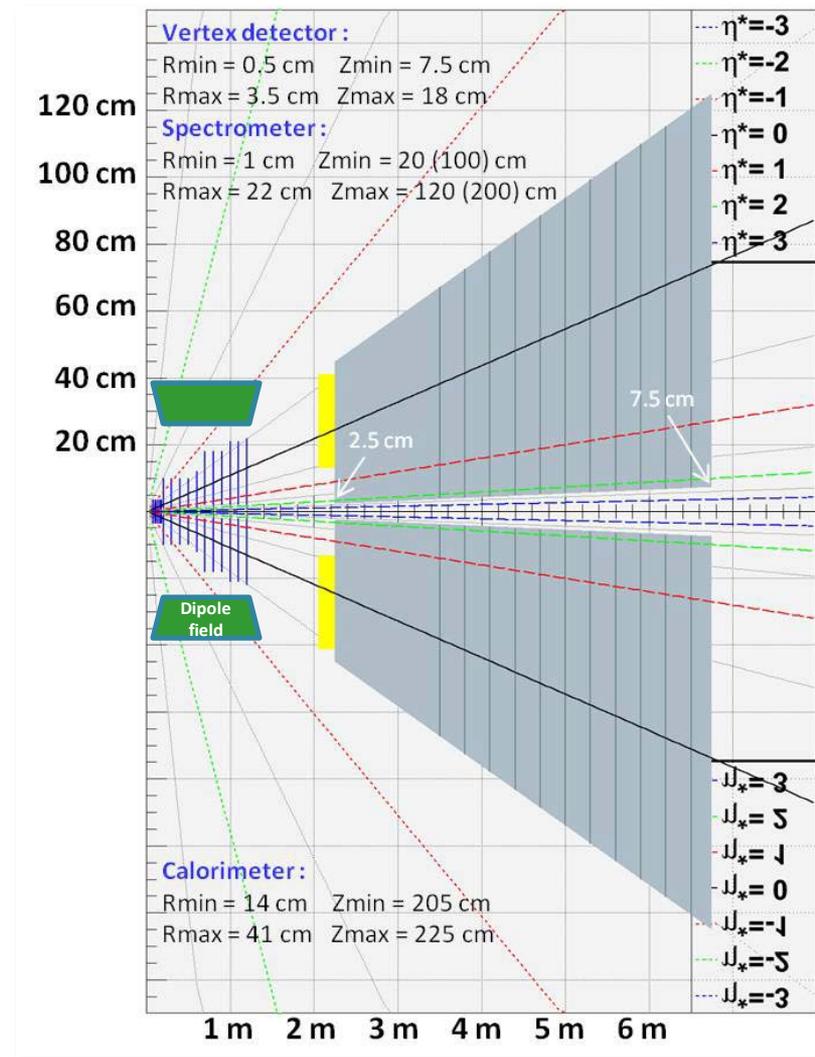
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Note that direct  $J/\Psi$  can be experimentally estimated  
 $Yield_{incl. J/\Psi} - Yield_{\chi_c \rightarrow J/\Psi + \gamma} - Yield_{\Psi'} \sim Yield_{direct J/\Psi}$

- **Primary goals :**
  - $\chi_c \rightarrow J/\Psi + \gamma \rightarrow \mu^+ \mu^- \gamma$  at  $y_{CMS} = 0$
  - $J/\Psi \rightarrow \mu^+ \mu^-$  in large  $y_{CMS}$  range
- **Detector features : very compact**
  - 1. Spectrometer**
    - Measure tracks before absorber  $\rightarrow \sigma_M \sim 20 \text{ MeV}/c^2$
    - Covers  $y_{CMS} [-0.5, 2]$   $\rightarrow$  need high segmentation
    - $\rightarrow$  Silicon technologies
  - 2. Calorimeter**
    - Measuring  $\gamma$  in high  $\pi^0$  multiplicity environment
    - $\rightarrow$  ultra-granular ECal (Calice)
  - 3. Absorber/trigger**
    - Using 4.5 m thick Fe to absorb  $\pi/K$  and low P  $\mu^{+/-}$
    - Can use smaller absorber if Fe magnetized
    - Trigger to be defined (expected rate = 0.3 kHz)
- **Expected performances**
  - 1. tracking :**

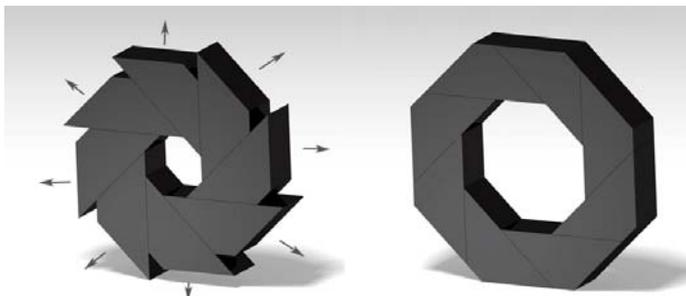
$$\frac{\Delta P}{P} \sim 1\% \text{ within } 1\text{m long } 2.5\text{T } \vec{B}$$
  - 2. Calorimetry :**

