

# Measurements of $\sin 2\beta$ in $B$ decays

William T. Ford

University of Colorado

FPCP03, 3 June, 2003

- Time evolution of  $B^0$  decays
- Connection with standard model
- “Golden mode” charmonium decays
- Interpretation of measurements in rare decays
- Open charm modes
- Penguin decays
- Summary

## Quark weak couplings in the standard model

Wolfenstein parameterization of the CKM matrix  $V$ :

$$V = \begin{pmatrix} V_{ud} = 1 - \frac{1}{2}\lambda^2 & V_{us} = \lambda & V_{ub} = A\lambda^3(\rho - i\eta) \\ V_{cd} = -\lambda & V_{cs} = 1 - \frac{1}{2}\lambda^2 & V_{cb} = A\lambda^2 \\ V_{td} = A\lambda^3(1 - \rho - i\eta) & V_{ts} = -A\lambda^2 & V_{tb} = 1 \end{pmatrix}$$

$$\begin{array}{rcl} \lambda & \simeq & \sin \theta_c \simeq 0.22 \\ A & \sim & 1 \end{array}$$

The irreducible phase generates  $CP$  non conservation:

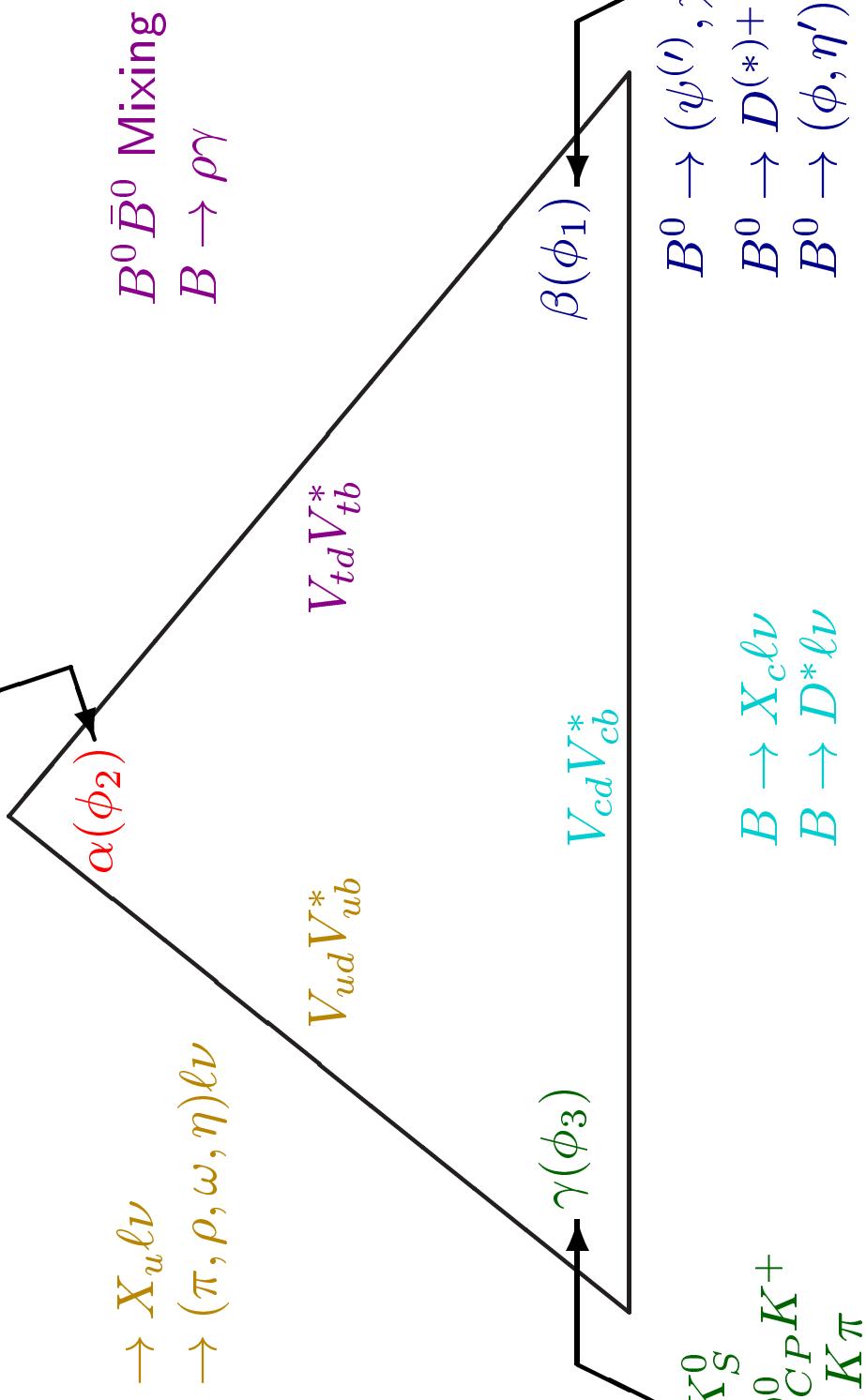
$$(CP)^{-1} HCP = H^* \neq H$$

## Unitarity triangle

$$V^\dagger V = 1 \Rightarrow V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

$$B \rightarrow \pi^+ \pi^-, \rho^+ \pi^-$$

$$\begin{aligned} B &\rightarrow X_u \ell\nu \\ B &\rightarrow (\pi, \rho, \omega, \eta) \ell\nu \end{aligned}$$



$$\begin{aligned} B_S &\rightarrow \rho K_S^0 \\ B^+ &\rightarrow D_{CP}^0 K^+ \\ B &\rightarrow \pi\pi, K\pi \end{aligned}$$

$B^0 \rightarrow (\psi^{(*)}, \chi_c) K_S^0$

$B^0 \rightarrow D^{(*)+} D^{(*)-}$

$B^0 \rightarrow (\phi, \eta') K_S^0$

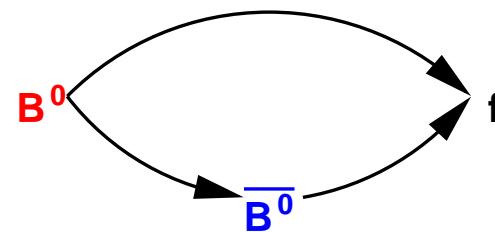
## Time evolution and $CP$ violation

- The decay amplitude for  $B^0 \rightarrow f$  is

$$\langle f | H | B^0_{\text{phys}}(t) \rangle = e^{-imt} e^{-\Gamma t/2} \left[ A_f \cos \frac{1}{2} \Delta m t + i \frac{q}{p} \bar{A}_f \sin \frac{1}{2} \Delta m t \right]$$

$$A_f \equiv \langle f | H | B^0 \rangle, \quad \bar{A}_f \equiv \langle f | H | \bar{B}^0 \rangle \quad |B^0_{L,H}\rangle = p |B^0\rangle \pm q |\bar{B}^0\rangle$$

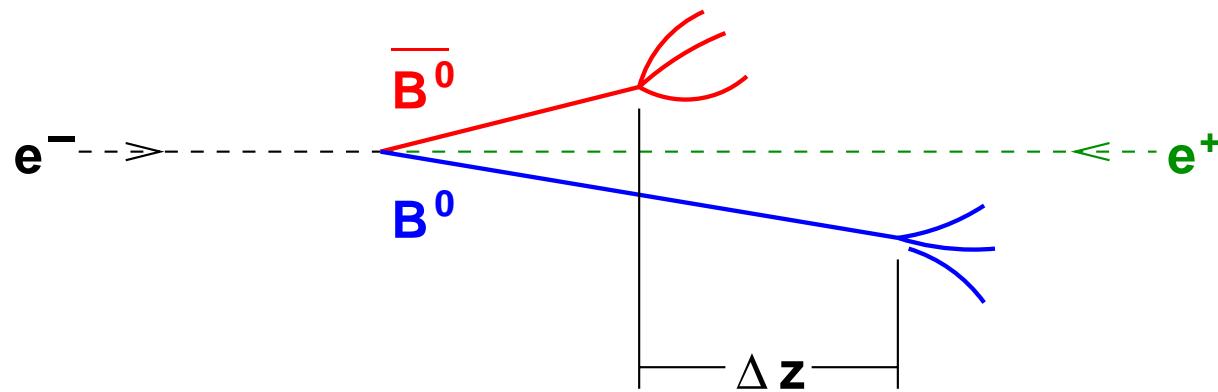
- $CP$  violation appears through the Interference between mixing ( $\frac{q}{p}$ ) and decay ( $\frac{\bar{A}_f}{A_f}$ )



## **$B$ meson pairs from boosted $\Upsilon(4S)$**

In  $\Upsilon(4S)$  decay  $B^0\bar{B}^0$  pair created in a  $C = -1$  eigenstate

These oscillate coherently between  $B^0$  and  $\bar{B}^0$  until one decays  
(Einstein-Podolsky-Rosen effect)



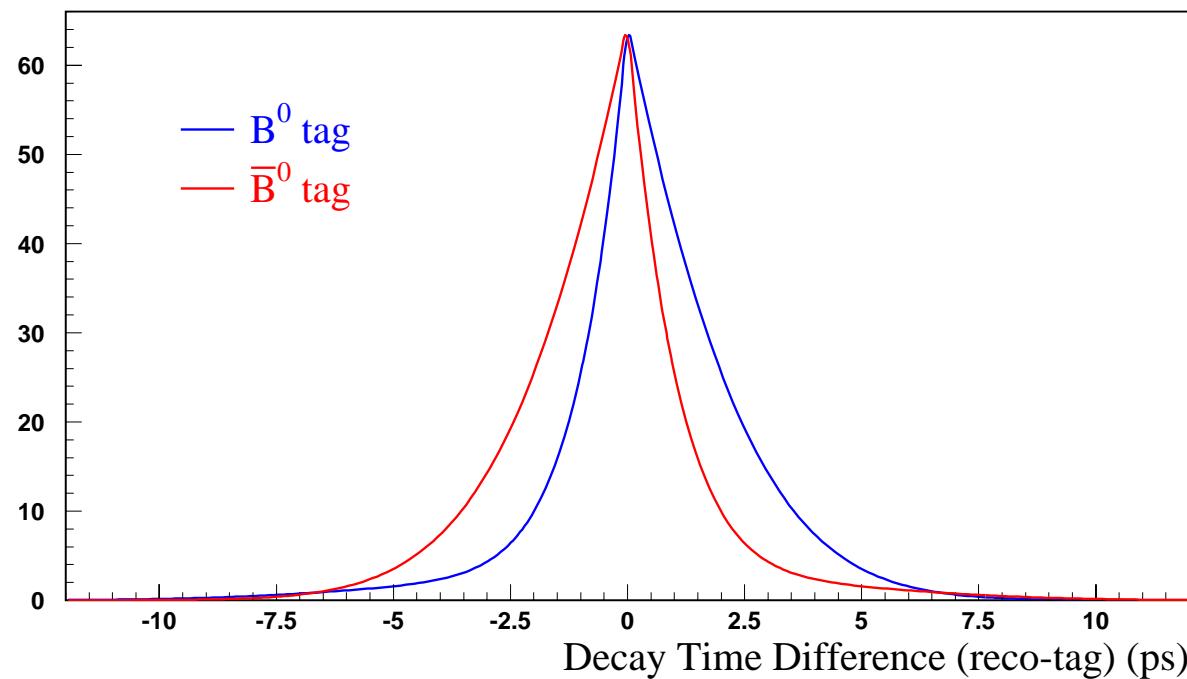
$$\Delta z \simeq \beta\gamma c \Delta t$$

$\beta\gamma = 0.56$  (PEP-II), 0.425 (KEKB)

## Measurement of time evolution: $\Delta t$

Start the  $\Delta t$  clock on the decay of one  $B$  to a flavor eigenstate (“tag”)  
 Stop it on the decay of the other  $B$  to  $CP$  eigenstate  $f$

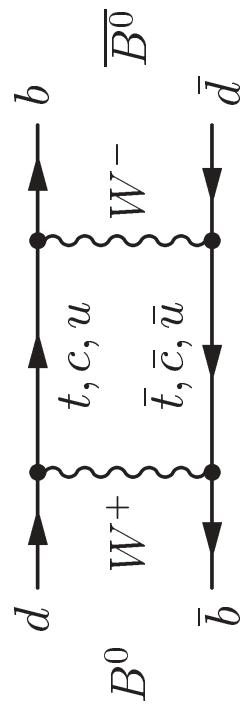
$$\frac{d\Gamma(\Delta t)}{d\Delta t} \propto e^{-|\Delta t|/\tau} (1 \pm \mathcal{I}m\lambda_f \sin \Delta m \Delta t), \quad \lambda_f \equiv \frac{q}{p} \frac{\bar{A}_f}{A_f}, \quad (|\lambda_f| = 1)$$



## The $B\overline{B}$ mixing factor

$$\lambda_f \equiv \frac{q}{p} \frac{\bar{A}_f}{A_f} \quad |B^0_{L,H}\rangle = p|B^0\rangle \pm q|\bar{B}^0\rangle$$