$$\frac{\Delta E}{E} \sim \frac{20\%}{\sqrt{E}}$$



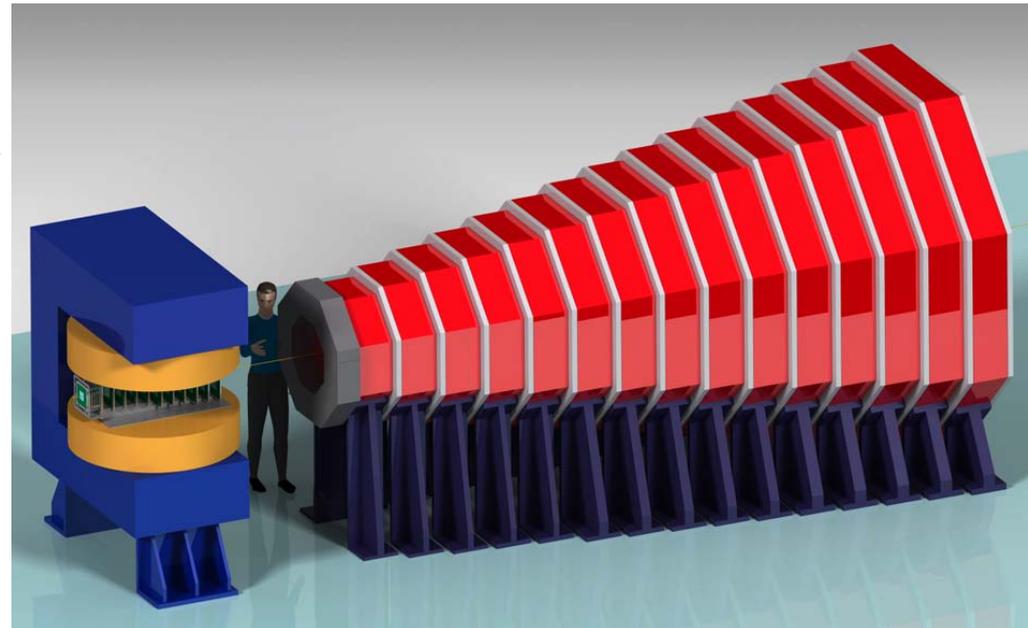
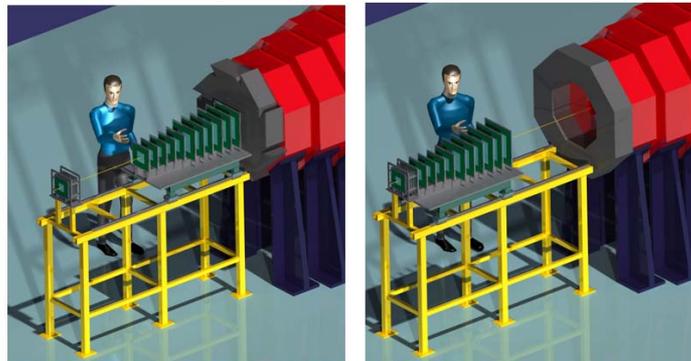
- CHIC: Experimental setup flexibility

**Very compact detector**  
(full detector simulation ongoing) →



**Forward rapidity**

**Mid rapidity**

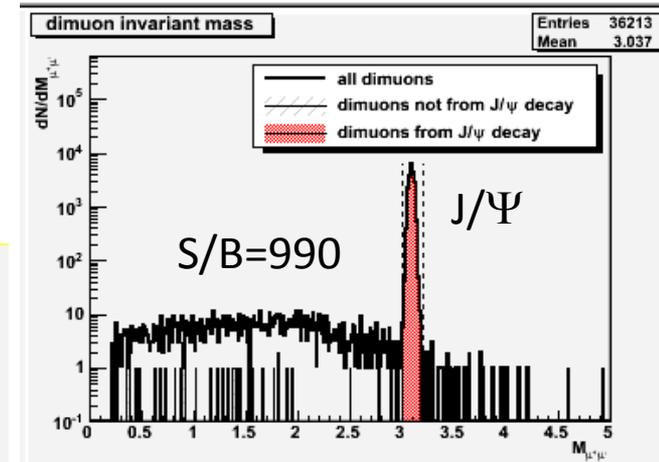
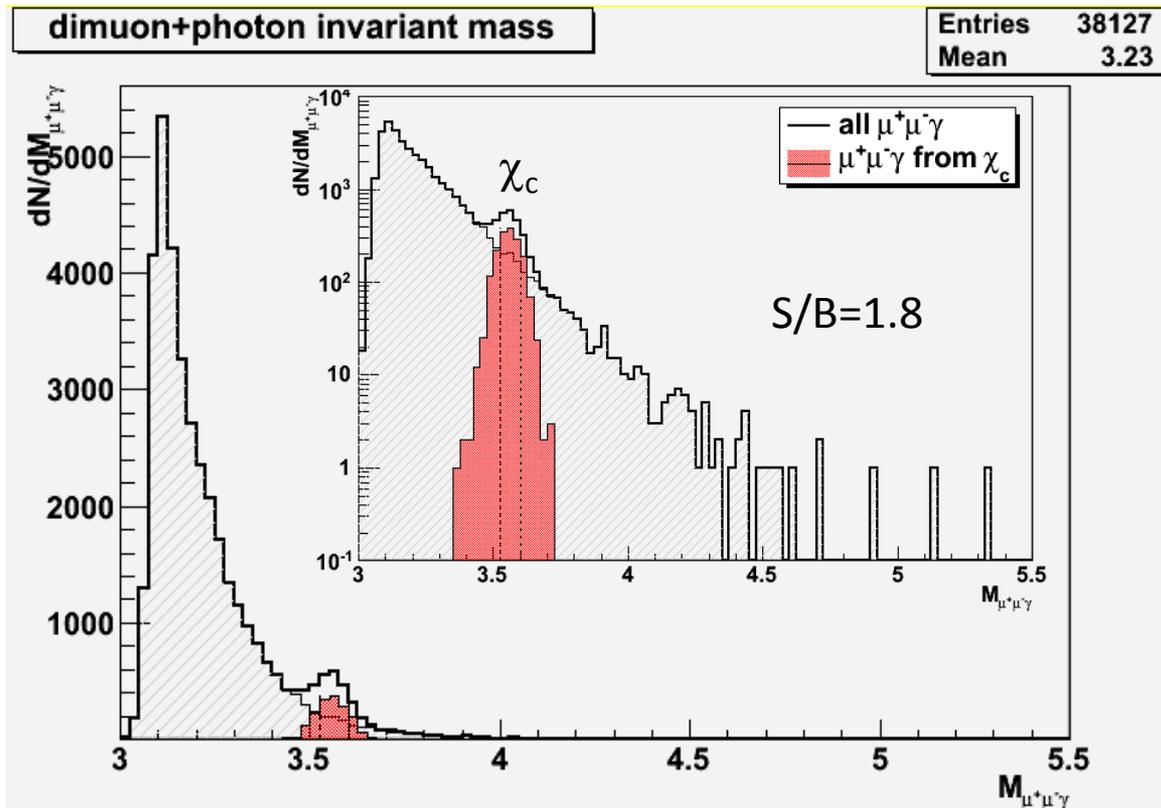


**Large rapidity coverage**

- fixed target mode → high flexibility
- displace tracker to access large rapidity
- modify calorimeter to access large rapidity

- **Typical mass plots**

- 200 000 Pb+Pb minBias EPOS events
  - 140 000 events with  $J/\Psi$  embedded (70%)
  - 60 000 events with  $\chi_c$  embedded (30%)



After acceptance and selection cuts:

- 35 000  $J/\Psi$   
 → **acc x eff = 17.4%**
- 1700  $\chi_c$   
 → **acc x eff = 2.8 %**

- **Typical one month Pb+Pb run with a 4mm thick target**

- ~ 200 000 inclusive  $J/\Psi \rightarrow \mu^+\mu^-$  expected

- 2 extreme scenarios:

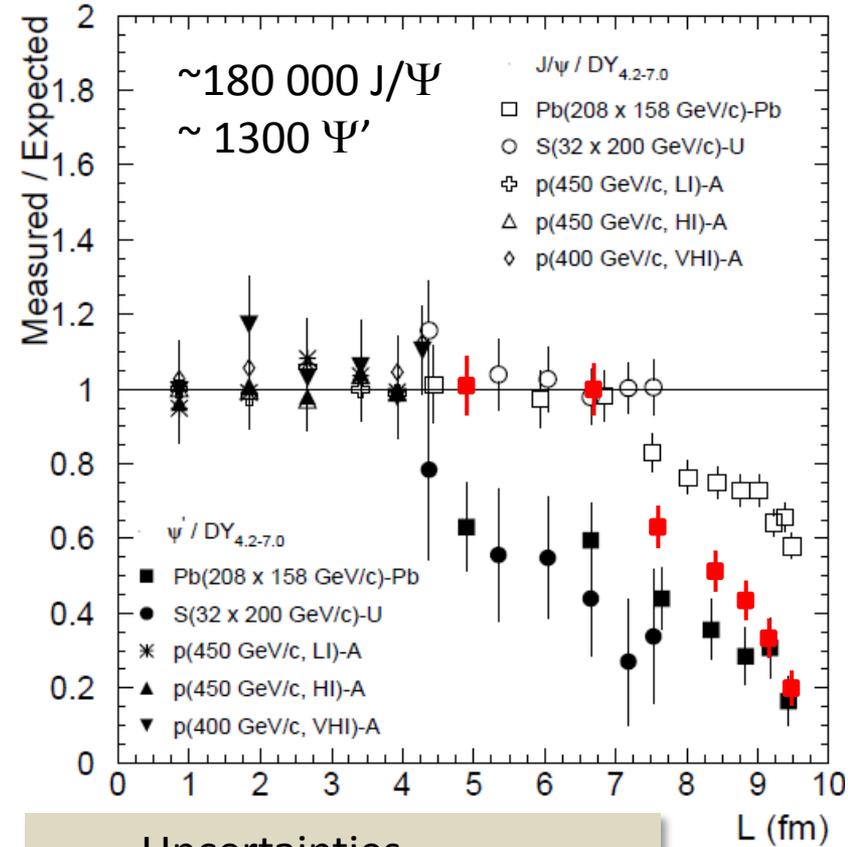
- If  $\chi_c$  suppressed as  $J/\Psi$   $\frac{\chi_c \text{ yield}}{J/\Psi \text{ yield}} \sim 4\%$

$$\Rightarrow \left( \begin{array}{c} \text{most periph.} \\ \chi_c \text{ yield} \end{array} \right) = 16942 \times 4\% = 677$$