Parton level calculation is reliable (short-distance dominated) for heavy mesons

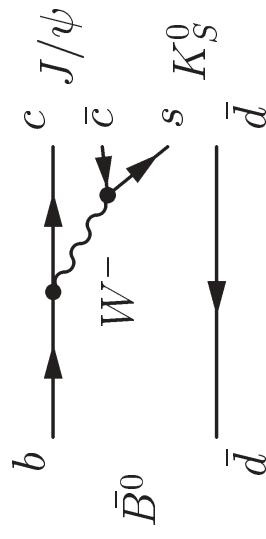


Massive  $t$  quark dominates (frustrated GIM mechanism); in Wolfenstein phase convention

$$\frac{q}{p} = \frac{V_{tb}^* V_{td}}{V_{tb} V_{td}^*} = e^{-2i\beta}, \quad \Rightarrow \lambda_f = e^{-2i\beta} \frac{\bar{A}_f}{A_f}$$

## $\sin 2\beta$ from charmonium $K^{0(*)}$ modes

$$\lambda_f \equiv \frac{q}{p} \frac{\bar{A}_f}{A_f}$$



$$\frac{\bar{A}_f}{A_f} = \eta_f \left( \frac{V_{cb} V_{cs}^*}{V_{cb}^* V_{cs}} \right) \left( \frac{V_{cs}^* V_{cb}}{V_{cs} V_{cb}^*} \right) \left( \frac{V_{cd}^* V_{cs}}{V_{cd} V_{cs}^*} \right) = \eta_f \quad (\text{or } \xi_f)$$

$\eta_f = \pm 1$  for a  $CP$  even/odd final state  $f$ ; for this example,  $\eta_f = -1$

$$\lambda_f = \eta_f e^{-2i\beta}, \quad \mathcal{Im}\lambda_f = -\eta_f \sin 2\beta$$

## Reconstruction of $B$ candidates at the $\Upsilon(4S)$

- $\Upsilon(4S) \rightarrow (B^0 \bar{B}^0, B^+ B^-)$     nearly at rest ( $p_B^* \simeq 325$  MeV/c)
- $m_B \simeq 5.3$  GeV/ $c^2$  =  $m_{\text{recoil}}$   
 $\Rightarrow$

$$E_B^* = E_{\text{beam}}^*, \text{ i.e., } \Delta E \equiv E_B^* - E_{\text{beam}}^* = 0$$

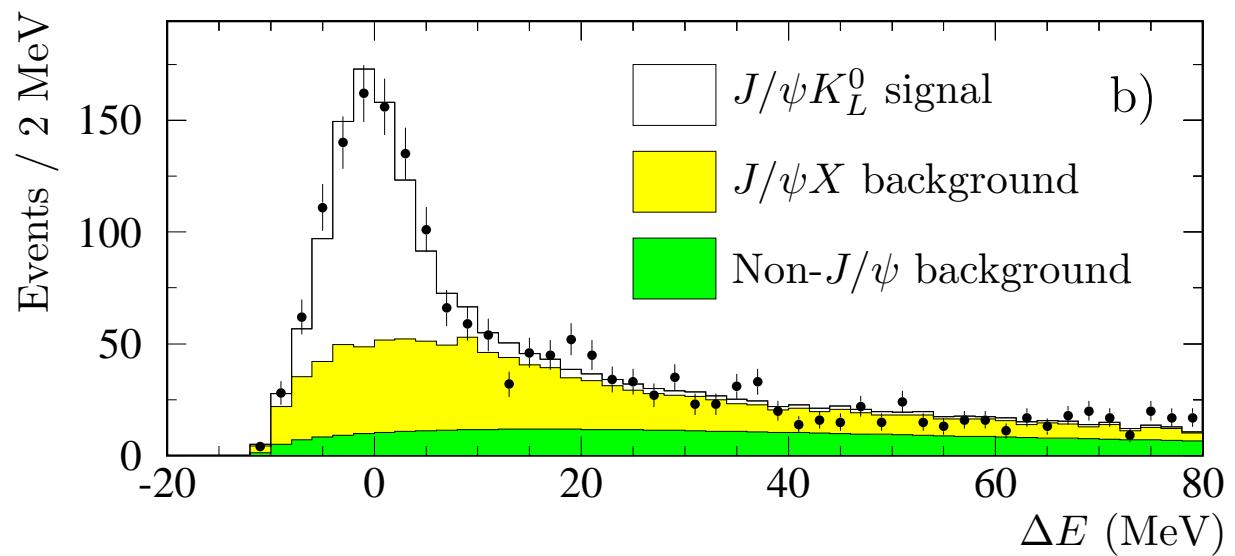
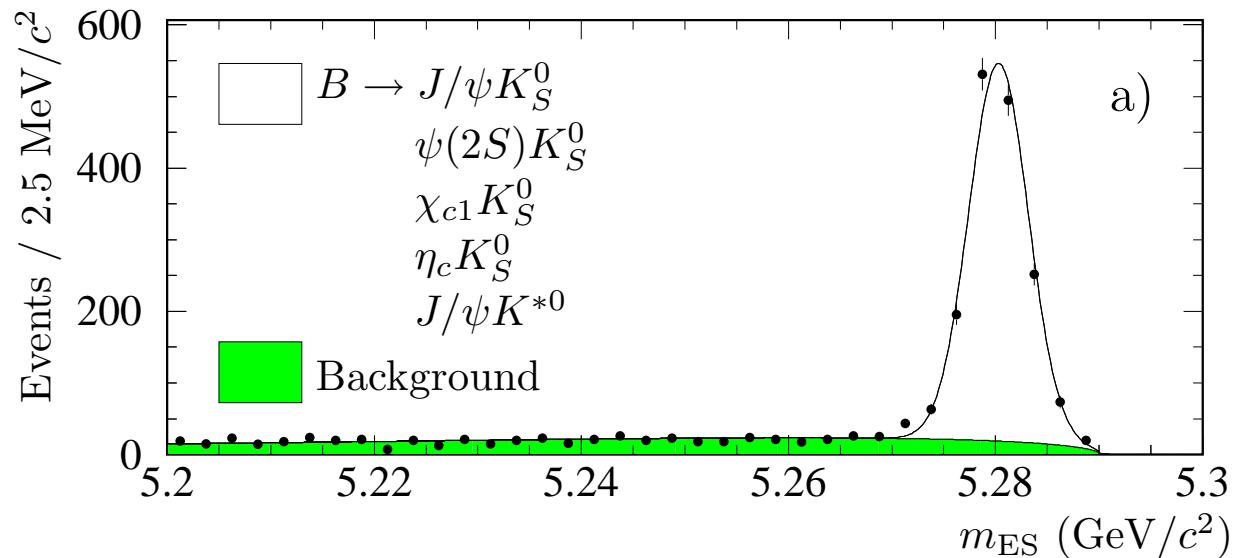
and

$$m_{ES} \equiv \sqrt{E_{\text{beam}}^{*2} - |\mathbf{p}_i^*|^2} = m_B$$

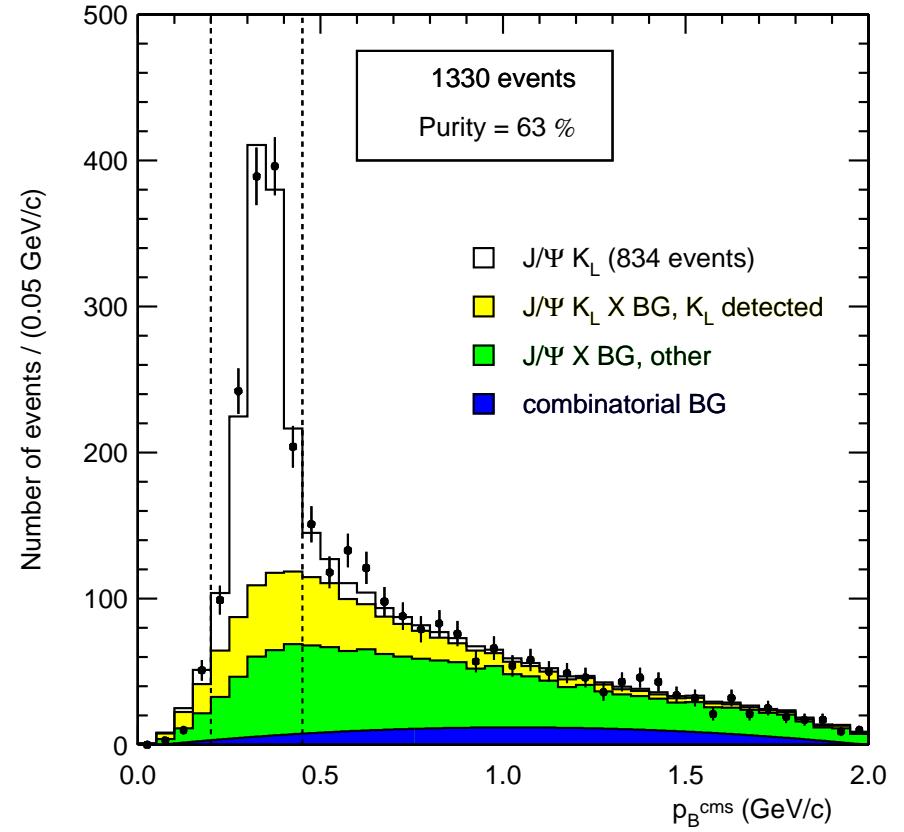
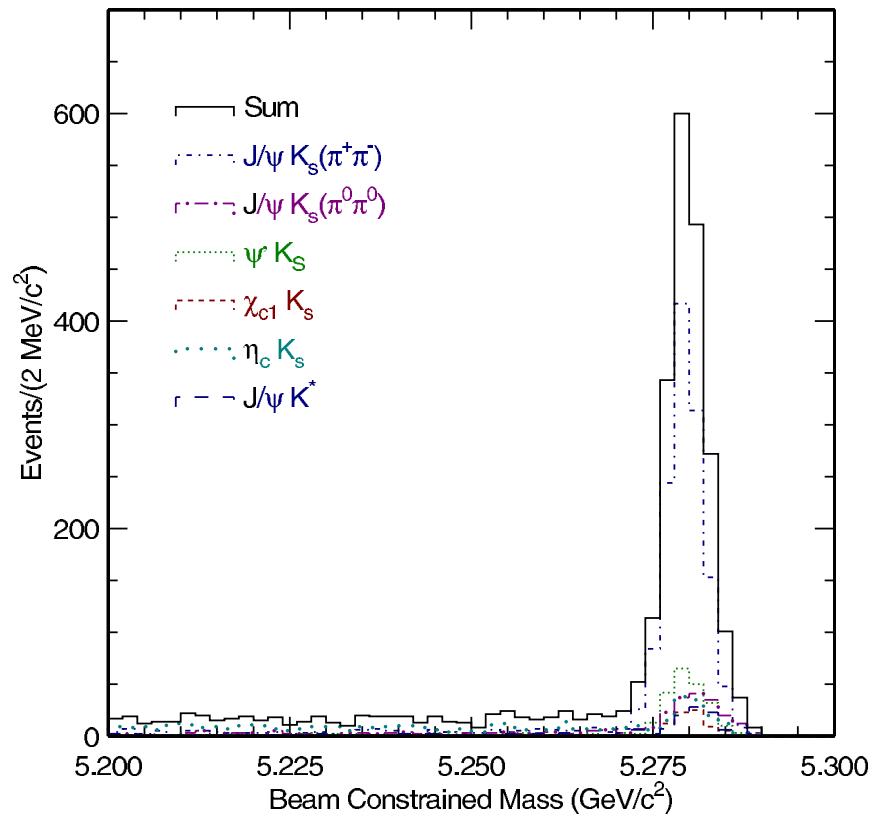
Typical resolution  $m_{ES} : 3$  MeV/ $c^2$ ,    $\Delta E : 15 - 50$  MeV

- For two-body
  - ◊ daughter  $E^* \simeq 2.6$  GeV
  - ◊ daughters nearly back-to-back

Tagged events  
 $(88 \times 10^6 B\bar{B})$



**$B_{CP}$  candidates**

Tagged events ( $85 \times 10^6$  produced  $B\bar{B}$  pairs)

## Flavor of the tag $B$

The other  $B$  is not fully reconstructed, but we need to know whether it's  $B^0$  or  $\overline{B}^0$ .

Tagging signatures:

$$B^0 \rightarrow \ell^+, \overline{B}^0 \rightarrow \ell^-$$

$$B^0 \rightarrow K^+, \overline{B}^0 \rightarrow K^-$$

Inclusive flavor signatures

etc.

Efficiency  $\epsilon$ , mistag rate  $w$  measured with reconstructed flavor eigenstate decays.