- If  $\chi_c$  suppressed as  $\Psi'$   $\frac{\chi_c \text{ yield}}{\Psi' \text{ yield}} = 2.18$

$$\left( \begin{array}{c} \text{most periph.} \\ \chi_c \text{ yield} \end{array} \right) = 16942 \times 4\% \times 0.6 = 406$$

$E_T$ range (GeV)	$\psi'$	$J/\psi$	$\chi_c$ as $J/\Psi$	$\chi_c$ as $\Psi'$
3–20	$186 \pm 25$	$16942 \pm 146$	677	406
20–35	$243 \pm 31$	$25229 \pm 181$	1010	530
35–50	$227 \pm 35$	$27276 \pm 192$	1091	495
50–65	$193 \pm 36$	$27681 \pm 196$	1107	421
65–80	$154 \pm 36$	$27315 \pm 200$	1093	336
80–95	$159 \pm 37$	$25111 \pm 193$	1004	347
95–150	$110 \pm 40$	$28570 \pm 209$	1143	240
			<b>7125</b>	<b>2775</b>



Uncertainties  
 $\chi_c$  stat > 2 x  $\Psi'$  stat  
 $\Rightarrow \chi_c$  error <  $\Psi'$  error /  $\sqrt{2}$

- **Conclusion**

- Core benchmark : **unique test of  $\chi_c$  in heavy ion collisions**
- What we didn't discuss :
  - **CHIC p+A program**
    - **9 months of proton beam available** – to be compared to the usual one month
    - **capability to access  $x_f = 1$**
    - physics of saturation : shadowing, CGC, energy loss (Arléo, Peigné)
    - charmonium hadronisation time
    - charmonium absorption cross section
  - **Drell-Yan studies**
  - **Open charm studies**
  - **Charged/neutral hadrons studies**
  - **Photons studies**
  - **Low mass dileptons**

# Backup slides

- Sequential suppression **in a QGP**

If QGP at work → **threshold effect**

Above threshold  $\frac{\epsilon}{T^4} = cte$

$$\Rightarrow \frac{\epsilon_d^{\Psi'}}{(T_d^{\Psi'})^4} = \frac{\epsilon_d^{\chi_c}}{(T_d^{\chi_c})^4}$$

Temperatures of dissociation :

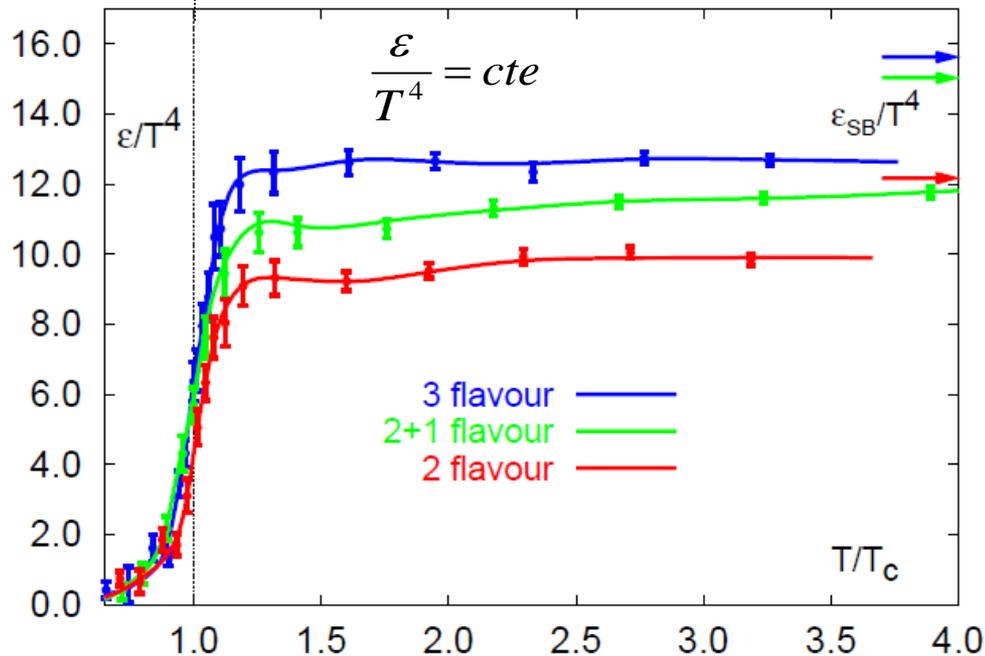
H. Satz, J. Phys. G 32 (2005)

state	J/ψ(1S)	χ <sub>c</sub> (1P)	ψ'(2S)	Υ(1S)	χ <sub>b</sub> (1P)	Υ(2S)	χ <sub>b</sub> (2P)	Υ(3S)
T <sub>d</sub> /T <sub>c</sub>	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

$$\frac{\epsilon_c^{\chi_c}}{(1.16 \times T_c)^4} = \frac{\epsilon_c^{\Psi'}}{(1.12 \times T_c)^4} \Rightarrow \epsilon_c^{\chi_c} = \left(\frac{1.16}{1.12}\right)^4 \epsilon_c^{\Psi'} \Rightarrow \epsilon_c^{\chi_c} = 1.15 \times \epsilon_c^{\Psi'}$$

No QGP ← QGP  
threshold

[F. Karsch, Lect. Notes Phys. 583 \(2002\) 209](#)



## Sequential suppression in a QGP

Theoretically, expect

$$\varepsilon_c^{\chi_c} = \left(\frac{1.16}{1.12}\right)^4 \varepsilon_c^{\Psi'} = 1.15 \times \varepsilon_c^{\Psi'}$$

Experimentally,  $\Psi'$  suppression starts at

L (fm)	4.37
$\varepsilon$ (GeV/fm <sup>3</sup> )	1.04

Theoretically,  $\chi_c$  suppression should start at

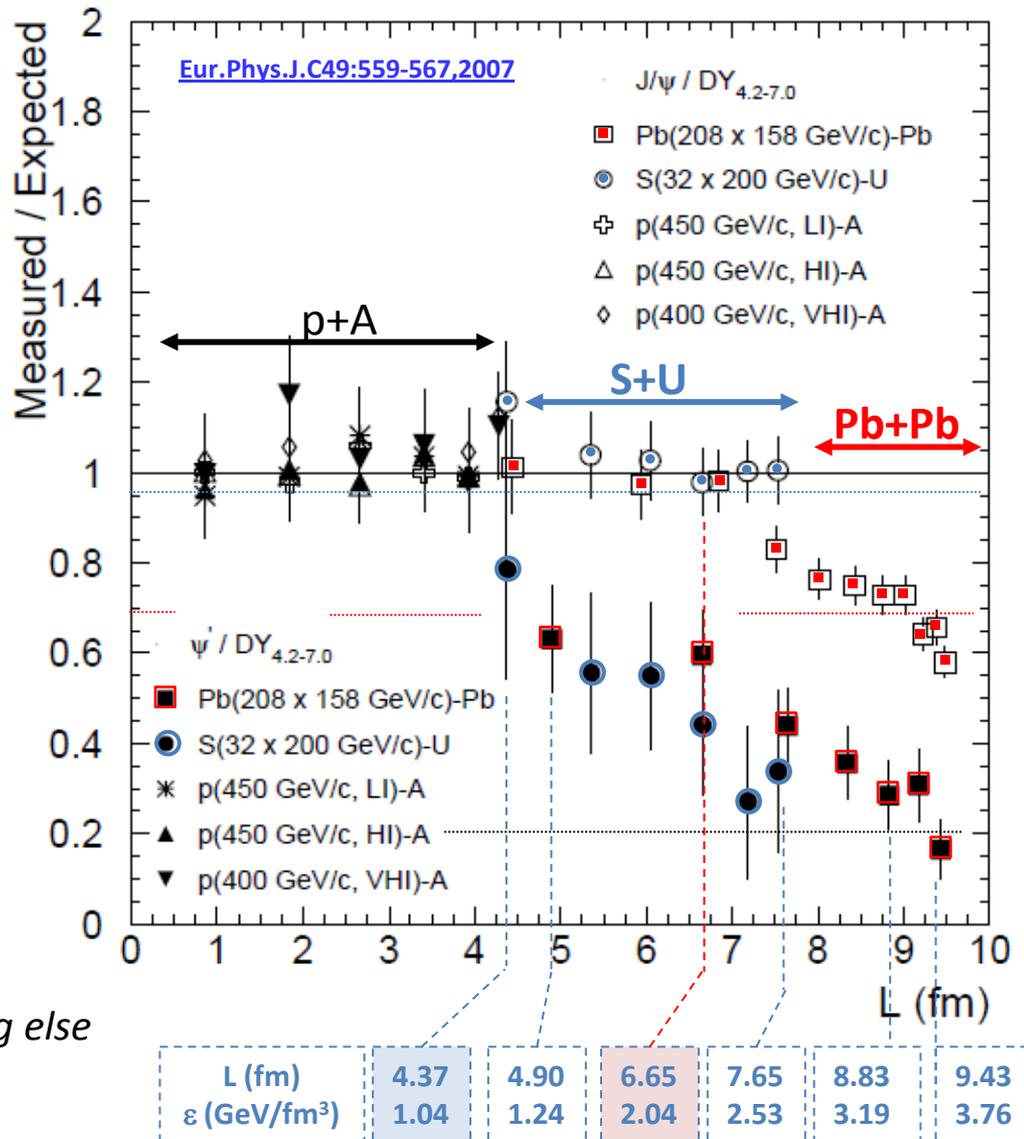
L (fm)	4.9
$\varepsilon$ (GeV/fm <sup>3</sup> )	1.2

Experimentally,  $J/\Psi$  suppression starts at

L (fm)	6.65
$\varepsilon$ (GeV/fm <sup>3</sup> )	2.04

### Conclusion

either theoretical predictions are wrong,  
or  $\Psi'$  is previously suppressed by something else



- Sequential suppression **by comovers**

[\(Eur.Phys.J.C58:437-444,2008\)](#)

- Suppression by comovers:

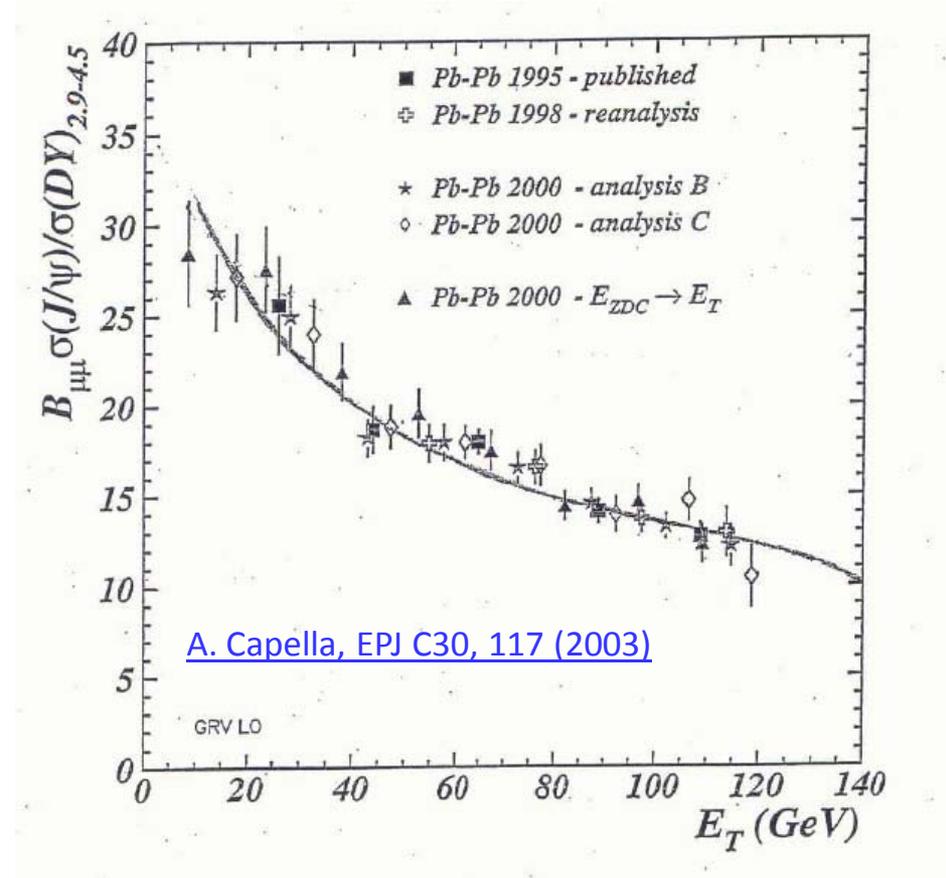
- quarkonia can be broken by interaction with comoving partons/hadrons



$$\tau \frac{dN_{J/\psi}}{d\tau}(b, s, y) = -\sigma_{co} N^{co}(b, s, y) N_{J/\psi}(b, s, y)$$

- Two parameters

- Hadron density  $N^{co}$
- Interaction cross section  $\sigma_{co}$



- **Sequential suppression by comovers**



- Suppression by comovers:

- quarkonia can be broken by interaction with comoving partons/hadrons

- Two parameters  $\tau \frac{dN_{J/\psi}}{d\tau}(b, s, y) = -\sigma_{co} N^{co}(b, s, y) N_{J/\psi}(b, s, y)$

- Hadron density
- Interaction cross section  $\sigma_{co}$

- **There is a hierarchy in the suppression**

- $\sigma_{co}$  is linked to the quarkonium binding energy
- The larger the binding energy, the smaller the  $\sigma_{co}$
- But  $\sigma_{co}$  is theoretically unknown (must be fitted on the data)

- Sequential suppression

- $\Delta E(J/\Psi) > \Delta E(\chi_c) > \Delta E(\Psi')$

- $\sigma_{J/\Psi-co} \leq \sigma_{\chi_c-co} \leq \sigma_{\Psi'-co}$

state	$\eta_c$	$J/\psi$	$\chi_{c0}$	$\chi_{c1}$	$\chi_{c2}$	$\psi'$
mass [GeV]	2.98	3.10	3.42	3.51	3.56	3.69
$\Delta E$ [GeV]	0.75	0.64	0.32	0.22	0.18	0.05

Quarkonium binding energy  
 $(\Delta E = M_{\text{quarkonium}} - 2M_D)$

- Sequential suppression  
**by comovers**

[Eur.Phys.J.C58:437-444,2008](#)