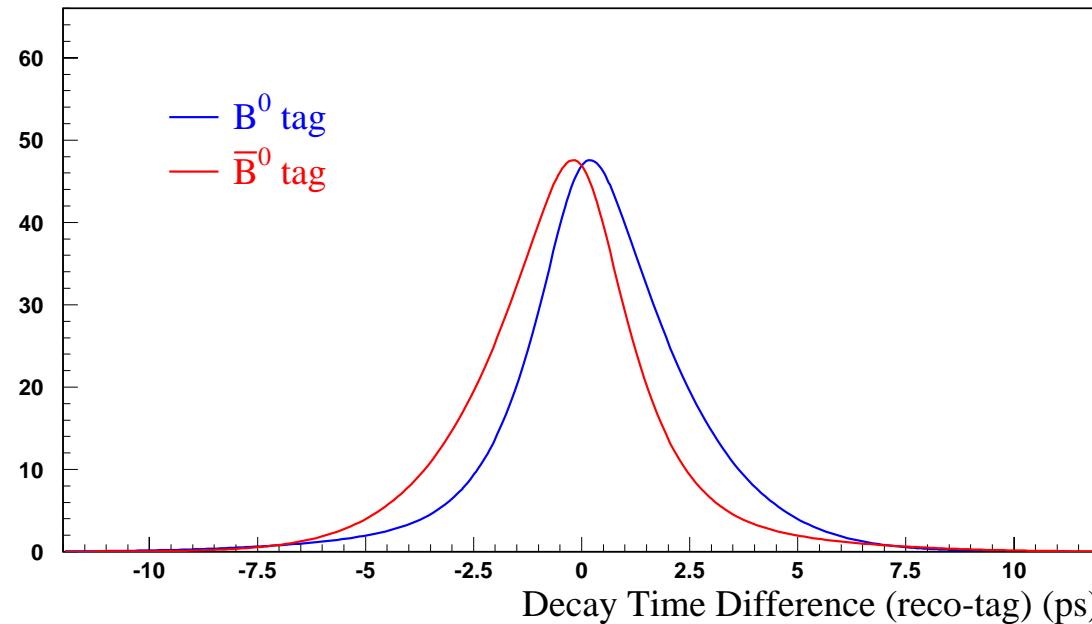
The effective efficiency is

$$Q = \langle \epsilon(1 - 2w)^2 \rangle = (28.1 \pm 0.7)\% \text{ (BABAR)}, (28.8 \pm 0.6)\% \text{ (Belle)}$$

## Decay rate with resolution and realistic tagging

$$\frac{1}{\Gamma} \frac{d\Gamma(\Delta t)}{d\Delta t} = \frac{e^{-|\Delta t|/\tau}}{4\tau} (1 \pm (1 - 2w) \mathcal{I}m\lambda_f \sin \Delta m \Delta t) \otimes \mathcal{R}$$

Vertex resolution (largest contribution from tag side)  $\simeq 180 \text{ }\mu\text{m}$ ,  
or  $\sim 1.25 \text{ ps}$



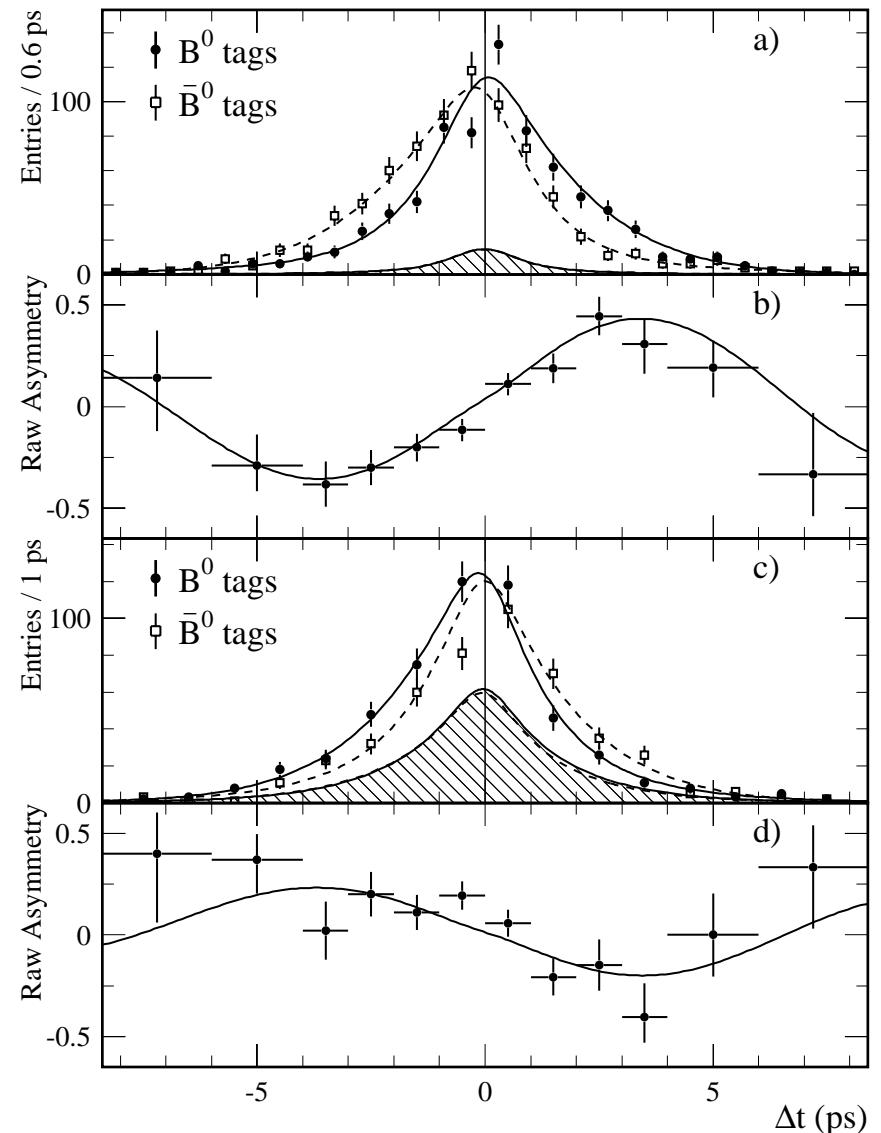
## Time development fit for $B_{CP}$



$CP$  Odd modes, with asymmetry, above  
 $CP$  Even  $K_L^0$  mode below  
 34 parameter likelihood fit

$$\sin 2\beta = 0.741 \pm 0.067 \pm 0.034$$

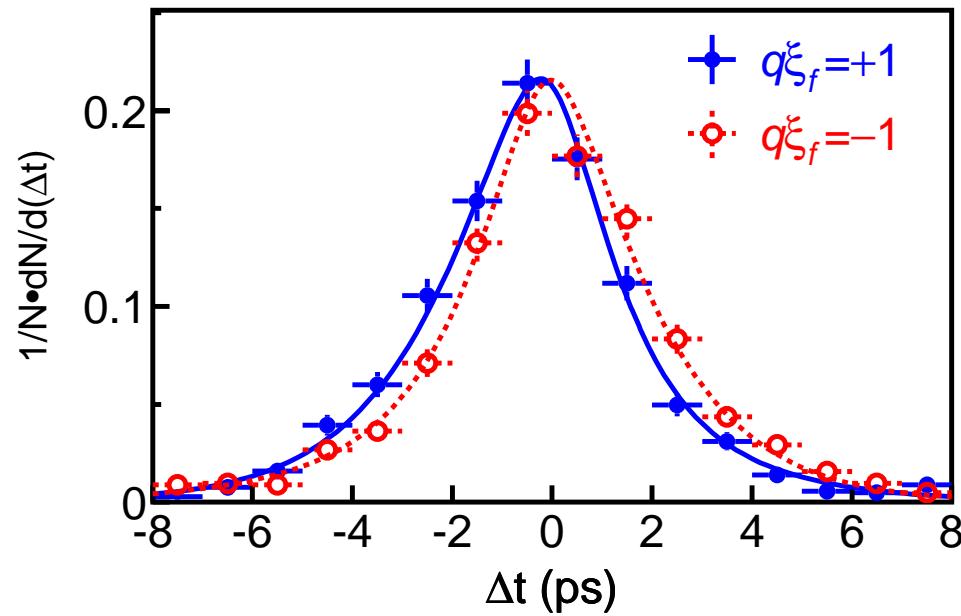
2641 tagged events, 78% purity  
 $(88 \times 10^6 B\bar{B}$  pairs)



## Time development fit for $B_{CP}$

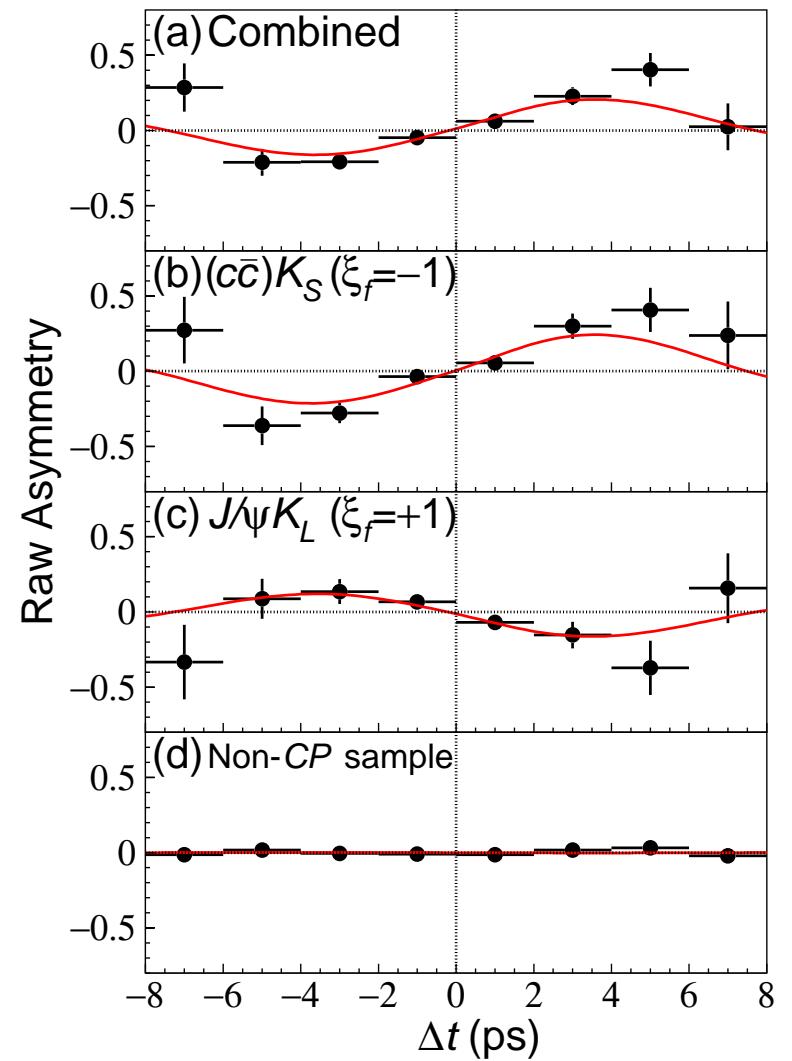


$q = \text{tag sign}, B^0 = +$

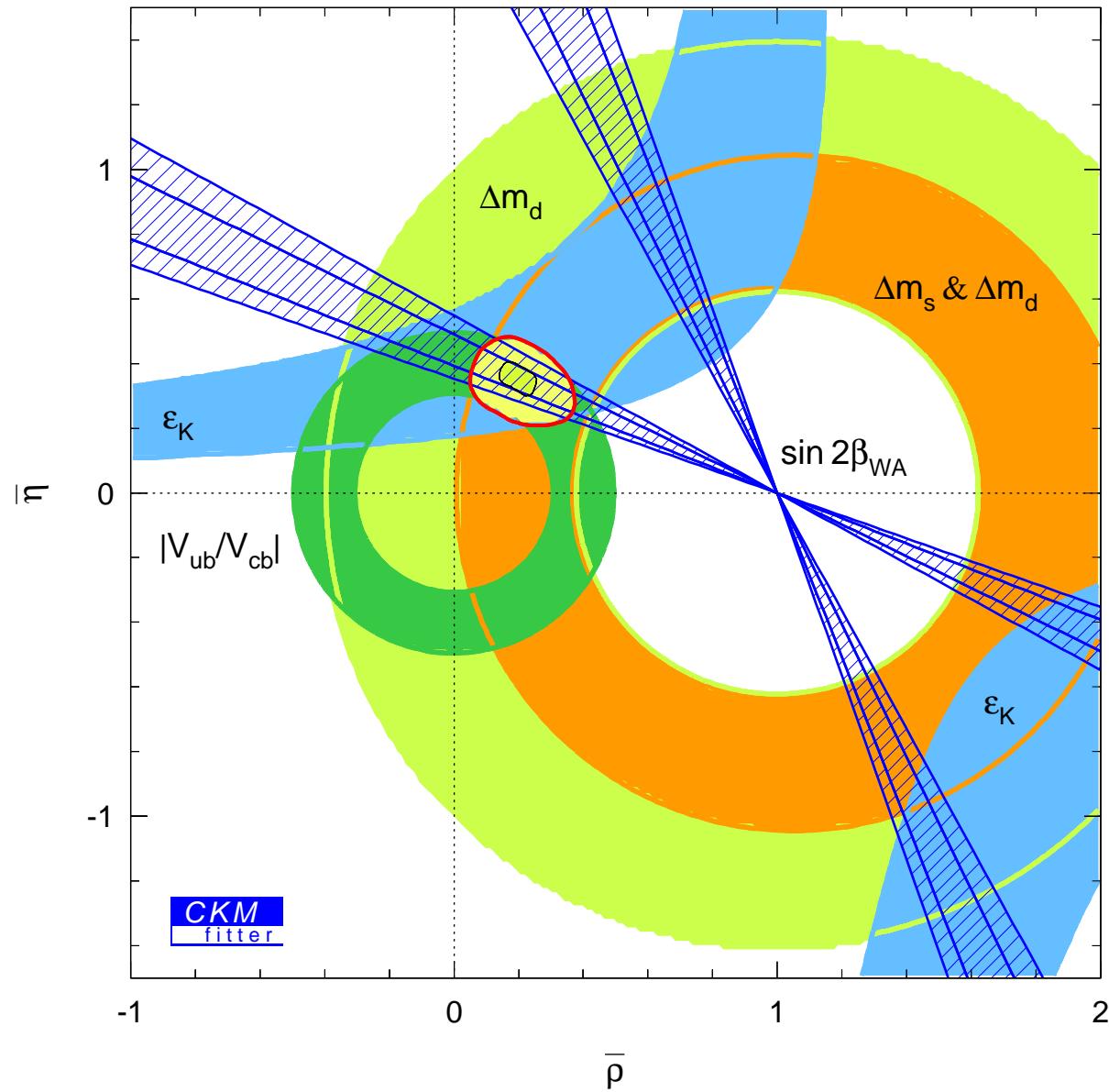


$$\sin 2\beta = 0.719 \pm 0.074 \pm 0.035$$

2958 tagged events, 81% purity  
 $(85 \times 10^6 B\bar{B}$  pairs)



## Unitarity triangle and CP measurements



## Further Investigations

- Rarer  $\mathcal{O}(\lambda^3)$   $B$  decays
  - ◊  $b \rightarrow c\bar{c}d$  (Cabibbo-suppressed; charmonium  $\pi^0$ , open charm pair)
  - ◊  $b \rightarrow s\bar{q}q$  (gluonic penguin;  $\phi K_S^0$ ,  $\eta' K_S^0$ )
- Sensitive to new physics
  - ◊ Smaller amplitudes may reveal NP through interference terms
  - ◊ Virtual particles (e.g., SUSY) in penguin loops
- These experiments are harder
  - ◊ Lower rates, higher backgrounds
  - ◊ tree, (multi-) penguin amplitudes complicate interpretation
  - ◊ Uncertainties from short-distance effects

For these decays we remove the assumption  $|\lambda_f| = 1$ ,

Cast the decay time dependence in terms of

sine-like ( $S_f$ ) and cosine-like ( $C_f$ ) coefficients.

For final  $CP$  eigenstate  $f$

(defining  $\Delta w = w(B^0) - w(\overline{B}^0)$ , still assuming  $\Gamma_H - \Gamma_L = 0$ )

$$\frac{d\Gamma(\Delta t)}{d\Delta t} \propto \frac{e^{-|\Delta t|/\tau}}{4\tau} [1 \mp \Delta w \pm (1 - 2\langle w \rangle) (S_f \sin(\Delta m_d \Delta t) - C_f \cos(\Delta m_d \Delta t))]$$

where

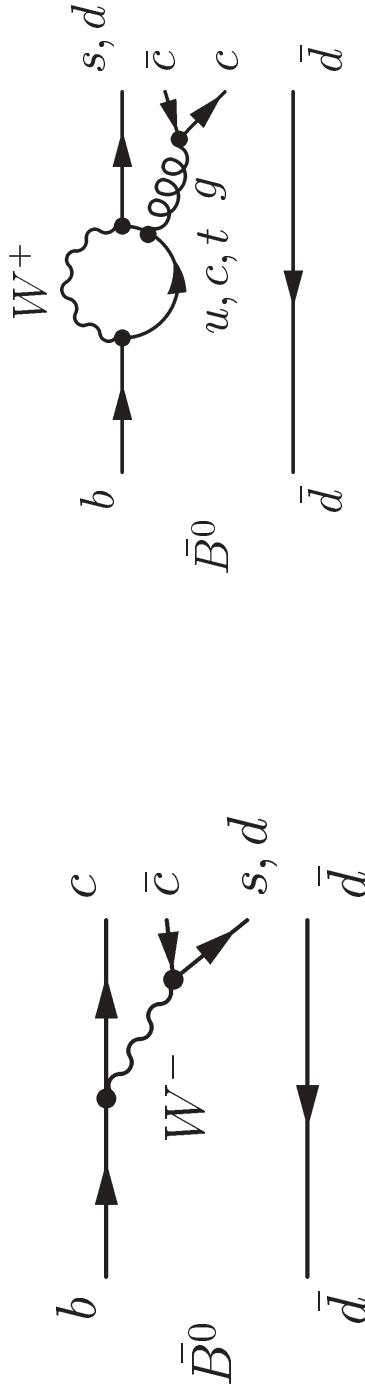
$$S_f = \frac{2\mathcal{I}m\lambda_f}{1+|\lambda_f|^2} (= \sin 2\beta \text{ for } B^0 \rightarrow \psi K_S^0)$$

and

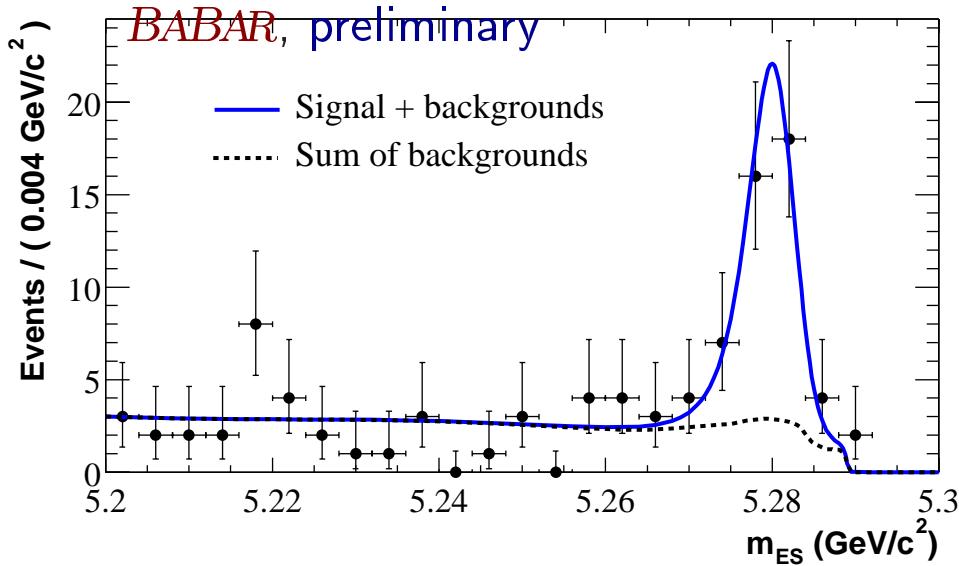
$$-A_f = C_f = \frac{1-|\lambda_f|^2}{1+|\lambda_f|^2} (= 0 \text{ for } B^0 \rightarrow \psi K_S^0)$$

## $b \rightarrow c\bar{c}d$ decays

- In  $b \rightarrow c\bar{c}(s, d)$  the color-suppressed tree competes with penguins having in the loop:
  - ◊  $c - t$  (same  $CP$  phase as the tree)
  - ◊  $u - t$  (different  $CP$  phase)



- For  $b \rightarrow c\bar{c}s$  (e.g.,  $J/\psi K_S^0$ ),  $(u, t)$  penguin/tree =  $\mathcal{O}(\lambda^4/\lambda^2)$
- For (also Cabibbo-suppressed)  $B \rightarrow c\bar{c}d$  (e.g.,  $J/\psi \pi^0$ ), both are  $\mathcal{O}(\lambda^3)$
- The P/T ratio becomes a major theoretical systematic for interpretation of  $b \rightarrow c\bar{c}d$  decays.



$$S_{J/\psi\pi^0} = 0.05 \pm 0.49 \pm 0.16$$

$$C_{J/\psi\pi^0} = 0.38 \pm 0.41 \pm 0.09$$

$(40 \pm 7 \text{ signal events, } 88 \times 10^6 B\bar{B} \text{ pairs})$

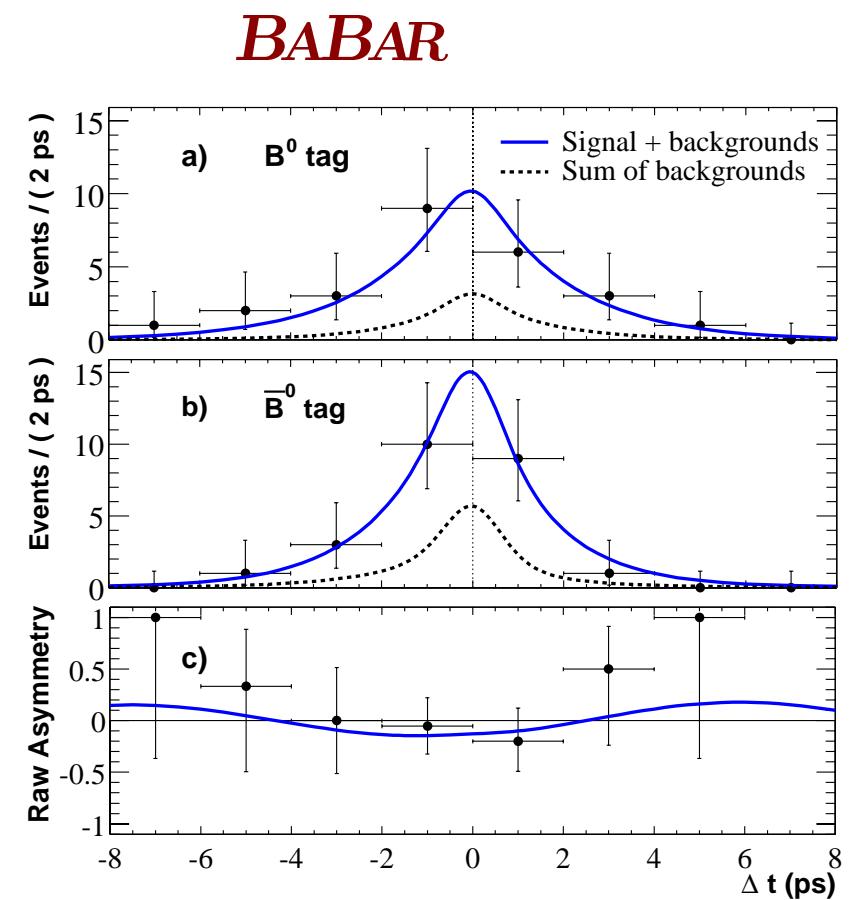
 **preliminary**

$$S_{J/\psi\pi^0} = 0.93 \pm 0.49 \pm 0.11^{+0.27}_{-0.03}$$

$$-C_{J/\psi\pi^0} = \mathcal{A} = -0.25 \pm 0.39 \pm 0.06$$

$(57 \text{ total events, } 86\% \text{ purity, } 85 \times 10^6 B\bar{B})$

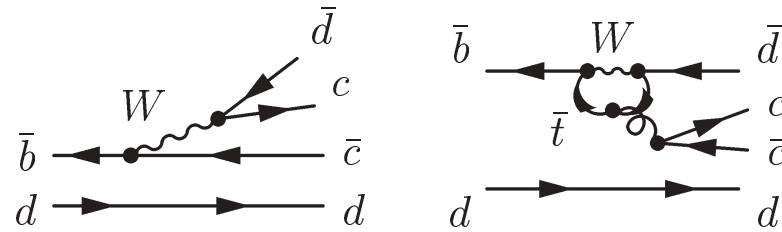
## sin2 $\beta$ from $B^0 \rightarrow J/\psi\pi^0$



Both measurements consistent with  $S_{J/\psi\pi^0} = \sin 2\beta$ ,  $C_{J/\psi\pi^0} = 0$ , within large errors.

$$B^0 \rightarrow D^{*\pm} D^\mp$$

The penguin with  $t$  or  $u$  brings in a second weak phase, . . .

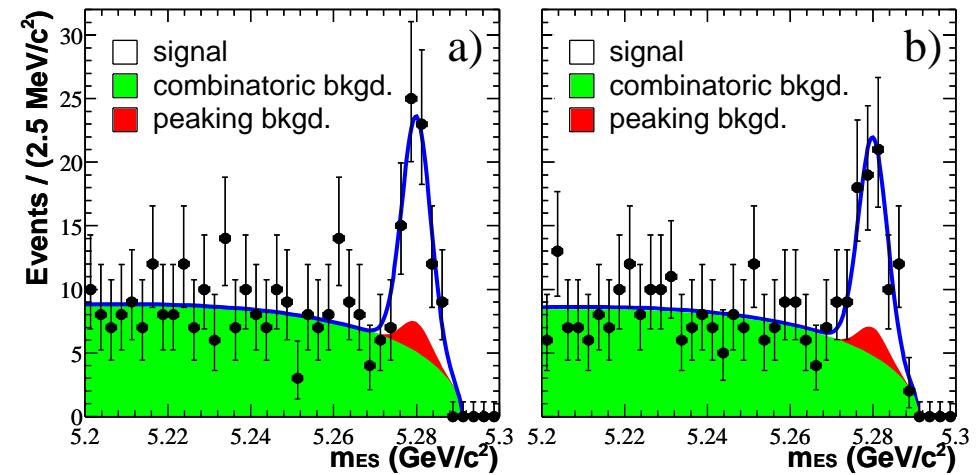


. . . expected to be small ( $\Delta\beta \sim 0.1$ ).  
 These are not  $CP$  eigenstates, but are accessible from  $B^0$  and  $\bar{B}^0$ .