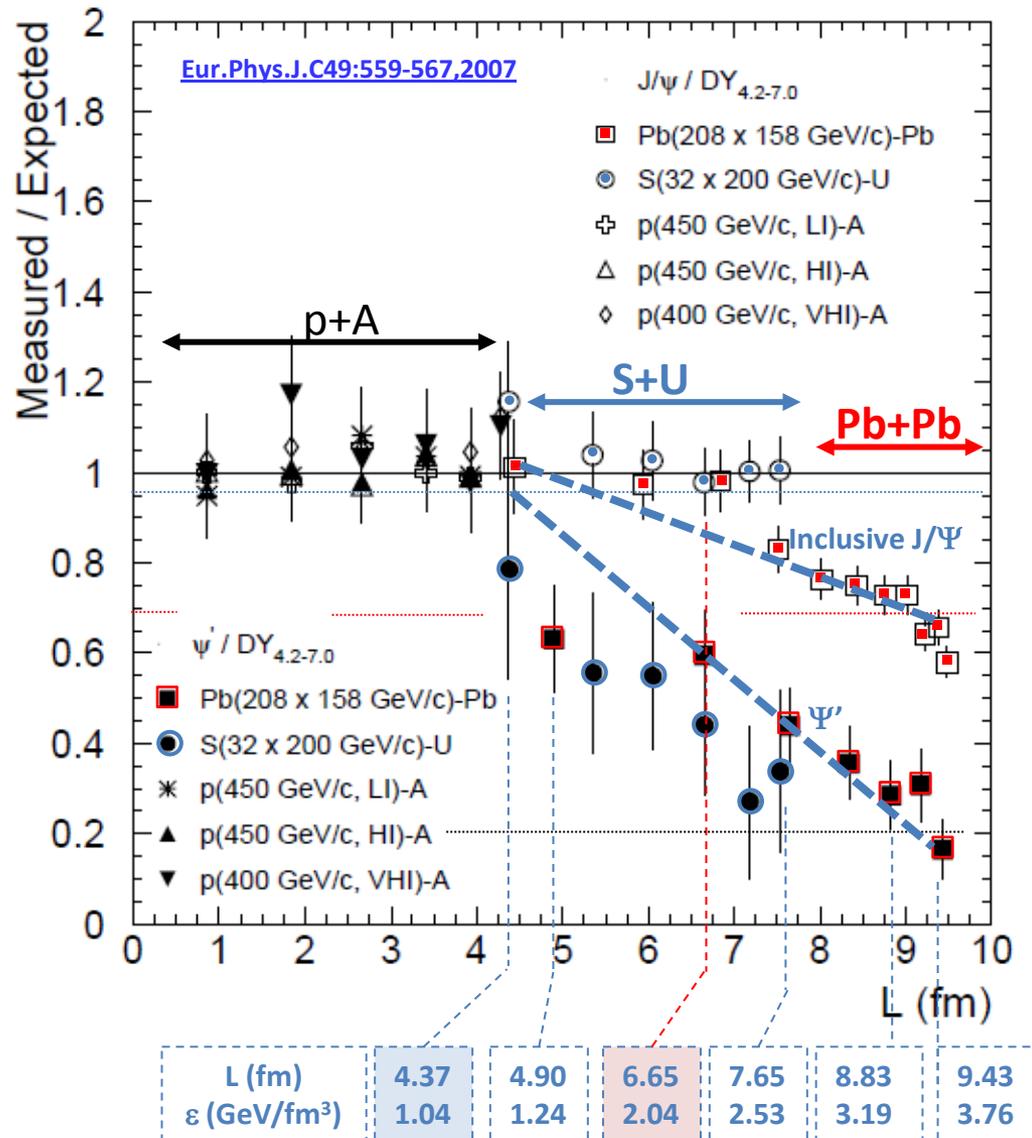
If comovers at work → **smooth suppression**  
(reminder: If QGP at work → **threshold effect**)

Experimentally,

- $\Psi'$  suppression pattern slightly steeper than  $J/\Psi$  one (theoretically  $\sigma_{J/\Psi-co} \leq \sigma_{\Psi'-co}$ )
- If comovers at work,  $\chi_c$  suppression pattern should stand within  $\Psi'$  and  $J/\Psi$  suppression patterns

### Conclusion

Need to measure  $\chi_c$  pattern to test comovers scenario

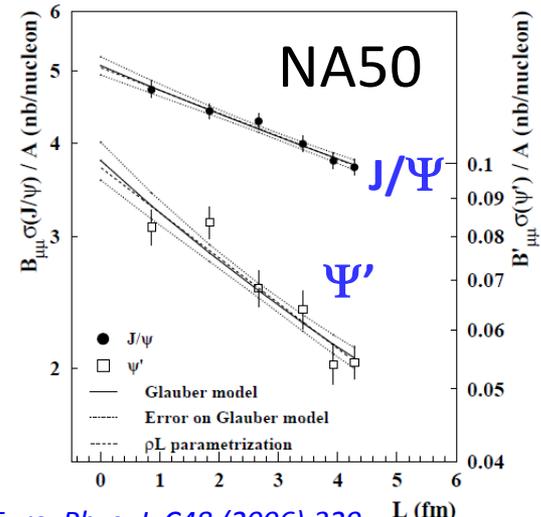


- Benchmark 2: Measure charmonium in p+A at SPS**

$J/\Psi$  and  $\Psi'$  suppression in p+A collisions as a function of L



→ Measuring different charmonium states gives key information on Cold Nuclear Matter and production mechanism.

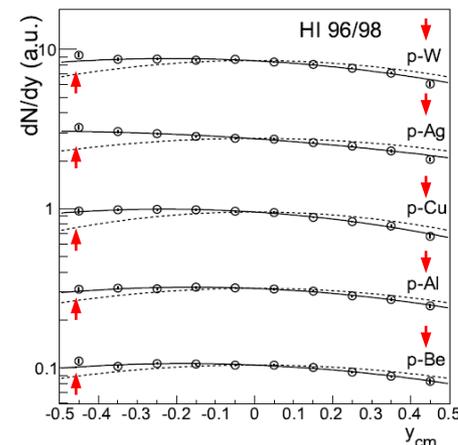


[Euro. Phys. J. C48 \(2006\) 329.](#)

$J/\Psi$  rapidity distribution in p+A collisions (asymetry wrt  $y_{cm}=0$ )



→ Measuring charmonium in a wide  $x_F$  range is important to identify possible (anti)shadowing effects