***BABAR***, preliminary

$$D^{*-} D^+$$

$$D^{*+} D^-$$



$$D^{*\pm} \rightarrow \pi^\pm D^0, (4 \text{ } D^0 \text{ modes})$$

$$D^+ \rightarrow K\pi\pi, K_S^0\pi$$

## sin2 $\beta$ from $B^0 \rightarrow D^{*\pm}D^\mp$

**BABAR** (preliminary,  $113 \pm 13$  signal events,  $88 \times 10^6 B\bar{B}$  pairs)

With notation  $S_{+-}$  for  $D^{*+}D^-$ , etc.,

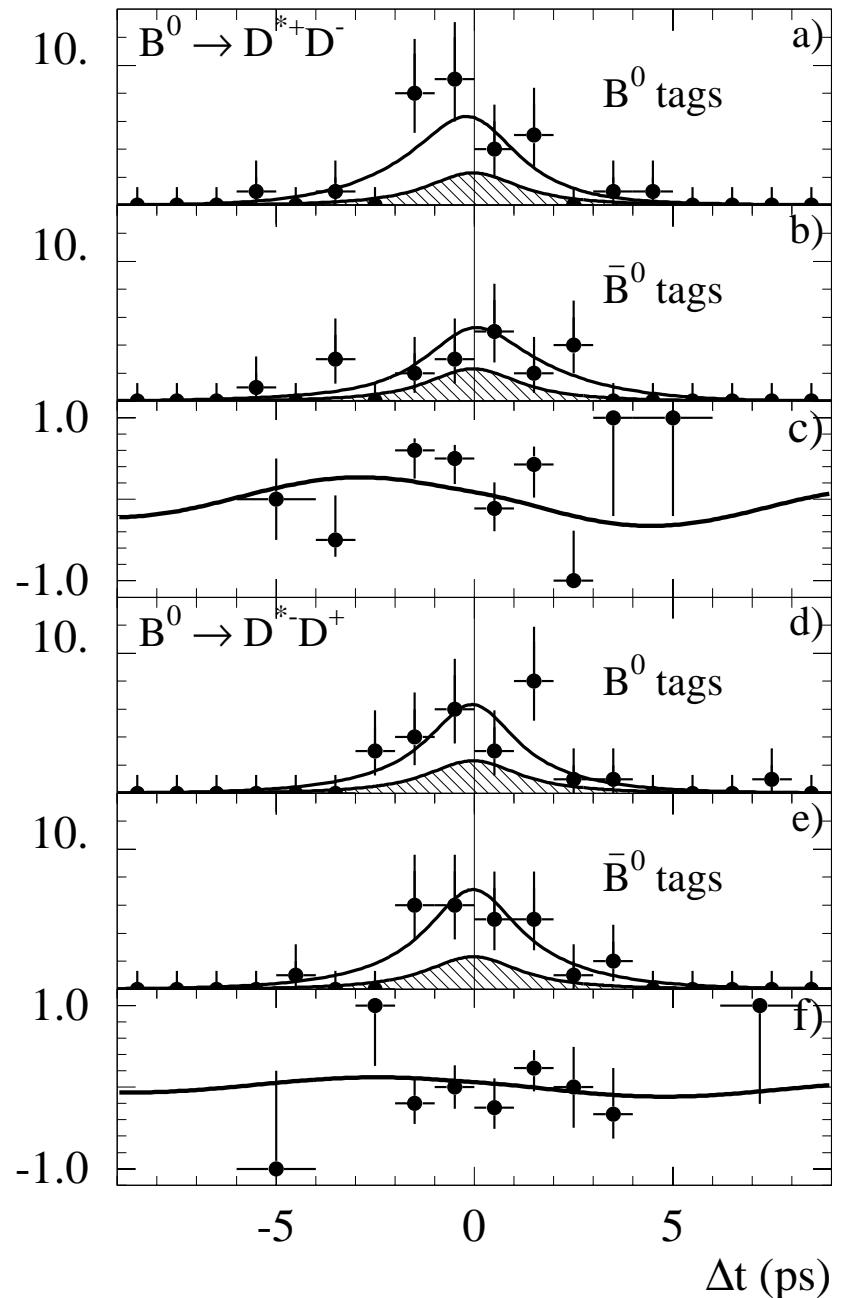
$$S_{-+} = -0.24 \pm 0.69 \pm 0.12$$

$$C_{-+} = -0.22 \pm 0.37 \pm 0.10$$

$$S_{+-} = -0.82 \pm 0.75 \pm 0.14$$

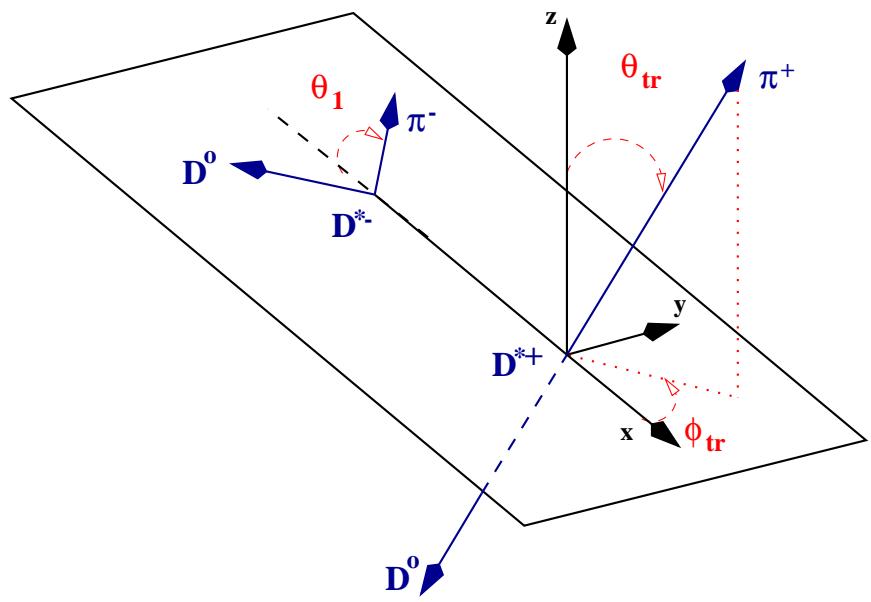
$$C_{+-} = -0.47 \pm 0.40 \pm 0.12$$

If equal amplitudes for  $D^{*-}D^+$ ,  $D^{*+}D^-$ , expect  $C_{-+} = C_{+-} = 0$   
and if penguins negligible,  
 $S_{-+} = S_{+-} = -\sin 2\beta = -0.7$



$$B^0 \rightarrow D^{*\pm} D^{*\mp}$$

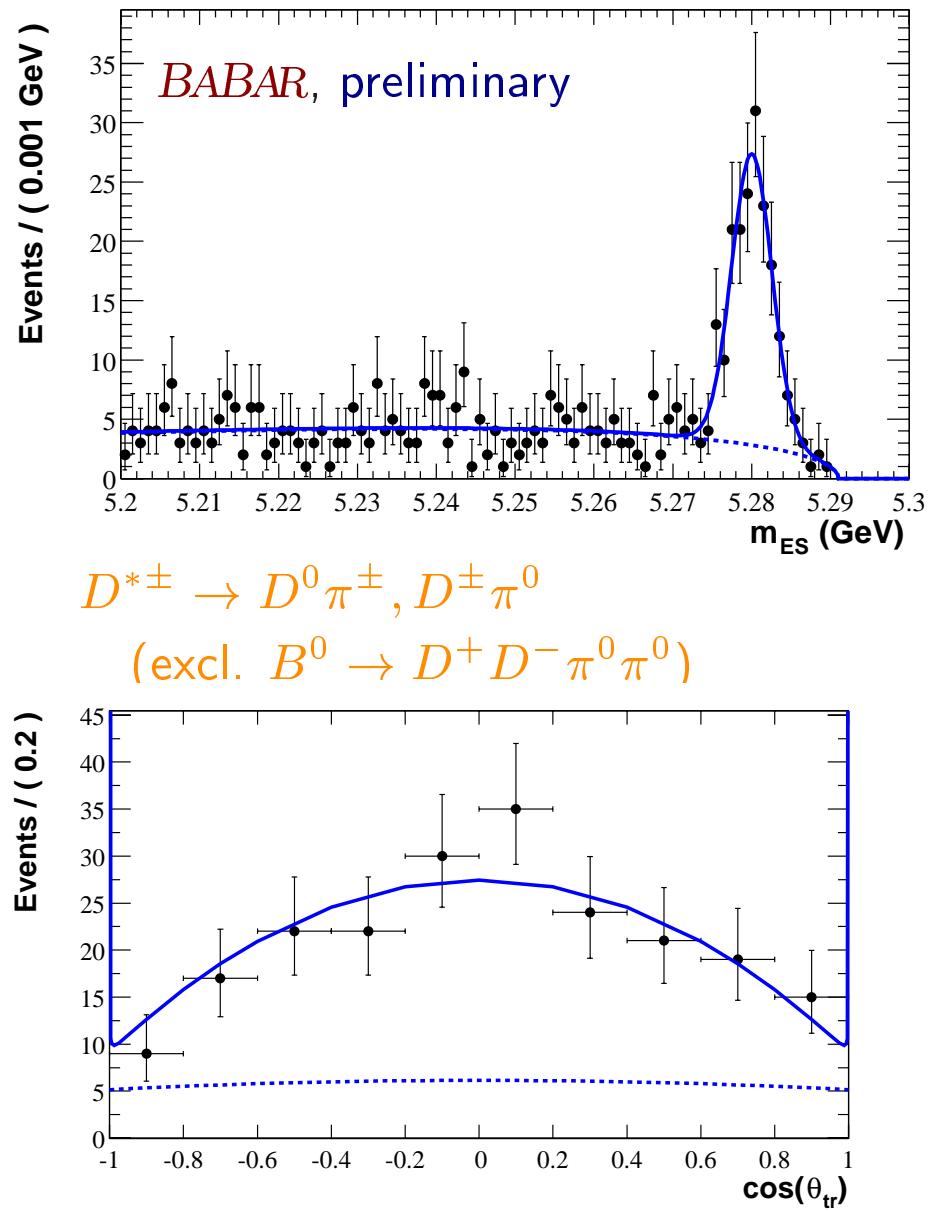
Vector-vector decay, with  $S, P, D$ -wave contributions.



Fit to  $\theta_{tr}$  yields  $CP$ -odd component  
 $R_\perp$

$$R_\perp = 0.063 \pm 0.055 \pm 0.009$$

$\Rightarrow \sim 94\% CP$  even



## $\sin 2\beta$ from $B^0 \rightarrow D^{*\pm} D^{*\mp}$

**BABAR**, preliminary

New, updated from ICHEP02

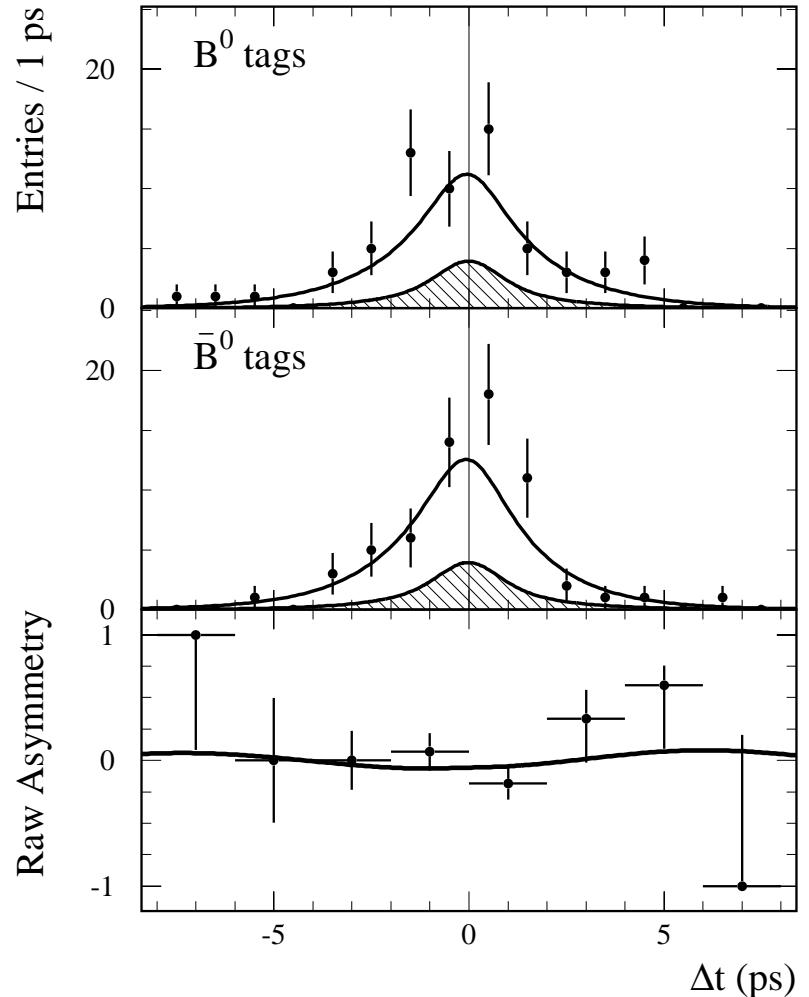
$156 \pm 14$  signal events (before tagging),  
 73% purity,  $88 \times 10^6 B\bar{B}$  pairs

Defining  $\lambda_{f+}$  for  $CP$ -even component

$$\text{Im}\lambda_{f+} = 0.05 \pm 0.29 \pm 0.10$$

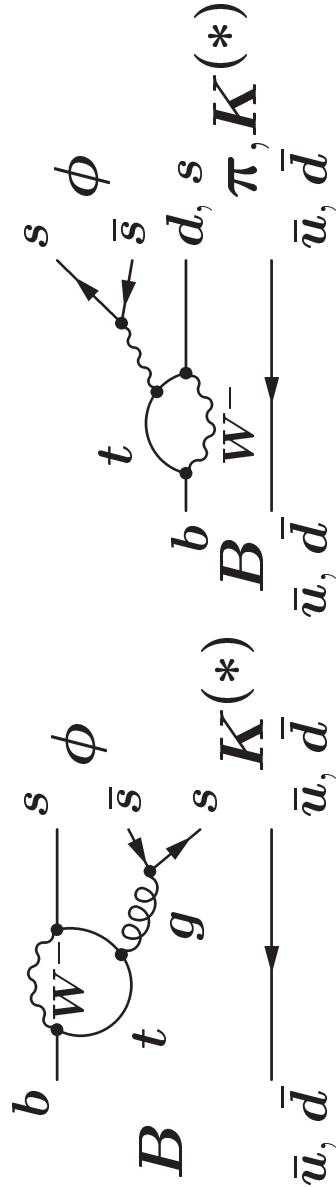
$$|\lambda_{f+}| = 0.75 \pm 0.19 \pm 0.02$$

Compare with tree-level expectation  $|\lambda_{f+}| = 1$ ,  $\text{Im}\lambda_{f+} = -\sin 2\beta$



$$B \rightarrow \phi K_s^0$$

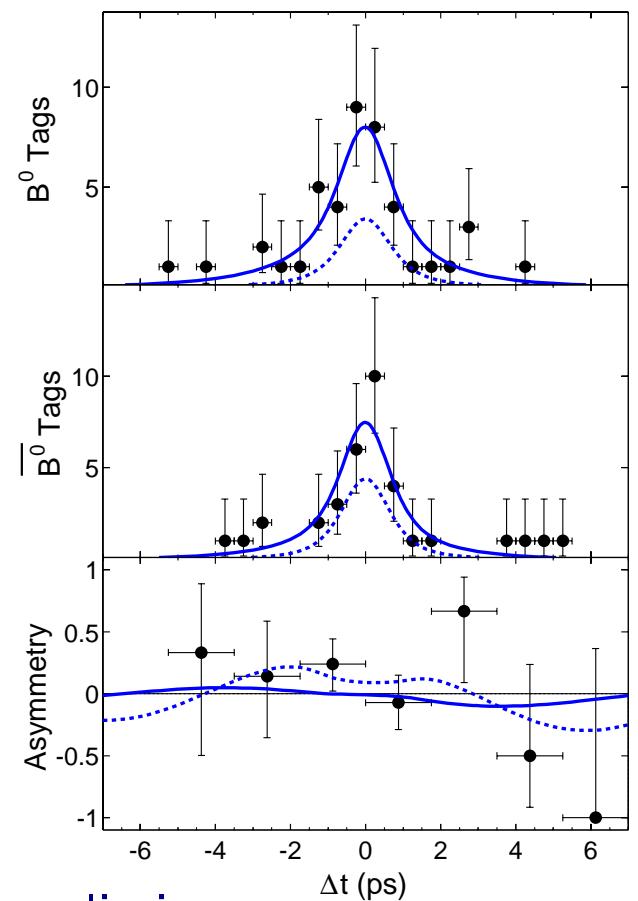
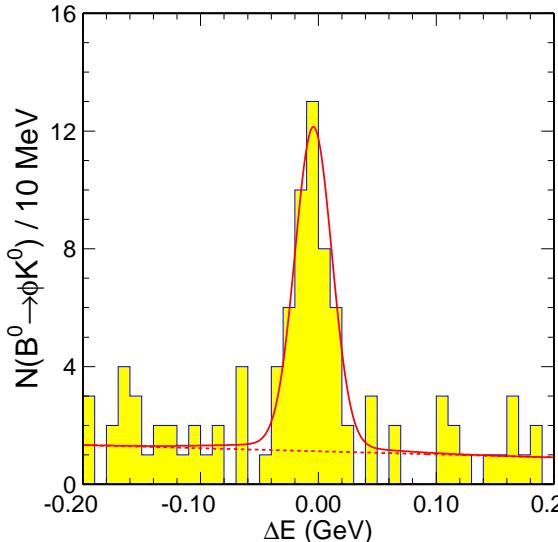
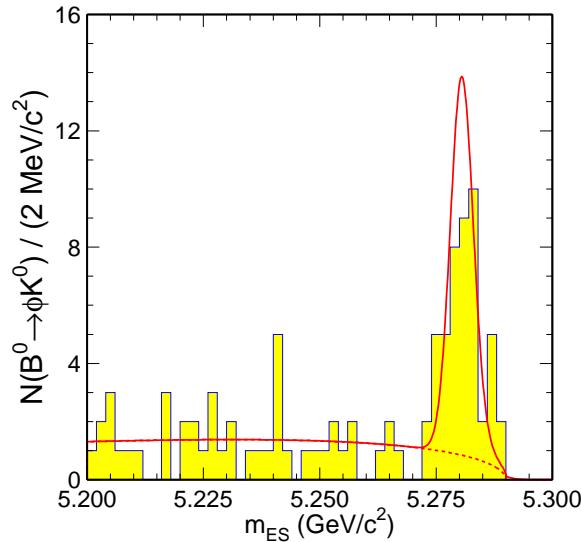
Pure  $b \rightarrow s\bar{s}s$  transition; Gluonic penguins dominate:



No tree, and like  $b \rightarrow c\bar{c}s$ ,  $u$ -loop with different weak phase is suppressed by  $\mathcal{O}(\lambda^2)$  (but here there is no penguin suppression),  $\Rightarrow$

**Deviation of  $S_{\phi K_S^0}$  from  $\sin 2\beta \Rightarrow$  new physics**

The naive estimate can eventually be replaced by bounds from  $SU(3)$  relations to channels not yet measured.  
(see Grossman, Ligeti, Nir, Quinn, hep-ph/0303171)

$$B^0 \rightarrow \phi K_S^0$$


preliminary

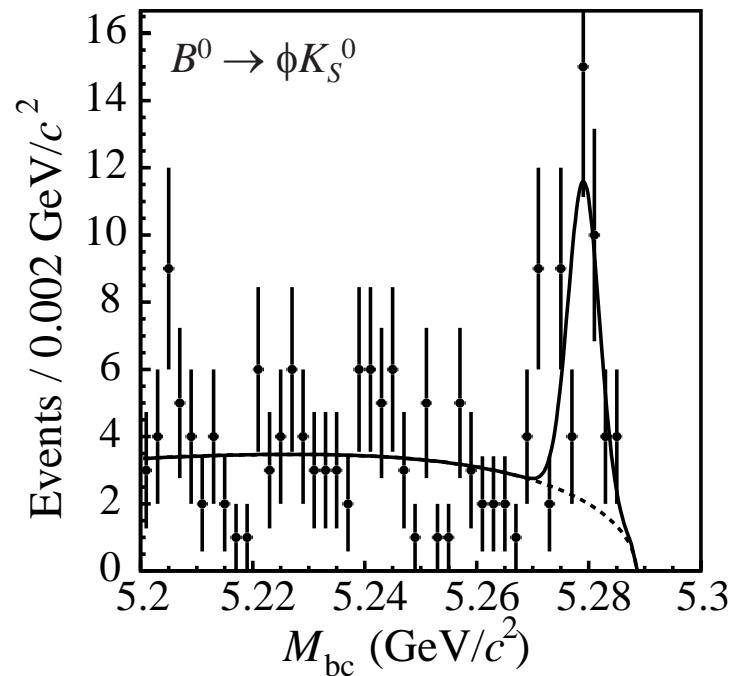
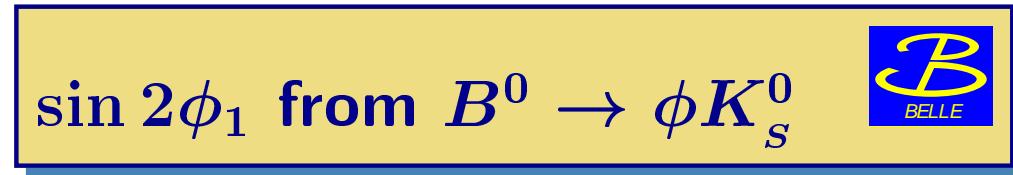
51 signal events (31 tagged),  $87 \times 10^6 B\bar{B}$  pairs

$$S_{\phi K_S^0} = -0.18 \pm 0.51 \pm 0.07$$

$$C_{\phi K_S^0} = -0.80 \pm 0.38 \pm 0.12$$

$$\text{fix } C_{\phi K_S^0} = 0, \quad S_{\phi K_S^0} = -0.26 \pm 0.51 \text{ (stat)}$$

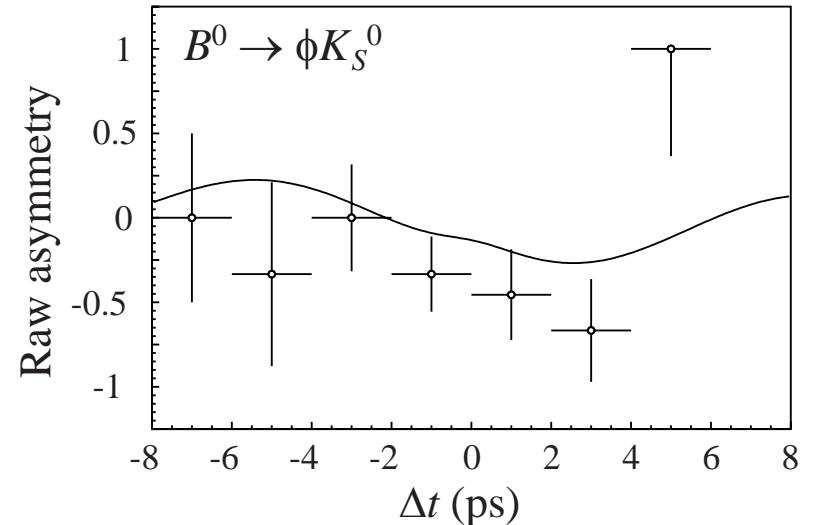
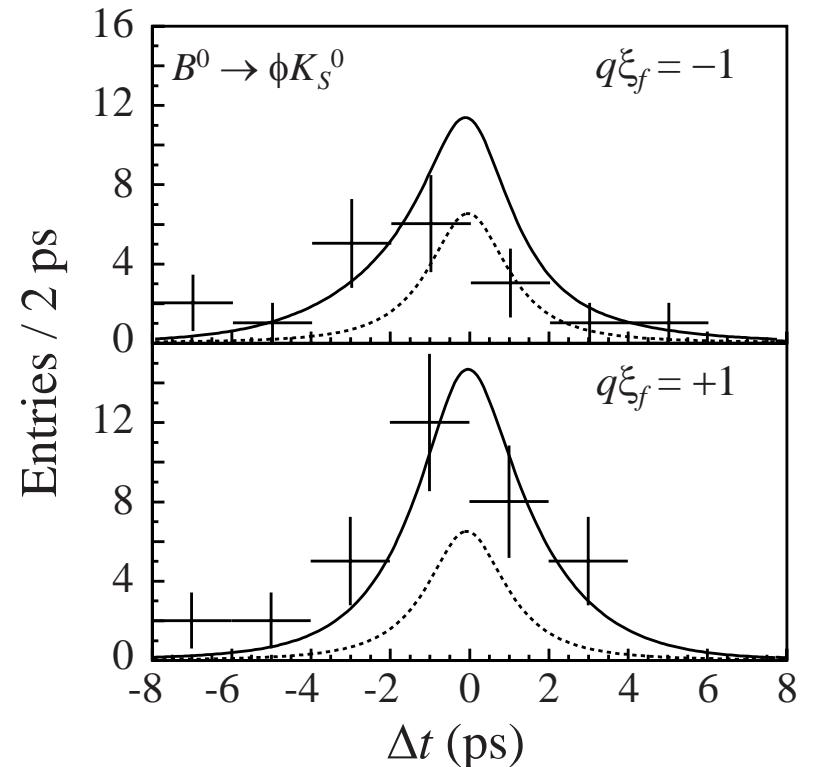
( $\sim$  null result, though consistent within errors with  $S_{\phi K_S^0} = \sin 2\beta = 0.7$ )

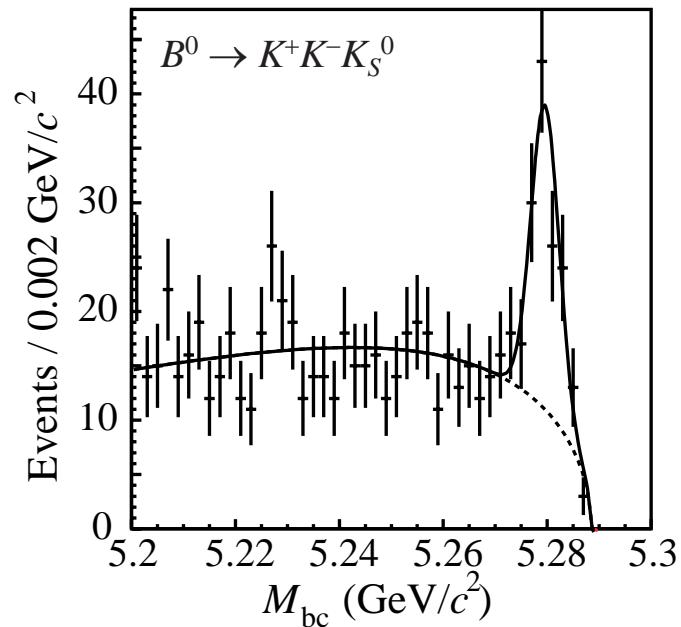
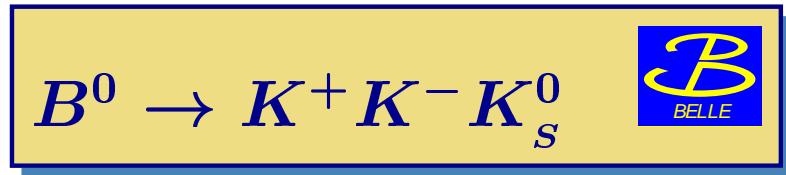