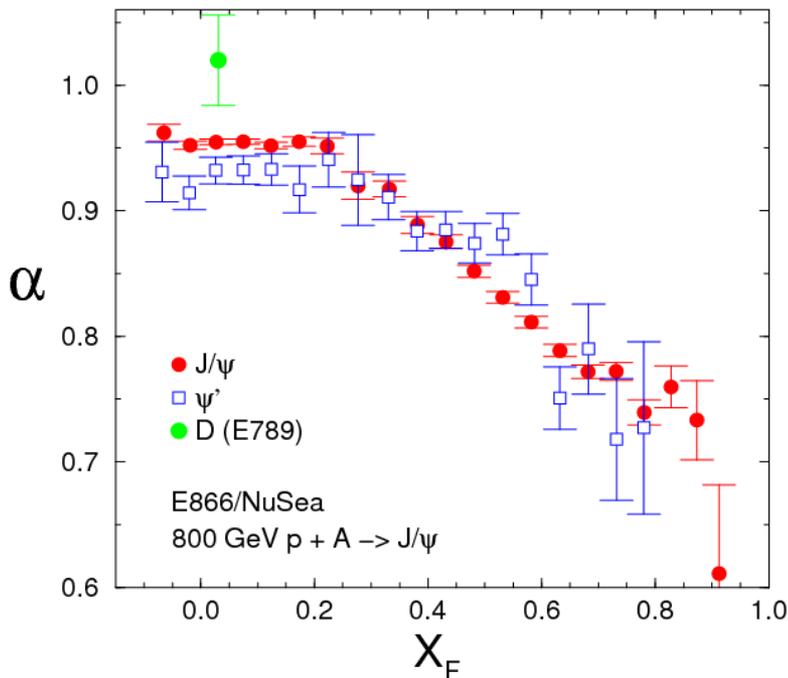


- Measure charmonium in p+A at SPS

→ Measuring charmonium in a wide  $x_F$  range is important to estimate possible (anti)shadowing effects

$$\sigma_A = \sigma_p * A^\alpha$$

E866, Phys. Rev. Lett. 84, 3256-3260 (2000)



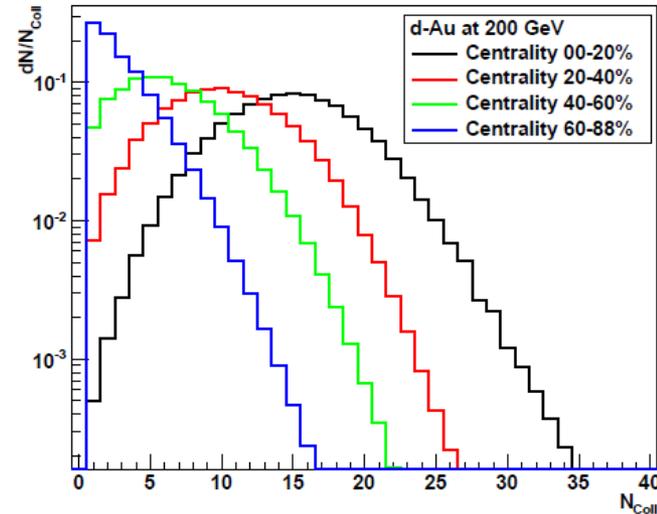
$$x_F = \frac{2M}{\sqrt{s}} \sinh y_{CMS}$$

With  $M=3.1 \text{ GeV}/c^2$  and  $\sqrt{s}=17.2 \text{ GeV}$  (158 GeV)  
 $x_F = 1 \rightarrow y_{CMS} = 1.7$

With  $M=3.1 \text{ GeV}/c^2$  and  $\sqrt{s}=29.1 \text{ GeV}$  (450 GeV)  
 $x_F = 1 \rightarrow y_{CMS} = 2.2$   
 $y_{CMS}=2 \rightarrow x_F = 0.8$

Possible to access large  $x_F$  if measuring charmonia at rapidity up to  $y_{CMS} \sim 2$

- **Cold Nuclear Matter studies**
  - Must be performed in p+A collisions
  - The more A versatility, the better
- **Collider mode**
  - Difficult to operate many A systems (for instance, since 2000, Phenix operated d+Au collisions only) → studies as a function of centrality
  - Constraints:
    1. **Centrality bin limitation:** due to the “small” number of particle produced in p+A, cannot make as many centrality bins as in A+A collisions
    2. **Glauber uncertainty :**  $\langle N_{coll} \rangle$ .vs.centricity through Glauber calculation → uncertainty on  $\langle N_{coll} \rangle$  (~7% for Phenix)
- **Fixed target mode**
  - Easy to operate many A systems
  - No bin limitation
  - No Glauber uncertainties



Phenix d+Au centrality bins  
[arXiv:1204.0777](https://arxiv.org/abs/1204.0777)

Collider mode:

$$R_{pA} = \frac{dN_{pA}^{J/\Psi}}{\langle N_{coll} \rangle dN_{pp}^{J/\Psi}} \times \frac{dN_{pp}^{MB}}{dN_{pA}^{MB}}$$

centrality	$\langle N_{coll} \rangle$
0-20%	$15.1 \pm 1.0$
20-40%	$10.2 \pm 0.7$
40-60%	$6.6 \pm 0.4$
60-88%	$3.2 \pm 0.2$

Fixed target mode: 
$$R_{pA} = \frac{\sigma_{pA}^{J/\Psi}}{A \sigma_{pp}^{J/\Psi}}$$