53 total events, purity 67%,  $78 \times 10^6 B\bar{B}$

$$\sin 2\beta = S_{\phi K_S^0} = -0.73 \pm 0.64 \pm 0.22$$

$$\mathcal{A}_{\phi K_S^0} = -C_{\phi K_S^0} = -0.56 \pm 0.41 \pm 0.16$$



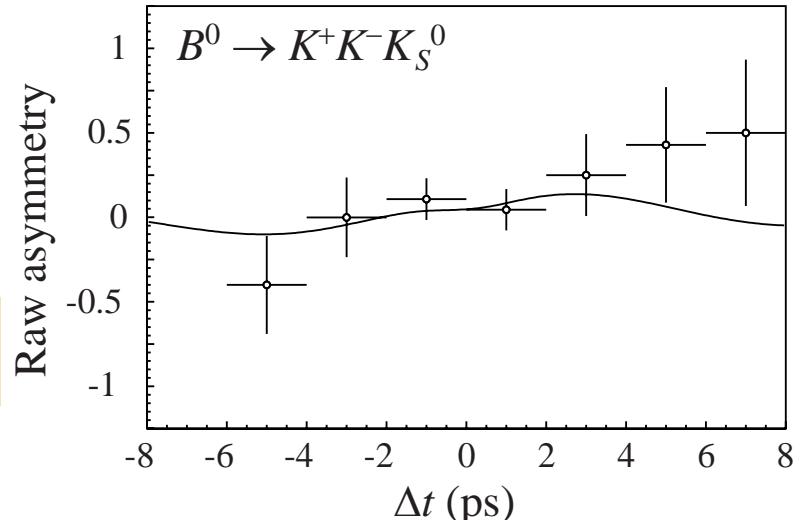
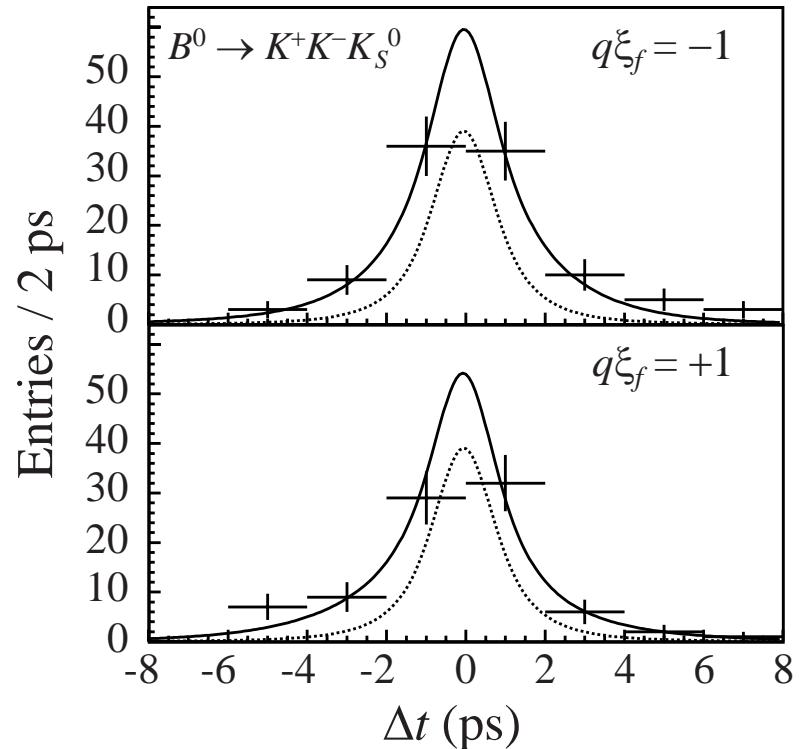


191 total events, purity 50%,  $78 \times 10^6 B\bar{B}$

Dalitz analysis  $\Rightarrow \xi_f = +1$  (97%)

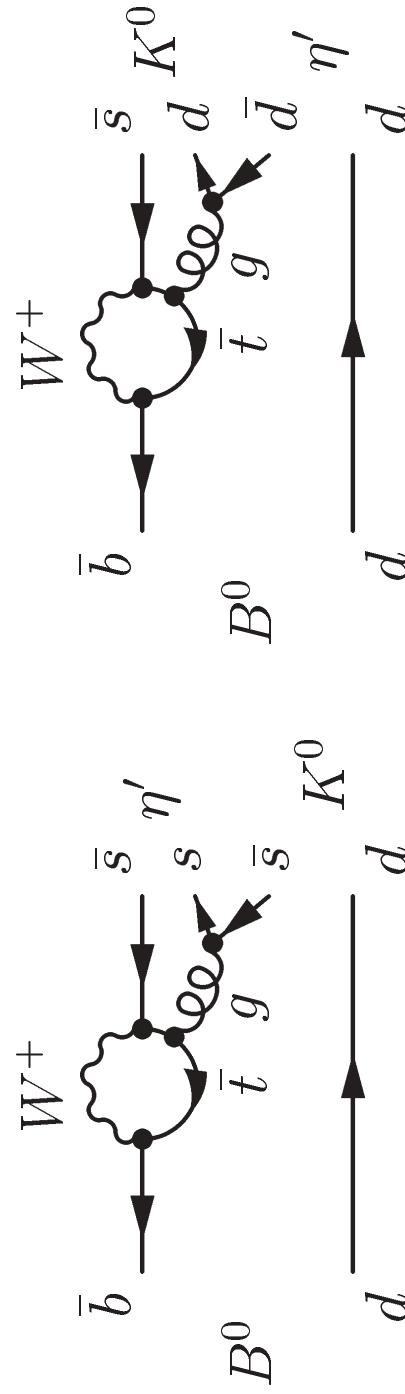
$$\sin 2\beta = -\xi_f S = 0.49 \pm 0.43 \pm 0.11^{+0.33}_{-0.00}$$

$$\mathcal{A} = -C = -0.40 \pm 0.33 \pm 0.10^{+0.00}_{-0.26}$$

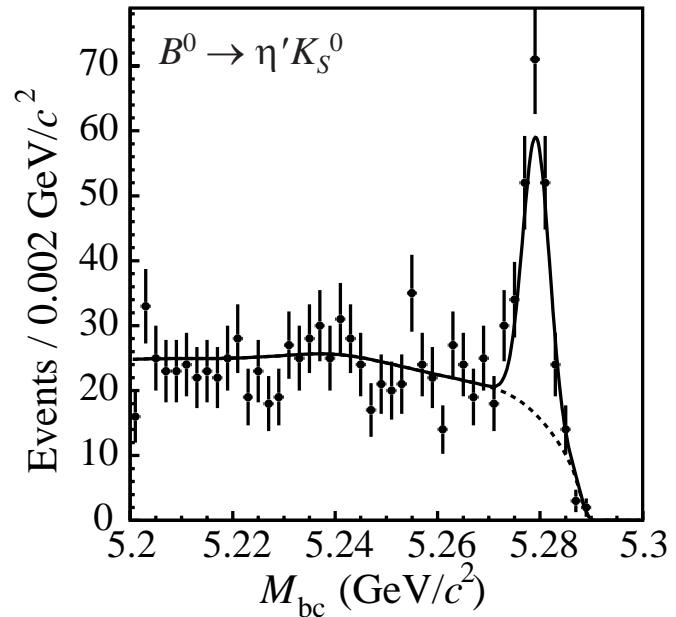


$$B^0 \rightarrow \eta' K_s^0$$

Again, gluonic penguins dominate:



Decay rate not well understood; a  $b \rightarrow u$  color-suppressed tree would bring in phase of  $V_{ub}$ , but with  $\mathcal{O}(\lambda^2)$  suppression.

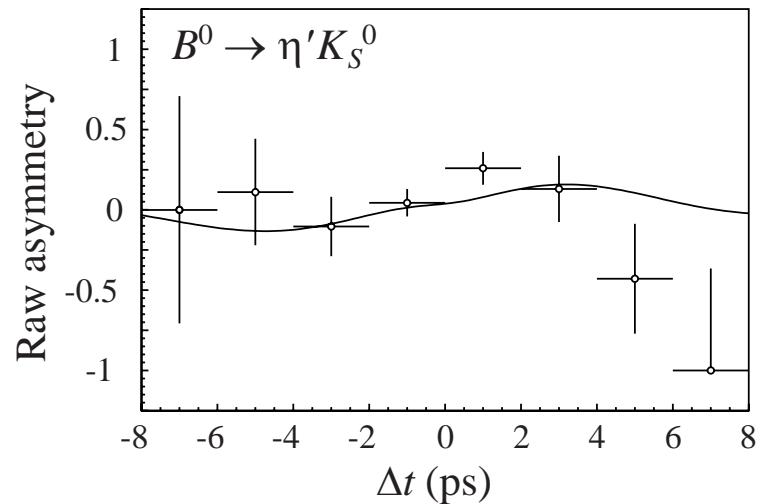
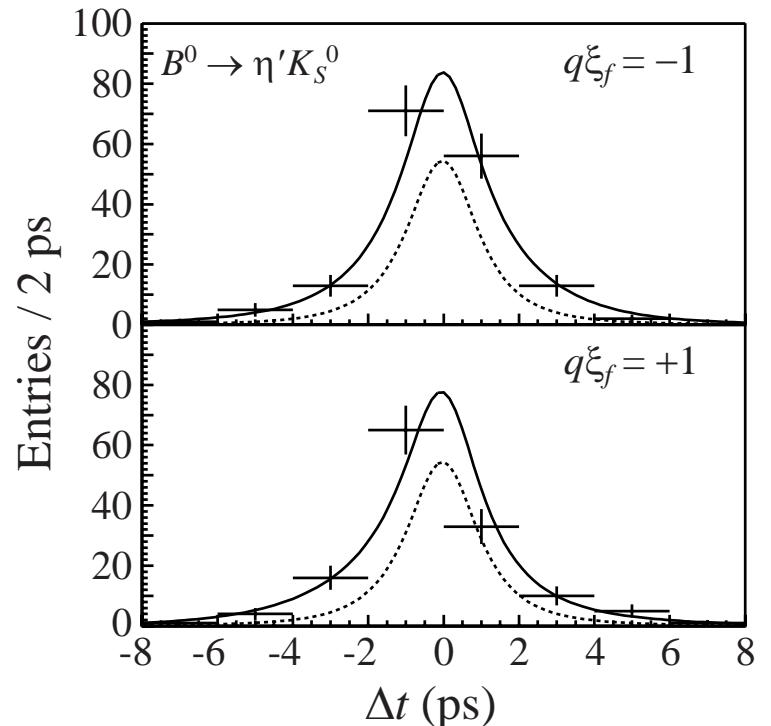


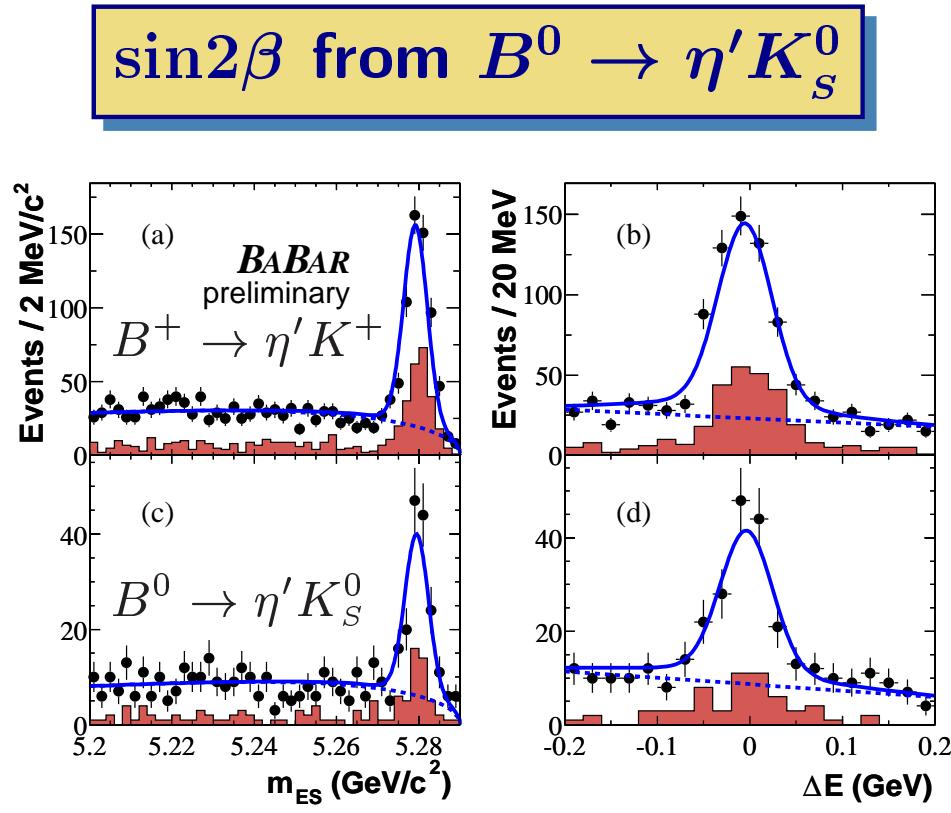
299 total events, purity 49%,  $78 \times 10^6 B\bar{B}$

$$\sin \phi_1 = S_{\eta' K_S^0} = +0.71 \pm 0.37^{+0.05}_{-0.06}$$

$$\mathcal{A} = -C_{\eta' K_S^0} = +0.26 \pm 0.22 \pm 0.03$$

(Consistent with charmonium result)

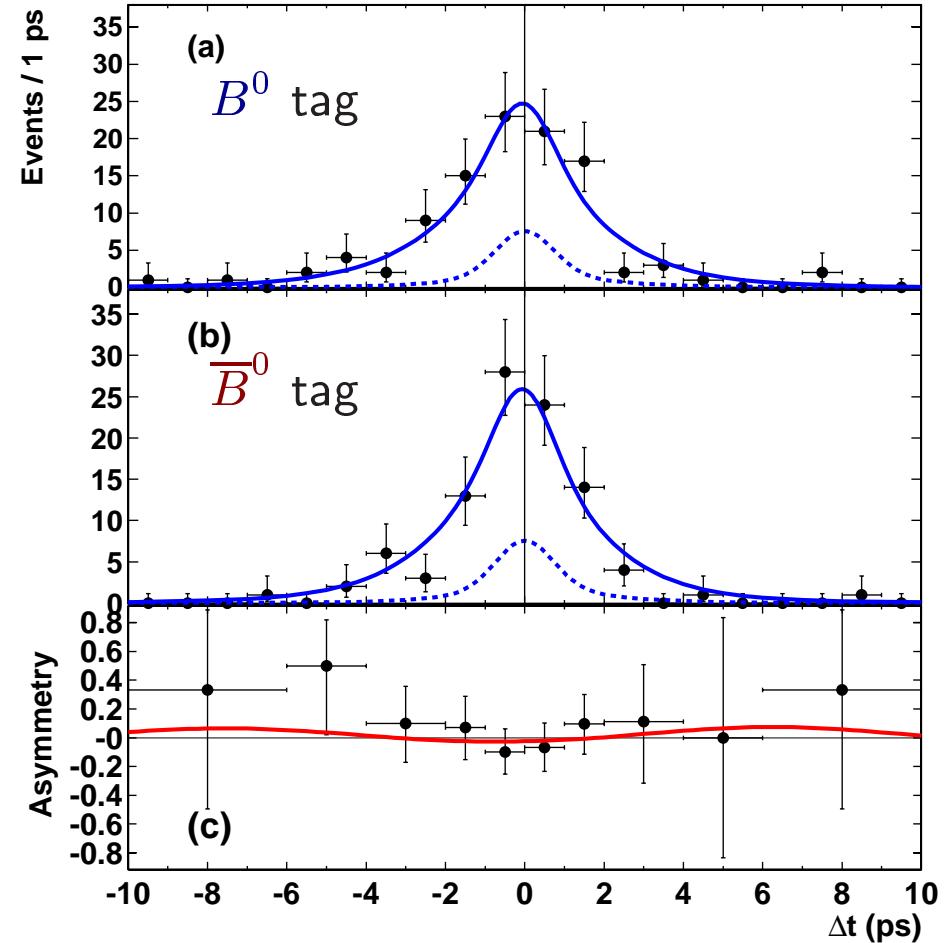




$\eta' \rightarrow \eta\pi\pi$  (shaded),  $\rho\gamma$

109.2 tagged signal events, purity 70%

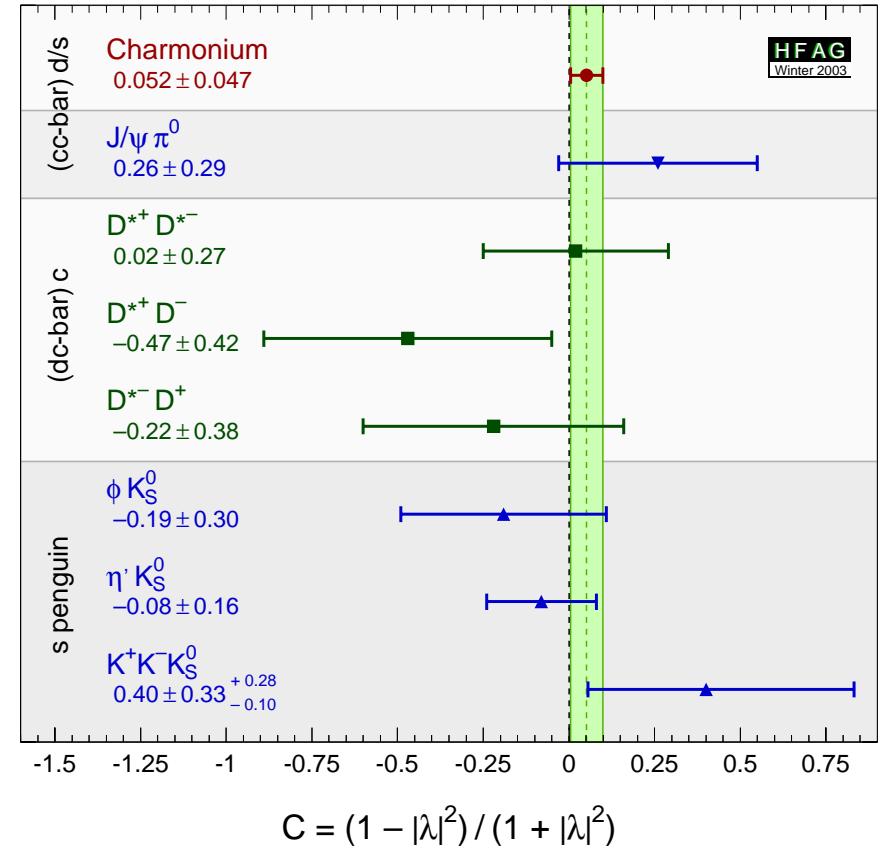
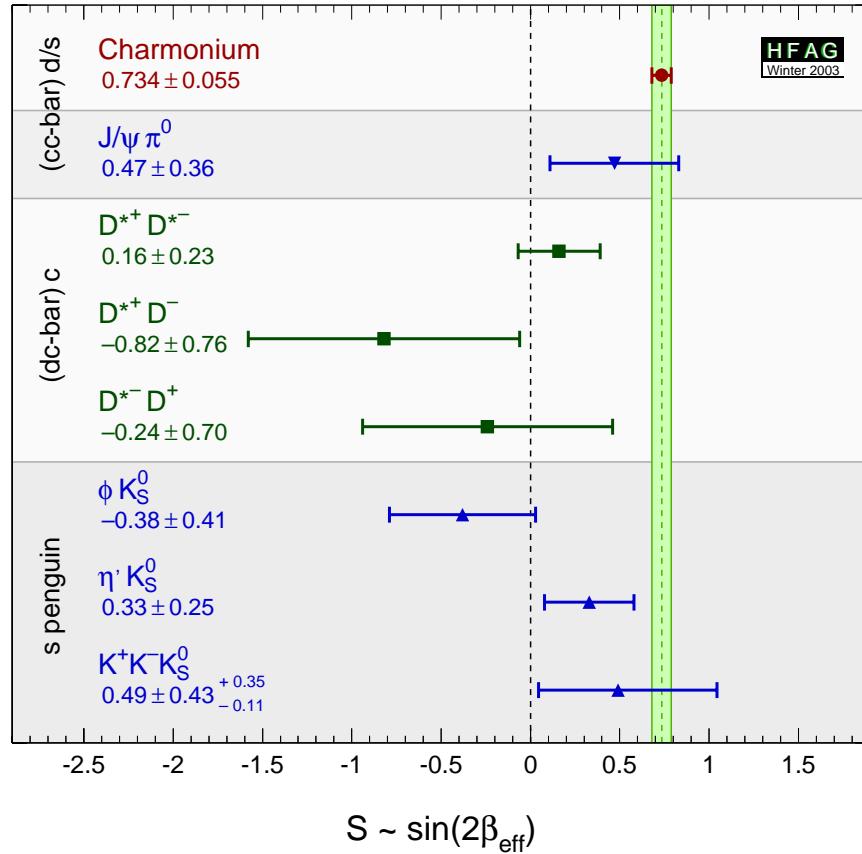
$89 \times 10^6 B\bar{B}$  pairs



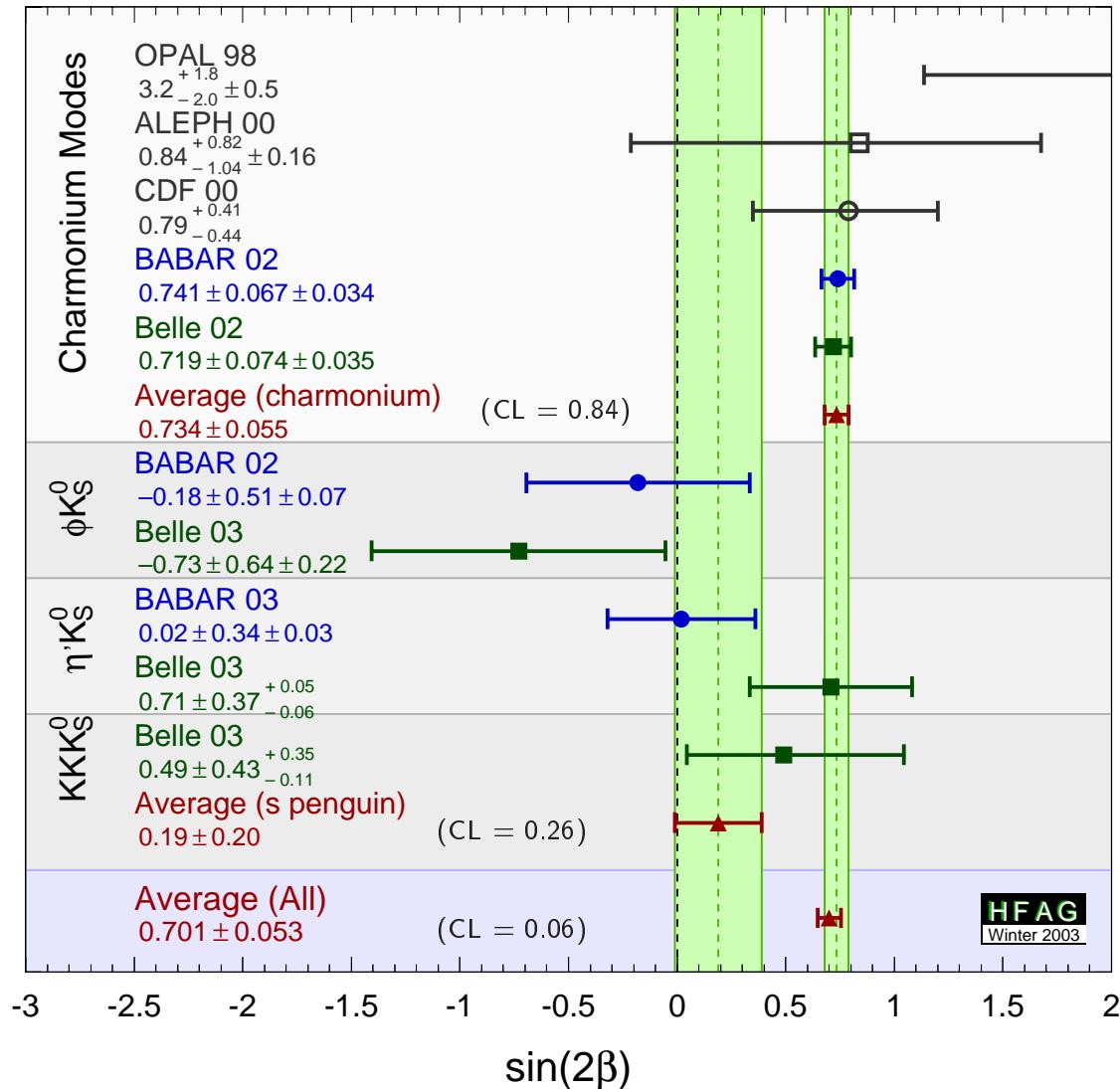
$$\sin 2\beta = S_{\eta' K_S^0} = 0.02 \pm 0.34 \pm 0.03, \quad C_{\eta' K_S^0} = 0.10 \pm 0.22 \pm 0.03$$

(Null result, though consistent within errors with  $\sin 2\beta = 0.7$ )

# World average measurements of $S, C$



## Measurements of $\sin 2\beta$



Future confirmation of the separation of the two bands would challenge the standard model.

Theoretical as well as experimental improvements needed.

## Conclusions

- $CP$  non conservation is well established in  $B^0$  decays
- The effect is large in the interference between mixing and decay
- The effect is well accommodated in the standard CKM model
- Just beginning to explore further
  - ◊ Anything new happening where S.M. effects are suppressed?
  - ◊ Hint of inconsistencies in  $\beta$  from  $\mathcal{O}(\lambda^3)$  decays
- Several channels now measured, needing only more data for definitive results